**Problem 4.8a.** Assuming that the Hamiltonian is invariant under time reversal, prove that the wave function for a spinless nondegenerate system at any given instant of time can always be chosen to be real.

Call  $|n\rangle$  the eigenstates of H with energy  $E_n$ . Then using that H is invariant under  $\Theta$ ,

$$H\Theta \left| n \right\rangle = \Theta H \left| n \right\rangle = E_n \Theta \left| n \right\rangle,$$

meaning that the energy of  $\Theta |n\rangle$  is  $E_n$ . We assumed that the system is nondegenerate, so this means  $\Theta |n\rangle$  and  $|n\rangle$  are the same state, differing by at most a phase independent of **r**. If  $\Theta |n\rangle = e^{i\delta} |n\rangle$ , we can redefine the energy eigenstates  $|\tilde{n}\rangle \equiv e^{i\delta/2} |n\rangle$  so that

$$\begin{split} \Theta \left| \tilde{n} \right\rangle &= \Theta(e^{i\delta/2} \left| n \right\rangle) \\ &= e^{-i\delta/2} e^{i\delta} \left| n \right\rangle \\ &= \left| \tilde{n} \right\rangle. \end{split}$$

Recall that

$$\Theta \psi_{\tilde{n}}(\mathbf{r}) = \psi_{\tilde{n}}^{*}(\mathbf{r})$$
$$\Rightarrow \Theta \langle \mathbf{r} | \tilde{n} \rangle = \langle \mathbf{r} | \tilde{n} \rangle^{*},$$

but we just showed that

$$\Theta \left< \mathbf{r} | \tilde{n} \right> = \left< \mathbf{r} | \tilde{n} \right>,$$

 $\mathbf{so}$ 

$$\langle \mathbf{r} | \tilde{n} \rangle = \langle \mathbf{r} | \tilde{n} \rangle^*,$$

therefore using this phase convention  $\psi_{\tilde{n}}(\mathbf{r})$  is real.

**Problem 4.8b.** The wave function for a plane-wave state at t = 0 is given by a complex function  $e^{i\mathbf{p}\cdot\mathbf{r}/\hbar}$ . Why does this not violate time-reversal invariance?

The statement in (a) doesn't apply here because  $e^{i\mathbf{p}\cdot\mathbf{r}/\hbar}$  is a degenerate state with  $e^{-i\mathbf{p}\cdot\mathbf{r}/\hbar}$ .

**Problem 4.9.** Let  $\phi(\mathbf{p})$  be the momentum-space wave function for state  $|\alpha\rangle$ , that is,  $\phi(\mathbf{p}) = \langle \mathbf{p} | \alpha \rangle$ . Is the momentum-space wave function for the time-reversed state  $\Theta | \alpha \rangle$  given by  $\phi(\mathbf{p}), \phi(-\mathbf{p}), \phi^*(\mathbf{p}), \text{ or } \phi^*(-\mathbf{p})$ ? Justify your answer.

First a few background tidbits: From class notes, we know  $\Theta \mathbf{P} \Theta^{-1} = -\mathbf{P}$ , in other words,  $\mathbf{P}$  is odd under time reversal. When we apply this to a state,

$$\mathbf{P}\Theta \left| \mathbf{p} \right\rangle = -\Theta \mathbf{P}\Theta^{-1}\Theta \left| \mathbf{p} \right\rangle = (-\mathbf{p})\Theta \left| \mathbf{p} \right\rangle$$

We see that  $\Theta |\mathbf{p}\rangle$  is the momentum eigenstate corresponding to eigenvalue  $-\mathbf{p}$ . Now on to solving the problem at hand:

We express  $\alpha$  as

$$|\alpha\rangle = \int \langle \mathbf{p} | \alpha \rangle^* | \mathbf{p} \rangle d^3 p.$$

So then let us write  $\Theta \left| \alpha \right\rangle$  as

$$\begin{split} \Theta \left| \alpha \right\rangle &= \int \Theta \left| \mathbf{p} \right\rangle \left\langle \mathbf{p} \right| \alpha \right\rangle^* d^3 p \\ &= \int \left| -\mathbf{p} \right\rangle \left\langle \mathbf{p} \right| \alpha \right\rangle d^3 p \end{split}$$

Theres nothing to keep us from moving the minus sign:

$$egin{aligned} \Theta \left| lpha 
ight
angle &= \int \left| \mathbf{p} 
ight
angle \left\langle -\mathbf{p} 
ight| lpha 
ight
angle^{st} d^{3}p \ &= \int \phi^{st}(-\mathbf{p}) \left| \mathbf{p} 
ight
angle d^{3}p, \end{aligned}$$

So we see that

$$\Theta\phi(\mathbf{p}) = \phi^*(-\mathbf{p}).$$

Does this make sense? Yes. Because in momentum-space time reversal means motion reversal; its intuitive that  $\mathbf{p}$  picks up the negative sign, and we know from the class notes that the time reversal operator involves taking the complex conjugate, as we see in our answer.