

# Chapter 5

## Symmetries and Conservation Laws

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# Parity Operator $\Pi$ - Overview

- Spatial reflection of the system  $\vec{r} \rightarrow -\vec{r}$ 
  - e.g. in 3D:  $x \rightarrow -x, \quad y \rightarrow -y, \quad z \rightarrow -z$
- Q.M. operators transform according to their (implicit) dependence on  $\vec{r}$ 
  - Odd under parity  $\rightarrow \Pi \hat{A} = -\hat{A}$
  - Even under parity  $\rightarrow \Pi \hat{A} = \hat{A}$
- Useful for solving matrix elements  $\langle \phi | \hat{A} | \psi \rangle$ 
  - If overall parity odd  $\rightarrow \langle \phi | \hat{A} | \psi \rangle = 0$
  - For sph. wave functions  $\rightarrow \Pi |l, m_l, s, m_s\rangle = (-1)^l |l, m_l, s, m_s\rangle$

# Parity Operator $\Pi$ - Examples

1.) Momentum operator

$$\hat{p} = -i\hbar\vec{\nabla} = -i\hbar\left(\frac{\partial}{\partial x}\hat{x} + \frac{\partial}{\partial y}\hat{y} + \frac{\partial}{\partial z}\hat{z}\right) \Rightarrow \Pi\hat{p} = -\hat{p} \quad \text{odd under parity}$$

$x \rightarrow -x \quad y \rightarrow -y \quad z \rightarrow -z$

2.) Angular momentum operator (pseudo vector)

$$\hat{L} = \hat{r} \times \hat{p} \quad \Rightarrow \quad \Pi\hat{L} = \hat{L} \quad \text{pseudo vectors are even under parity}$$

$\vec{r} \rightarrow -\vec{r} \quad \vec{p} \rightarrow -\vec{p}$

3.) Electromagnetic fields and potential

- pseudo vector  $\vec{B} = \vec{\nabla} \times \vec{A}$  is **even under parity**, therefore  $\vec{A}$  **odd under parity**
- Electric field  $\vec{E} = -\vec{\nabla}\Phi - \partial_t\vec{A}$  is **odd under parity**, therefore  $\Phi$  **even under parity**

# Parity Operator $\Pi$ - Examples

- 1.)  $\langle \alpha, l = 2, m = 1 | \hat{p}_x | \beta, l = 1, m = 0 \rangle \rightarrow$  overall even, might be non-zero  
           even                    odd            odd
  
- 2.)  $\langle \alpha, l = 2, m = 1 | \hat{x} \hat{p}_x | \beta, l = 1, m = 0 \rangle \rightarrow$  overall odd, matrix element zero  
           even                    odd/odd        odd  
   = even
  
- 3.)  $\langle \alpha, l = 3, m = 0 | \hat{z} \hat{L}_z | \beta, l = 1, m = 0 \rangle \rightarrow$  overall odd, matrix element zero  
           odd                    odd/even        odd  
   = odd

# Time Reversal Operator $\Theta$ - Overview

- Temporal reflection of the system  $t \rightarrow -t$ 
  - e.g.  $\vec{r} \rightarrow \vec{r}$ ,  $\vec{v} \rightarrow -\vec{v}$ ,  $\vec{E} \rightarrow \vec{E}$ ,  $\vec{B} \rightarrow -\vec{B}$ ,  $i\hbar\partial_t \rightarrow i\hbar\partial_t$
- Under time reversal, most operators are even or odd
  - For some operator,  $B$ :  $\Theta B \Theta^{-1} = \pm B$
- The time reversal operator takes the complex conjugate
  - Thus, it preserves the relation  $[J_i, J_j] = i\hbar\epsilon_{ijk}J_k$
- From this, we can say:
  - $\Theta J \Theta^{-1} = -J$

## Time Reversal Problem (Sakurai 4.12)

The Hamiltonian for a spin 1 system is given by:

$$H = AS_z^2 + B \left( S_x^2 - S_y^2 \right)$$

Solve this problem *exactly* to find the normalized energy eigenstates and eigenvalues. (A spin-dependent Hamiltonian of this kind actually appears in crystal physics.) Is this Hamiltonian invariant under time reversal? How do the normalized eigenstates you obtained transform under time reversal?

