

$$\int_{-\infty}^{\infty} dx e^{-x^2/(2a^2)} = a\sqrt{2\pi},$$

$$H = i\hbar\partial_t, \vec{P} = -i\hbar\nabla,$$

$$\sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix},$$

$$U(t, -\infty) = 1 + \frac{-i}{\hbar} \int_{-\infty}^t dt' V(t')U(t', -\infty),$$

$$\langle x|x'\rangle = \delta(x - x'), \langle p|p'\rangle = \frac{1}{2\pi\hbar}\delta(p - p'),$$

$$|p\rangle = \int dx |x\rangle e^{ipx/\hbar}, |x\rangle = \int \frac{dp}{2\pi\hbar} |p\rangle e^{-ipx/\hbar},$$

$$H = \frac{P^2}{2m} + \frac{1}{2}m\omega^2 x^2 = \hbar\omega(a^\dagger a + 1/2),$$

$$a^\dagger = \sqrt{\frac{m\omega}{2\hbar}}X - i\sqrt{\frac{1}{2\hbar m\omega}}P,$$

$$\psi_0(x) = \frac{1}{(\pi b^2)^{1/4}} e^{-x^2/2b^2}, \quad b^2 = \frac{\hbar}{m\omega},$$

$$\rho(\vec{r}, t) = \psi^*(\vec{r}_1, t_1)\psi(\vec{r}_2, t_2)$$

$$\vec{j}(\vec{r}, t) = \frac{-i\hbar}{2m}(\psi^*(\vec{r}, t)\nabla\psi(\vec{r}, t) - (\nabla\psi^*(\vec{r}, t))\psi(\vec{r}, t)) \\ - \frac{e\vec{A}}{mc}|\psi(\vec{r}, t)|^2.$$

$$H = \frac{(\vec{P} - e\vec{A}/c)^2}{2m} + e\Phi,$$

For  $V = \beta\delta(x - y)$ :  $-\frac{\hbar^2}{2m}\left(\frac{\partial}{\partial x}\psi(x)|_{y+\epsilon} - \frac{\partial}{\partial x}\psi(x)|_{y-\epsilon}\right) = -\beta\psi(y),$

$$\vec{E} = -\nabla\Phi - \frac{1}{c}\partial_t\vec{A}, \quad \vec{B} = \nabla \times \vec{A},$$

$$\omega_{\text{cyclotron}} = \frac{eB}{mc},$$

$$e^{A+B} = e^A e^B e^{-C/2}, \quad \text{if } [A, B] = C, \text{ and } [C, A] = [C, B] = 0,$$

$$Y_{0,0} = \frac{1}{\sqrt{4\pi}}, \quad Y_{1,0} = \sqrt{\frac{3}{4\pi}}\cos\theta, \quad Y_{1,\pm 1} = \mp\sqrt{\frac{3}{8\pi}}\sin\theta e^{\pm i\phi},$$

$$Y_{2,0} = \sqrt{\frac{5}{16\pi}}(3\cos^2\theta - 1), \quad Y_{2,\pm 1} = \mp\sqrt{\frac{15}{8\pi}}\sin\theta\cos\theta e^{\pm i\phi},$$

$$Y_{2,\pm 2} = \sqrt{\frac{15}{32\pi}}\sin^2\theta e^{\pm 2i\phi}, \quad Y_{\ell-m}(\theta, \phi) = (-1)^m Y_{\ell m}^*(\theta, \phi).$$

$$|N\rangle = |n\rangle - \sum_{m \neq n} |m\rangle \frac{1}{\epsilon_m - \epsilon_n} \langle m|V|n\rangle + \dots$$

$$E_N = \epsilon_n + \langle n|V|n\rangle - \sum_{m \neq n} \frac{|\langle m|V|n\rangle|^2}{\epsilon_m - \epsilon_n}$$

