

# Chapter 2 Review

Adams and Yeung

# Topics Covered in Chapter

Wave Equation

Momentum Space

Charge conservation and Continuity

Potential Problems in 1D

Harmonic Oscillator

# Past Problems

Not exactly (it's too easy!), but content shows up mostly in:

- Spherically symmetric bound state problems (in combination with chapter 4)
  - August 2021 problem 4, August 2020 problem 1, Spring 2020 problem 4 part d
- Harmonic oscillators frequently show up in Fermi's golden rule problems
- invariant

## Bound states

4. Consider a particle of mass  $m$  in a three-dimensional potential

$$V(\mathbf{r}) = -\beta\delta(\mathbf{r} - \mathbf{a}), \quad \beta > 0.$$

(a) (10 pts) In terms of  $\beta$  and  $m$ , what is the minimum value of  $a$  for which one has a bound state?

This gives a Schrödinger equation

$$\begin{aligned} -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial r^2} \psi(r) + V(r)\psi(r) &= E\psi(r) \\ -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial r^2} \psi(r) - \beta\delta(r - a)\psi(r) &= E\psi(r) \end{aligned}$$

## Bound states

$$\begin{aligned}\frac{\partial^2}{\partial r^2} u(r \neq a) &= -\frac{2mE}{\hbar^2} u(r \neq a) \\ &= \frac{2m|E|}{\hbar^2} u(r \neq a)\end{aligned}$$

This is positive because, for a bound state, the energy is negative. This implies the wavefunctions are exponentials. We can also tell that  $k$  is the same on both sides.

$$\begin{aligned}-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial r^2} \psi(r) - \beta \delta(r - a) \psi(r) &= E \psi(r) \\ \frac{\partial^2}{\partial r^2} \psi(r) &= -\frac{2m}{\hbar^2} (E + \beta \delta(r - a)) \psi(r)\end{aligned}$$

Integrating this over the interval  $[a - \varepsilon, a + \varepsilon]$  ( $\varepsilon \rightarrow 0$ ), we get the boundary condition

$$\left. \frac{\partial \psi}{\partial r} \right|_{r=a^+} - \left. \frac{\partial \psi}{\partial r} \right|_{r=a^-} = -\frac{2m\beta}{\hbar^2} u(a)$$

# Bound states

We also know that the wavefunction must be continuous across the boundary. Furthermore,

$$\begin{aligned}\psi(0) &= 0 \\ \psi(r \rightarrow \infty) &= 0\end{aligned}$$

From these boundary conditions, we can derive

$$\begin{aligned}\psi(r < a) &= A \sinh kr \\ \psi(r > a) &= B e^{-kr}\end{aligned}$$

# Bound states

At  $r = a$ , the wavefunction and its slope give, respectively,

$$A \sinh ka = B e^{-ka}$$

$$-k B e^{-ka} - k A \cosh ka = -\frac{2m\beta}{\hbar^2} A \sinh ka$$

$$-k A \sinh ka - k A \cosh ka = -\frac{2m\beta}{\hbar^2} A \sinh ka$$

$$\sinh ka + \cosh ka = \frac{2m\beta}{\hbar^2} \frac{\sinh ka}{k}$$

# Bound states

For the minimum number of states,  $k \rightarrow 0$  as we want the wavefunction to be as evenly spread throughout the space as possible.

$$\begin{aligned}\lim_{k \rightarrow 0} (\sinh ka + \cosh ka) &= \lim_{k \rightarrow 0} \frac{2m\beta a \sinh ka}{\hbar^2 ka} \\ 1 &= \frac{2m\beta a}{\hbar^2} \\ a &= \frac{\hbar^2}{2m\beta}\end{aligned}$$