## Time Reversal Problem (cont.)

- First, we need to build the  $S_x$ ,  $S_y$ , and  $S_z$  matrices

$$S_z \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} = \hbar \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}$$

$$S_z \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} = 0$$

$$S_z \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} = -\hbar \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

$$S_z = \hbar \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix}$$

## Time Reversal Problem (cont.)

- Then, we build our  $S_+$  and  $S_-$  matrices from the following:

$$\begin{array}{l} S_+ \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} \quad S_+ \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} \quad S_+ \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} = 0 \\ S_- \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} \quad S_- \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} \quad S_- \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} = 0 \end{array}$$



## Time Reversal Problem (cont.)

- We can then use the relations  $S_x = \frac{1}{2}(S_+ + S_-)$  and  $S_y = \frac{1}{2i}(S_+ - S_-)$

$$S_+ = \hbar\sqrt{2} \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 0 & 0 & 0 \end{pmatrix}$$

$$S_- = \hbar\sqrt{2} \begin{pmatrix} 0 & 0 & 0 \\ 1 & 0 & 0 \\ 0 & 1 & 0 \end{pmatrix}$$

$$S_x = \frac{\hbar}{\sqrt{2}} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}$$

$$S_y = \frac{\hbar}{\sqrt{2}} \begin{pmatrix} 0 & -i & 0 \\ i & 0 & -i \\ 0 & i & 0 \end{pmatrix}$$

## Time Reversal Problem (cont.)

- We can now represent our Hamiltonian in matrix form:

$$H = AS_z^2 + B(S_x^2 - S_y^2)$$

$$H = A\hbar^2 \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix} + \frac{B}{2\hbar^2} \left( \begin{pmatrix} 1 & 0 & 1 \\ 0 & 2 & 0 \\ 1 & 0 & 1 \end{pmatrix} - \begin{pmatrix} 1 & 0 & -1 \\ 0 & 2 & 0 \\ -1 & 0 & 1 \end{pmatrix} \right) = \hbar^2 \begin{pmatrix} A & 0 & B \\ 0 & 0 & 0 \\ B & 0 & A \end{pmatrix}$$

## Time Reversal Problem (cont.)

- For eigenvalues and eigenvectors, we use  $\det(H-\lambda I) = 0$  to find:

$$\lambda_1 = A + B$$

$$\lambda_2 = 0$$

$$\lambda_3 = A - B$$

$$|\lambda_1 \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 0 \\ 1 \end{pmatrix}$$

$$|\lambda_2 \rangle = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}$$

$$|\lambda_3 \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 0 \\ -1 \end{pmatrix}$$

## Time Reversal Problem (cont.)

- Now, is the Hamiltonian invariant under time reversal?

$$\begin{aligned}\Theta H \Theta^{-1} &= A \Theta S_z^2 \Theta^{-1} + B (\Theta S_x^2 \Theta^{-1} + \Theta S_y^2 \Theta^{-1}) \\ &= A \Theta S_z \Theta^{-1} \Theta S_z \Theta^{-1} + B (\Theta S_x \Theta^{-1} \Theta S_x \Theta^{-1} + \Theta S_y \Theta^{-1} \Theta S_y \Theta^{-1}) \\ &= A (-S_z) (-S_z) + B ((-S_x) (-S_x) + (-S_y) (-S_y)) \\ &= A S_z^2 + B (S_x^2 + S_y^2)\end{aligned}$$

$$\Theta H \Theta^{-1} = H$$

## Time Reversal Problem (cont.)

- Are the eigenvectors invariant? First, rewrite them in the  $S_z$  basis

$$|j = 1, m = 1 \rangle = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}$$

$$|\lambda_1 \rangle = \frac{1}{2} (|j = 1, m = 1 \rangle + |j = 1, m = -1 \rangle)$$

$$|j = 1, m = 0 \rangle = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}$$

$$|\lambda_2 \rangle = |j = 1, m = 0 \rangle$$

$$|j = 1, m = -1 \rangle = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

$$|\lambda_3 \rangle = \frac{1}{2} (|j = 1, m = 1 \rangle - |j = 1, m = -1 \rangle)$$

## Time Reversal Problem (cont.)

- From Sakurai, for a spin-1 system,  $\Theta|j, m\rangle = (-1)^m|j, -m\rangle$

$$\Theta|\lambda_1\rangle = \frac{1}{2}(\Theta|1, 1\rangle + \Theta|1, -1\rangle) = \frac{1}{2}(-|1, -1\rangle - |1, 1\rangle)$$

$$\Theta|\lambda_1\rangle = -|\lambda_1\rangle$$

$$\Theta|\lambda_2\rangle = \Theta|1, 0\rangle = |1, 0\rangle$$

$$\Theta|\lambda_2\rangle = |\lambda_2\rangle$$

$$\Theta|\lambda_3\rangle = \frac{1}{2}(\Theta|1, 1\rangle - \Theta|1, -1\rangle) = \frac{1}{2}(-|1, -1\rangle + |1, 1\rangle)$$

$$\Theta|\lambda_3\rangle = |\lambda_3\rangle$$