$$j_0(x) = \frac{\sin x}{x}, \quad n_0(x) = -\frac{\cos x}{x}, \quad j_1(x) = \frac{\sin x}{x^2} - \frac{\cos x}{x}, \quad n_1(x) = -\frac{\cos x}{x^2} - \frac{\sin x}{x}$$

$$j_2(x) = \left(\frac{3}{x^3} - \frac{1}{x}\right) \sin x - \frac{3}{x^2} \cos x, \quad n_2(x) = -\left(\frac{3}{x^3} - \frac{1}{x}\right) \cos x - \frac{3}{x^2} \sin x,$$

$$\frac{d}{dt} P_{i \rightarrow n}(t) = \frac{2\pi}{\hbar} |V_{ni}|^2 \delta(E_n - E_i),$$

$$\frac{d\sigma}{d\Omega} = \frac{m^2}{4\pi^2 \hbar^4} \left| \int d^3r \mathcal{V}(r) e^{i(\vec{k}_f - \vec{k}_i) \cdot \vec{r}} \right|^2,$$

$$\sigma = \frac{(2S_R + 1) 4\pi (\hbar\Gamma_R/2)^2}{(2S_1 + 1)(2S_2 + 1) k^2 (\epsilon_k - \epsilon_r)^2 + (\hbar\Gamma_R/2)^2},$$

$$\frac{d\sigma}{d\Omega} = \left( \frac{d\sigma}{d\Omega} \right)_{\text{single}} \tilde{S}(\vec{q}), \quad \tilde{S}(\vec{q}) = \left| \sum_{\vec{a}} e^{i\vec{q} \cdot \vec{a}} \right|^2,$$

$$e^{i\vec{k} \cdot \vec{r}} = \sum_{\ell} (2\ell + 1) i^{\ell} j_{\ell}(kr) P_{\ell}(\cos \theta),$$

$$P_{\ell}(\cos \theta) = \sqrt{\frac{4\pi}{2\ell + 1}} Y_{\ell, m=0}(\theta, \phi),$$

$$P_0(x) = 1, \quad P_1(x) = x, \quad P_2(x) = (3x^2 - 1)/3,$$

$$f(\Omega) \equiv \sum_{\ell} (2\ell + 1) e^{i\delta_{\ell}} \sin \delta_{\ell} \frac{1}{k} P_{\ell}(\cos \theta)$$

$$\psi_{\vec{k}}(\vec{r})|_{R \rightarrow \infty} = e^{i\vec{k} \cdot \vec{r}} + \frac{e^{ikr}}{r} f(\Omega),$$

$$\frac{d\sigma}{d\Omega} = |f(\Omega)|^2, \quad \sigma = \frac{4\pi}{k^2} \sum_{\ell} (2\ell + 1) \sin^2 \delta_{\ell}, \quad \delta \approx -ak$$

$$L_{\pm} |\ell, m\rangle = \sqrt{\ell(\ell + 1) - m(m \pm 1)} |\ell, m \pm 1\rangle,$$

$$C_{m_{\ell}, m_s; JM}^{\ell, s} = \langle \ell, s, J, M | \ell, s, m_{\ell}, m_s \rangle,$$

$$\langle \tilde{\beta}, J, M | T_q^k | \beta, \ell, m_{\ell} \rangle = C_{qm_{\ell}; JM}^{k\ell} \frac{\langle \tilde{\beta}, J || T^{(k)} || \beta, \ell, J \rangle}{\sqrt{2J + 1}},$$

$$n = \frac{(2s + 1)}{(2\pi)^d} \int_{k < k_f} d^d k, \quad d \text{ dimensions,}$$

$$\{\Psi_s(\vec{x}), \Psi_{s'}^{\dagger}(\vec{y})\} = \delta^3(\vec{x} - \vec{y}) \delta_{ss'},$$

$$\Psi_s^{\dagger}(\vec{r}) = \frac{1}{\sqrt{V}} \sum_{\vec{k}} e^{i\vec{k} \cdot \vec{r}} a_s^{\dagger}(\vec{k}), \quad \{\Psi_s(\vec{x}), a_{\alpha}^{\dagger}\} = \phi_{\alpha, s}(\vec{x}).$$

