1. Do not write your name on the exam. Your exam has a "secret student number" on each page. If you include any pages beyond those included with the exam, be sure to write that number on EACH additional page of the solutions you submit.
2. This exam is closed-book, closed-mouth, closed-notes and closed-internet. You are not permitted to use mathematical software, e.g. mathematica. You are not to communicate with any other individuals during the exam, aside from the proctor.

$$
\begin{aligned}
& \int_{-\infty}^{\infty} d x e^{-x^{2} /\left(2 a^{2}\right)}=a \sqrt{2 \pi}, \\
& \boldsymbol{H}=i \hbar \boldsymbol{\partial}_{t}, \overrightarrow{\boldsymbol{P}}=-i \hbar \boldsymbol{\nabla}, \\
& \sigma_{z}=\left(\begin{array}{ll}
1 & 0 \\
0 & -1
\end{array}\right), \sigma_{x}=\left(\begin{array}{ll}
0 & 1 \\
1 & 0
\end{array}\right), \quad \sigma_{y}=\left(\begin{array}{ll}
0 & -i \\
i & 0
\end{array}\right), \\
& U(t,-\infty)=1+\frac{-i}{\hbar} \int_{-\infty}^{t} d t^{\prime} V\left(t^{\prime}\right) U\left(t^{\prime},-\infty\right), \\
& \left\langle x \mid x^{\prime}\right\rangle=\delta\left(x-x^{\prime}\right),\left\langle p \mid p^{\prime}\right\rangle=\frac{1}{2 \pi \hbar} \delta\left(p-p^{\prime}\right), \\
& |p\rangle=\int d x|x\rangle e^{i p x / \hbar}, \quad|x\rangle=\int \frac{d p}{2 \pi \hbar}|p\rangle e^{-i p x / \hbar}, \\
& H=\frac{P^{2}}{2 m}+\frac{1}{2} m \omega^{2} x^{2}=\hbar \omega\left(a^{\dagger} a+1 / 2\right), \\
& a^{\dagger}=\sqrt{\frac{m \omega}{2 \hbar}} X-i \sqrt{\frac{1}{2 \hbar m \omega}} P, \\
& \psi_{0}(x)=\frac{1}{\left(\pi b^{2}\right)^{1 / 4}} e^{-x^{2} / 2 b^{2}}, \quad b^{2}=\frac{\hbar}{m \omega}, \\
& \rho(\vec{r}, t)=\psi^{*}\left(\vec{r}_{1}, t_{1}\right) \psi\left(\vec{r}_{2}, t_{2}\right) \\
& \vec{j}(\vec{r}, t)=\frac{-i \hbar}{2 m}\left(\psi^{*}(\vec{r}, t) \nabla \psi(\vec{r}, t)-\left(\nabla \psi^{*}(\vec{r}, t)\right) \psi(\vec{r}, t)\right) \\
& -\frac{e \vec{A}}{m c}|\psi(\vec{r}, t)|^{2} . \\
& H=\frac{(\vec{P}-e \vec{A} / c)^{2}}{2 m}+e \Phi, \\
& \text { For } V=\beta \delta(x-y): \quad-\frac{\hbar^{2}}{2 m}\left(\left.\frac{\partial}{\partial x} \psi(x)\right|_{y+\epsilon}-\left.\frac{\partial}{\partial x} \psi(x)\right|_{y-\epsilon}\right)=-\beta \psi(y) \text {, } \\
& \vec{E}=-\nabla \Phi-\frac{1}{c} \partial_{t} \vec{A}, \quad \vec{B}=\nabla \times \vec{A}, \\
& \omega_{\text {cyclotron }}=\frac{e B}{m c}, \\
& e^{A+B}=e^{A} e^{B} e^{-C / 2}, \quad \text { if }[A, B]=C, \text { and }[C, A]=[C, B]=0, \\
& Y_{0,0}=\frac{1}{\sqrt{4 \pi}}, \quad Y_{1,0}=\sqrt{\frac{3}{4 \pi}} \cos \theta, \quad Y_{1, \pm 1}=\mp \sqrt{\frac{3}{8 \pi}} \sin \theta e^{i \pm \phi}, \\
& Y_{2,0}=\sqrt{\frac{5}{16 \pi}}\left(3 \cos ^{2} \theta-1\right), \quad Y_{2, \pm 1}=\mp \sqrt{\frac{15}{8 \pi}} \sin \theta \cos \theta e^{ \pm i \phi}, \\
& Y_{2, \pm 2}=\sqrt{\frac{15}{32 \pi}} \sin ^{2} \theta e^{ \pm 2 i \phi}, \quad Y_{\ell-m}(\theta, \phi)=(-1)^{m} Y_{\ell m}^{*}(\theta, \phi) .
\end{aligned}
$$