# Aug 2020 Problem 1

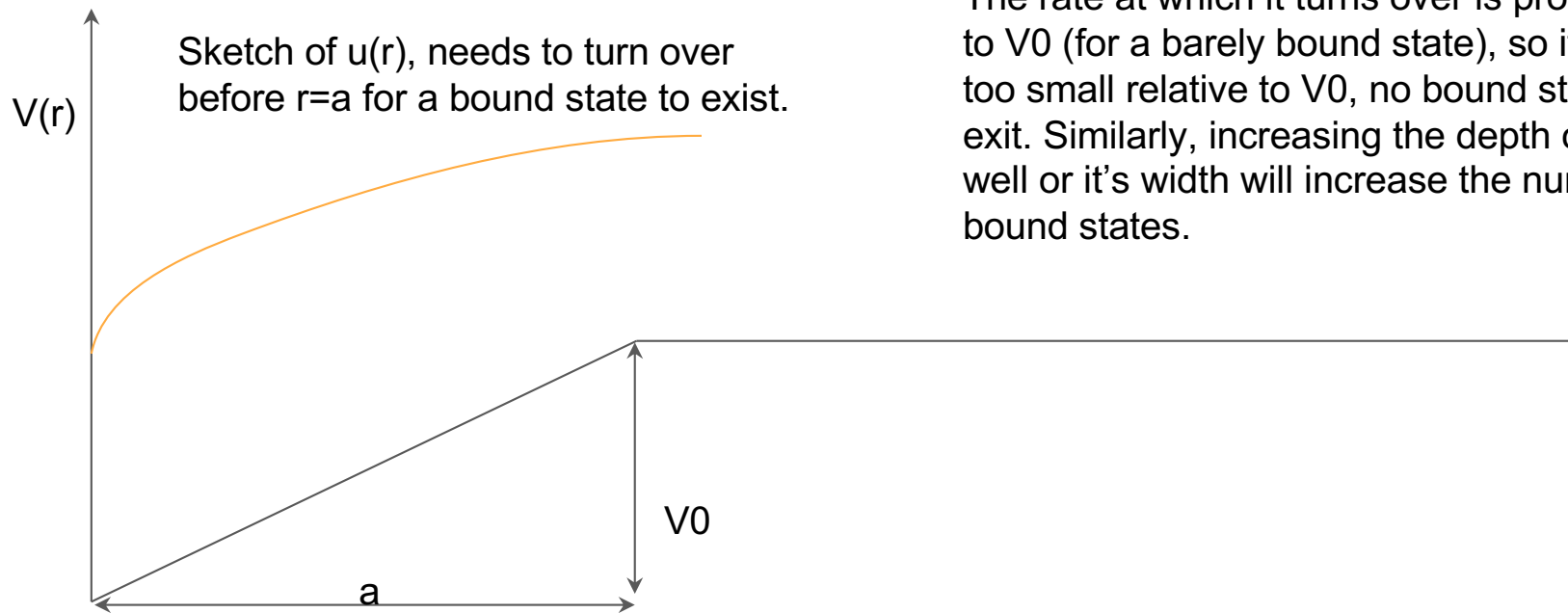
1. (10 pts) A particle of mass  $m$  interacts with a spherically symmetric attractive potential,

$$V(r) = \begin{cases} -V_0 + \frac{V_0 r}{a}, & r < a \\ 0, & r > a \end{cases}$$

Circle the true statements below:

- Fixing  $a$ , and making  $V_0$  very small, but non-zero, there will always be at least one bound state.
- Fixing  $V_0$ , and making  $a$  very small, but non-zero, there will always be at least one bound state.
- Fixing  $a$ , as the magnitude of  $V_0$  increases, the number of bound states will increase.
- Fixing  $V_0$ , as the magnitude of  $a$  increases, the number of bound states will increase.

# August 2020 Problem 1: Solution



The rate at which it turns over is proportional to  $V_0$  (for a barely bound state), so if “ $a$ ” is too small relative to  $V_0$ , no bound state will exist. Similarly, increasing the depth of the well or its width will increase the number of bound states.

# Harmonic Oscillator Overlap

Fermi's golden rule problems sometimes require calculating matrix elements such as  $\langle 1|x|0 \rangle$  where these are the ground and first excited states of the harmonic oscillator. In these cases, you can avoid integrating by using  $x = \sqrt{\frac{\hbar}{2m\omega}} (a + a^\dagger)$   
Then:

$$\begin{aligned}\langle 1|x|0 \rangle &= \langle 1|\sqrt{\frac{\hbar}{2m\omega}} (a + a^\dagger) |0 \rangle \\ &= \sqrt{\frac{\hbar}{2m\omega}} \langle 1|a^\dagger|0 \rangle = \sqrt{\frac{\hbar}{2m\omega}}\end{aligned}$$