1. Type  $\alpha$ ,  $\beta$  and  $\gamma$  particles exist in a TWO-DIMENSIONAL world. The  $\alpha$  particle has mass  $M_\alpha$  and is described by the two-dimensional field operator within a large area  $A$ ,

$$\begin{aligned}\Phi_\alpha(\vec{r}, t) &= \frac{1}{\sqrt{A}} \sum_{\vec{k}} e^{-iE_k t/\hbar + i\vec{k}\cdot\vec{r}} a_{\vec{k}}, \\ \Phi_\alpha^\dagger(\vec{r}, t) &= \frac{1}{\sqrt{A}} \sum_{\vec{k}} e^{iE_k t/\hbar - i\vec{k}\cdot\vec{r}} a_{\vec{k}}^\dagger, \\ E_k &= M_\alpha c^2 + \frac{\hbar^2 k^2}{2M_\alpha}.\end{aligned}$$

The  $\beta$  and  $\gamma$  particles are massless and described by the operators,

$$\begin{aligned}\Psi_\beta(\vec{r}, t) &= \frac{1}{\sqrt{A}} \sum_{\vec{q}} e^{-iE_q t/\hbar + i\vec{q}\cdot\vec{r}} b_{\vec{q}}, \\ \Psi_\beta^\dagger(\vec{r}, t) &= \frac{1}{\sqrt{A}} \sum_{\vec{q}} e^{iE_q t/\hbar - i\vec{q}\cdot\vec{r}} b_{\vec{q}}^\dagger, \\ \Psi_\gamma(\vec{r}, t) &= \frac{1}{\sqrt{A}} \sum_{\vec{q}} e^{-iE_q t/\hbar + i\vec{q}\cdot\vec{r}} c_{\vec{q}}, \\ \Psi_\gamma^\dagger(\vec{r}, t) &= \frac{1}{\sqrt{A}} \sum_{\vec{q}} e^{iE_q t/\hbar - i\vec{q}\cdot\vec{r}} c_{\vec{q}}^\dagger, \\ E_q &= \hbar c q.\end{aligned}$$

The massive  $\alpha$  particle can decay to an  $\alpha$  and a  $\beta$  particle via the interaction

$$H_{\text{int}} = g \int dx dy \left[ \Phi_\alpha(\vec{r}, t) \Psi_\beta^\dagger(\vec{r}, t) \Psi_\gamma^\dagger(\vec{r}, t) + \Phi_\alpha^\dagger(\vec{r}, t) \Psi_\beta(\vec{r}, t) \Psi_\gamma(\vec{r}, t) \right],$$

where the coupling constant  $g$  is small. The creation and destruction operators obey the commutation rules  $[a_{\vec{k}}, a_{\vec{k}'}] = \delta_{\vec{k}\vec{k}'}$ ,  $[b_{\vec{q}}, b_{\vec{q}'}] = \delta_{\vec{q}\vec{q}'}$  and  $[c_{\vec{q}}, c_{\vec{q}'}] = \delta_{\vec{q}\vec{q}'}$ . FYI: In two dimensions  $\delta_{\vec{k}\vec{k}'} = \delta_{k_x k'_x} \delta_{k_y k'_y}$ .

- (5 pts) Evaluate the commutator  $[\Phi_\alpha(\vec{r}, t), \Phi_\alpha^\dagger(\vec{r}', t)]$ .
- (5 pts) What is the dimensionality of  $g$ ?
- (20 pts) Calculate the rate at which an  $\alpha$  particle at rest decays into a  $\beta$  and a  $\gamma$  particle.
- (5 pts) How would your answer change if the  $\beta$  and  $\gamma$  particles were identical?