$$
\begin{aligned}
& |N\rangle=|n\rangle-\sum_{m \neq n}|m\rangle \frac{1}{\epsilon_{m}-\epsilon_{n}}\langle m| V|n\rangle+\cdots \\
& \boldsymbol{E}_{N}=\epsilon_{n}+\langle n| V|n\rangle-\sum_{m \neq n} \frac{|\langle m| V| n\rangle\left.\right|^{2}}{\epsilon_{m}-\epsilon_{n}} \\
& j_{0}(x)=\frac{\sin x}{x}, n_{0}(x)=-\frac{\cos x}{x}, j_{1}(x)=\frac{\sin x}{x^{2}}-\frac{\cos x}{x}, n_{1}(x)=-\frac{\cos x}{x^{2}}-\frac{\sin x}{x} \\
& j_{2}(x)=\left(\frac{3}{x^{3}}-\frac{1}{x}\right) \sin x-\frac{3}{x^{2}} \cos x, n_{2}(x)=-\left(\frac{3}{x^{3}}-\frac{1}{x}\right) \cos x-\frac{3}{x^{2}} \sin x, \\
& \frac{d}{d t} P_{i \rightarrow n}(t)=\frac{2 \pi}{\hbar}\left|V_{n i}\right|^{2} \delta\left(E_{n}-E_{i}\right), \\
& \frac{d \sigma}{d \Omega}=\frac{m^{2}}{4 \pi^{2} \hbar^{4}}\left|\int d^{3} r \mathcal{V}(r) e^{i\left(\vec{k}_{f}-\vec{k}_{i}\right) \cdot \vec{r}}\right|^{2}, \\
& \sigma=\frac{\left(2 S_{R}+1\right)}{\left(2 S_{1}+1\right)\left(2 S_{2}+1\right)} \frac{4 \pi}{k^{2}} \frac{\left(\hbar \Gamma_{R} / 2\right)^{2}}{\left(\epsilon_{k}-\epsilon_{r}\right)^{2}+\left(\hbar \Gamma_{R} / 2\right)^{2}}, \\
& \frac{d \sigma}{d \Omega}=\left(\frac{d \sigma}{d \Omega}\right)_{\text {single }} \tilde{S}(\vec{q}), \quad \tilde{S}(\vec{q})=\left|\sum_{\vec{a}} e^{i \vec{q} \cdot \vec{a}}\right|^{2}, \\
& e^{i \vec{k} \cdot \vec{r}}=\sum_{\ell}(2 \ell+1) i^{\ell} j_{\ell}(k r) P_{\ell}(\cos \theta), \\
& P_{\ell}(\cos \theta)=\sqrt{\frac{4 \pi}{2 \ell+1}} Y_{\ell, m=0}(\theta, \phi), \\
& P_{0}(x)=1, P_{1}(x)=x, P_{2}(x)=\left(3 x^{2}-1\right) / 3, \\
& f(\Omega) \equiv \sum_{\ell}(2 \ell+1) e^{i \delta_{\ell}} \sin \delta_{\ell} \frac{1}{k} P_{\ell}(\cos \theta) \\
& \left.\psi_{\vec{k}}(\vec{r})\right|_{R \rightarrow \infty}=e^{i \vec{k} \cdot \vec{r}}+\frac{e^{i k r}}{r} f(\Omega), \\
& \frac{d \sigma}{d \Omega}=|f(\Omega)|^{2}, \quad \sigma=\frac{4 \pi}{k^{2}} \sum_{\ell}(2 \ell+1) \sin ^{2} \delta_{\ell}, \quad \delta \approx-a k \\
& L_{ \pm}|\ell, m\rangle=\sqrt{\ell(\ell+1)-m(m \pm 1)}|\ell, m \pm 1\rangle, \\
& C_{m_{\ell}, m_{s} ; J M}^{\ell, s}=\left\langle\ell, s, J, M \mid \ell, s, m_{\ell}, m_{s}\right\rangle, \\
& \langle\tilde{\beta}, J, M| T_{q}^{k}\left|\beta, \ell, m_{\ell}\right\rangle=C_{q m_{\ell} ; J M}^{k \ell} \frac{\langle\tilde{\beta}, J|\left|T^{(k)}\right||\beta, \ell, J\rangle}{\sqrt{2 J+1}}, \\
& n=\frac{(2 s+1)}{(2 \pi)^{d}} \int_{k<k_{f}} d^{d} k, \quad d \text { dimensions }, \\
& \left\{\Psi_{s}(\vec{x}), \Psi_{s^{\prime}}^{\dagger}(\vec{y})\right\}=\delta^{3}(\vec{x}-\vec{y}) \delta_{s s^{\prime}}, \\
& \Psi_{s}^{\dagger}(\vec{r})=\frac{1}{\sqrt{V}} \sum_{\vec{k}} e^{i \vec{k} \cdot \vec{r}} a_{s}^{\dagger}(\vec{k}), \quad\left\{\Psi_{s}(\vec{x}), a_{\alpha}^{\dagger}\right\}=\phi_{\alpha, s}(\vec{x}) .
\end{aligned}
$$

1. Consider a two-component system with the original Hamiltonian,

$$
H_{0}=A \sigma_{z}, \quad A>0
$$

A time $\boldsymbol{t}=\mathbf{0}$ one adds an additional potential

$$
\boldsymbol{V}=\boldsymbol{g} \sigma_{x}
$$

(a) (10 pts) If at time $\boldsymbol{t}=\mathbf{0}$ the system is in the ground state of $\boldsymbol{H}_{\mathbf{0}}$, find the probability for being in that original ground state as a function of time.
(b) (10 pts) To second order in perturbation theory (where $\boldsymbol{V}$ is the perturbation), what is the ground state energy?
(c) (10 pts) What is the exact solution for the eigen-energies of the full Hamiltonian?
2. Consider the states

$$
\begin{aligned}
|\boldsymbol{\eta}\rangle & =\exp \left\{\boldsymbol{i}\left(\boldsymbol{\eta} \boldsymbol{a}^{\dagger}+\boldsymbol{\eta}^{*} \boldsymbol{a}\right)\right\}|0\rangle \\
|\gamma\rangle & =\exp \left\{\boldsymbol{i}\left(\gamma \boldsymbol{a}^{\dagger}+\boldsymbol{\gamma}^{*} \boldsymbol{a}\right)\right\}|0\rangle
\end{aligned}
$$

Here, $\boldsymbol{\eta}$ and $\boldsymbol{\gamma}$ are complex numbers, and $\boldsymbol{a}^{\dagger}$ and $\boldsymbol{a}$ are creation and annihilation operators respectively.
(a) (10 pts) Calculate $\langle\boldsymbol{\eta} \mid \boldsymbol{\eta}\rangle$.
(b) (10 pts) Calculate $\langle\boldsymbol{\gamma} \mid \boldsymbol{\eta}\rangle$.
3. (20 pts) Two point charges are positioned in a line along the $\boldsymbol{z}$ axis, with a distance $\boldsymbol{a}$ separating the charges. A beam with momentum $\boldsymbol{p}$ is incident on the charges along the $\boldsymbol{z}$ axis. The scattering can be considered as a perturbative process. In order to determine the distance $\boldsymbol{a}$, you measure the directions at which the differential cross section is the smallest. In terms of $\boldsymbol{a}$ and $\boldsymbol{p}$, list the angles for which all the scattering is smallest. Use $\boldsymbol{\theta}$ for the polar angle, the angle relative to the $\boldsymbol{z}$ axis, and $\boldsymbol{\phi}$ for the azimuthal angle.

Extra space for \#3
4. Consider a particle of mass $\boldsymbol{m}$ in a three-dimensional potential

$$
V(r)=-\beta \delta(r-a), \quad \beta>0
$$

(a) (10 pts) In terms of $\boldsymbol{\beta}$ and $\boldsymbol{m}$, what is the minimum value of $\boldsymbol{a}$ for which one has a bound state?
(b) (10 pts) Assuming $\boldsymbol{\beta}$ is above the value above, what is the $\boldsymbol{s}$-wave phase shift as a function of the magnitude of the momentum, $\hbar \boldsymbol{k}$ ?

Extra space for \#4
5. (30 pts) Consider a ONE-DIMENSIONAL world, where a non-relativistic particle of mass $\boldsymbol{M}$ is in the ground state of a harmonic oscillator characterized by frequency $\boldsymbol{\omega}_{\mathbf{0}}$. These massive particles are created and destroyed with field operators, $\boldsymbol{\Psi}(\boldsymbol{x})$. These particles, referred to as $\Psi$ particles, can transform into $\boldsymbol{\Phi}$ particles, which have exactly the same mass but differ in that they do not feel the harmonic oscillator potential. The $\boldsymbol{\Phi}$ particles are created and destroyed with field operators $\boldsymbol{\Phi}(\boldsymbol{x})$. The $\boldsymbol{\Psi}(\boldsymbol{x})$ and $\boldsymbol{\Phi}(\boldsymbol{x})$ field operators obey the commutation relations

$$
\left[\Psi(x), \Psi^{\dagger}(y)\right]=\delta(x-y), \quad\left[\Phi(x), \Phi^{\dagger}(y)\right]=\delta(x-y)
$$

The perturbative interaction responsible for the transformation is

$$
V=g \int d x\left(\Psi^{\dagger}(x) \Phi(x)+\Phi^{\dagger}(x) \Psi(x)\right)
$$

Using Fermi's golden rule, calculate the rate of decay of $\boldsymbol{\Psi}$ particles into $\boldsymbol{\Phi}$ particles.

Extra space for \#5