(Extra work space for #1)

$$\begin{aligned}
 \text{a) } [\Phi(\vec{r}, t), \Phi^\dagger(\vec{r}', t)] &= \frac{1}{A} \sum_{\vec{k}, \vec{k}'} e^{-iE_{\vec{k}}t/\hbar + i\vec{k}\cdot\vec{r}} e^{iE_{\vec{k}'}t/\hbar - i\vec{k}'\cdot\vec{r}'} [a_{\vec{k}}, a_{\vec{k}'}^\dagger] \\
 &= \frac{1}{A} \sum_{\vec{k}} e^{i\vec{k}\cdot(\vec{r}-\vec{r}')} = \frac{A}{A(2\pi)^2} \int dk_x dk_y e^{ik_x(x-x')} e^{-ik_y(y-y')} \\
 &= \delta(x-x') \delta(y-y') = \delta(\vec{r}-\vec{r}')
 \end{aligned}$$

$$\begin{aligned}
 \text{b) } [E] &= [g] \cdot [L^2] \cdot [\Phi \Phi^\dagger] \\
 \text{dimension of } \Phi \text{ or } \Psi &= \left[ \frac{1}{L} \right] \\
 \text{from commutation relation}
 \end{aligned}$$

$$[E] = [g] \left[ \frac{1}{L} \right]$$

$$[g] = \text{energy} \cdot \text{length}$$

$$\text{c) } \Gamma = \sum_{q, q'} \frac{2\pi}{\hbar} |m_{qq'}|^2 \delta(E_k - E_q - E_{q'})$$

$$\begin{aligned}
 m_{qq'} &= \langle 0 | a_q a_{q'} H_{int} a_{k=0}^\dagger | 0 \rangle \\
 &= \frac{g}{A^{3/2}} \int d^2r e^{-i\vec{q}\cdot\vec{r} - i\vec{q}'\cdot\vec{r}}
 \end{aligned}$$

$$= \frac{g}{A^{1/2}} \delta_{q, -q'}$$

$$|m_{qq'}|^2 = \frac{g^2}{A} \delta_{q, -q'}$$

(Extra work space for #1)

$$\begin{aligned}
 \Gamma &= \frac{2\pi}{\hbar} \frac{g^2}{A} \sum_q \delta(M_a c^2 - 2\hbar c q) \\
 &= \frac{2\pi}{\hbar} \frac{g^2}{A} \cdot \frac{A}{(2\pi)^2} \int 2\pi q dq \delta(M_a c^2 - 2\hbar c q) \\
 &= \frac{g^2}{\hbar} \int q dq \delta(M_a c^2 - 2\hbar c q) \\
 &= \frac{g^2}{2\hbar^2 c} q, \quad q = \frac{M_a c^2}{2\hbar c} = \frac{M_a c}{2\hbar} \\
 &= \frac{g^2}{2\hbar^2 c} \frac{M_a c}{2\hbar} = \frac{g^2 M_a}{4\hbar^3}
 \end{aligned}$$

check units

$$\left[ \frac{1}{t} \right] \stackrel{?}{=} [E^2] [L^2] \frac{[M]}{[E^3] [t^3]} = \frac{1}{[t]} \frac{1}{[E]} \left[ \frac{ML^2}{t^2} \right]$$

$\uparrow$   
 $= [E]$

$$\stackrel{!}{=} \frac{1}{t} \quad \checkmark$$

d) Matrix element would double  
 so  $|M|^2$  increases by 4.  
 But  $\rightarrow$  only sum over  $\frac{1}{2}$  as  
 many  $q$  values ( $q, -q$ ) are same  
 state. So overall

$\Gamma$  would double

2. Imagine you had calculated the following matrix element,

$$\mathcal{M} = \langle \alpha, J_f = 1, M_f = 0 | z^2 | \beta, J_i = 3, M_i = 0 \rangle.$$

For the following matrix elements, first state whether they are zero, and if not, express them in terms of  $\mathcal{M}$ . Evaluate any Clebsch-Gordan coefficients in your answers. You can use tables or online resources to get the values. (5 pts. each)

- $\langle \alpha, J_f = 1, M_f = 0 | x^2 + y^2 | \beta, J_i = 3, M_i = 0 \rangle$

I know  $\langle J=1 | x^2 + y^2 + z^2 | J=3 \rangle = 0$   
 so  
 $\langle \alpha | J_f=1, M_f=0 | x^2 + y^2 | \beta, J_i=3, M_i=0 \rangle = \underline{-\mathcal{M}}$

- $\langle \alpha, J_f = 1, M_f = 0 | z^2 | \beta, J_i = 3, M_i = 2 \rangle$

$z^2 \sim T_{q=0}^2$  so  
 $= 0$

- $\langle \alpha, J_f = 1, M_f = 0 | xz | \beta, J_i = 3, M_i = 0 \rangle$

$xz \sim T_1^2$  or  $T_{-1}^2$ , so  
 $= 0$

- $\langle \alpha, J_f = 1, M_f = 0 | xz | \beta, J_i = 3, M_i = 2 \rangle$

same reason

$= 0$

- $\langle \alpha, J_f = 1, M_f = 0 | xy | \beta, J_i = 3, M_i = 2 \rangle$

$\mathcal{M} = \langle M_f=0 | z^2 | M_i=0 \rangle = \langle 0 | \frac{2}{3} T_0^2 + \frac{1}{3} T_0^2 | 0 \rangle$   
 $\langle M_f=0 | xy | M_i=2 \rangle = \langle 0 | \frac{i}{\sqrt{6}} (T_2^2 - T_{-2}^2) | 0 \rangle$   
 $= \frac{i}{\sqrt{6}} \frac{\langle 1,0 | 2,-2, 3,2 \rangle}{\langle 1,0 | 2,0, 3,0 \rangle} \cdot \frac{3}{2} \mathcal{M}$   
 $= \frac{3i}{2\sqrt{6}} \mathcal{M} \frac{1/\sqrt{7}}{3/\sqrt{35}} = i \frac{\sqrt{5}}{2\sqrt{6}} \mathcal{M}$

(Extra work space for #2)

3. Electrons of mass  $m$  are placed in a three-dimensional spherically symmetric harmonic oscillator with potential,  $V = m\omega^2 r^2/2$ .

(a) (5 pts) What is the ground state energy for a single electron?

$$\frac{3}{2} \hbar \omega$$

(b) (5 pts) What is the net ground state energy if 20 electrons are in the well?

$N = 0 \rightarrow 2 \text{ electrons}$	$E = 3/2 \hbar \omega$	$3 \hbar \omega$
$N = 1 \ (l=1) \rightarrow 6 \text{ electrons}$	$E = 5/2 \hbar \omega$	$15 \hbar \omega$
$N = 2 \ (l=2) \rightarrow 10 \text{ electrons}$	$E = 7/2 \hbar \omega$	$35 \hbar \omega$
$N = 2 \ (l=0) \rightarrow 2 \text{ electrons}$	$E = 7/2 \hbar \omega$	$7 \hbar \omega$
		<b><math>60 \hbar \omega</math></b>

Now add an interaction with an external magnetic field,

$$V_B = -\mu \vec{B} \cdot (\vec{L} + 2\vec{S}),$$

where the strength of the magnetic field is adjusted so that  $\mu B = \omega$ .

(c) (10 pts) What is the new single-particle ground state energy?

$E = \frac{1}{2} \hbar \omega$  for  $N=0, m_s = 1/2, l=0$   
 or any other  $N, l=N, m_s = 1/2, m_l = l$

$$E = (N + \frac{3}{2}) \frac{\hbar \omega}{2} - \mu_B \hbar (1 + N) = \frac{1}{2} \hbar \omega$$

(d) (10 pts) What is the net ground state energy if 20 electrons are in the well?

Each level if  $E = \frac{1}{2} \hbar \omega$  can hold one electron, so you can put all 20 electrons in different levels with that energy.

$$E_{\text{net}} = 20 \cdot \frac{1}{2} \hbar \omega = 10 \hbar \omega$$



(Extra work space for #3)