## PHY 851 – QUANTUM MECHANICS MIDTERM II October 23-25, 2020

- Submit your exam via GRADESCOPE. Be sure that no page contains responses for more than one problem. Having different parts 2b,2c··· of a given problem on a single page is fine. Individual problems can span multiple pages.
- Once the exam is opened, you are to upload your answer within six hours.
- Do not write your name on the exam
- This exam is open-book, open-notes, open-internet but closed-mouth. You are permitted to use mathematical software, e.g. Mathematica. You are not to communicate with any other individuals regarding the exam, during the exam.

$$\begin{split} &\int_{-\infty}^{\infty} dx \; e^{-x^2/(2a^2)} = a\sqrt{2\pi}, \\ &H = i\hbar\partial_t, \; \vec{P} = -i\hbar\nabla, \\ &\sigma_z = \left( \begin{array}{cc} 1 & 0 \\ 0 & -1 \end{array} \right) \;, \sigma_x = \left( \begin{array}{cc} 0 & 1 \\ 1 & 0 \end{array} \right) \;, \; \sigma_y = \left( \begin{array}{cc} 0 & -i \\ i & 0 \end{array} \right), \\ &U(t, -\infty) = 1 + \frac{-i}{\hbar} \int_{-\infty}^{t} dt' \; V(t') U(t', -\infty), \\ &\langle x | x' \rangle = \delta(x - x'), \; \langle p | p' \rangle = \frac{1}{2\pi\hbar} \delta(p - p'), \\ &|p\rangle = \int dx \; |x\rangle e^{ipx/\hbar}, \; |x\rangle = \int \frac{dp}{2\pi\hbar} |p\rangle e^{-ipx/\hbar}, \\ &H = \frac{P^2}{2m} + \frac{1}{2}m\omega^2 x^2 = \hbar\omega(a^{\dagger}a + 1/2), \\ &a^{\dagger} = \sqrt{\frac{m\omega}{2\hbar}} X - i\sqrt{\frac{1}{2\hbarm\omega}} P, \\ &\psi_0(x) = \frac{1}{(\pi b^2)^{1/4}} e^{-x^2/2b^2}, \; b^2 = \frac{\hbar}{m\omega}, \\ &\rho(\vec{r}, t) = \psi^*(\vec{r}_1, t_1)\psi(\vec{r}_2, t_2) \\ &\vec{j}(\vec{r}, t) = \frac{-i\hbar}{2m}(\psi^*(\vec{r}, t)\nabla\psi(\vec{r}, t) - (\nabla\psi^*(\vec{r}, t))\psi(\vec{r}, t)) \\ &- \frac{e\vec{A}}{mc} |\psi(\vec{r}, t)|^2. \\ &H = \frac{(\vec{P} - e\vec{A}/c)^2}{2m} + e\Phi, \\ \\ \text{For } V = \beta\delta(x - y): \; - \frac{\hbar^2}{2m} \left( \frac{\partial}{\partial x}\psi(x)|_{y+\epsilon} - \frac{\partial}{\partial x}\psi(x)|_{y-\epsilon} \right) = -\beta\psi(y), \\ &\vec{E} = -\nabla\Phi - \frac{1}{c}\partial_t\vec{A}, \; \vec{B} = \nabla \times \vec{A}, \\ &\omega_{\text{cyclotron}} = \frac{eB}{mc}, \\ &e^{A+B} = e^A e^B e^{-C/2}, \; \text{if } [A, B] = C, \; \text{and } [C, A] = [C, B] = 0, \\ &Y_{0,0} = \frac{1}{\sqrt{4\pi}}, \; Y_{1,0} = \sqrt{\frac{3}{4\pi}}\cos\theta, \; Y_{1,\pm 1} = \mp \sqrt{\frac{3}{8\pi}}\sin\theta e^{i\pm\phi}, \\ &Y_{2,\pm 2} = \sqrt{\frac{15}{16\pi}}(3cs^2\theta - 1), \; Y_{2,\pm 1} = \mp \sqrt{\frac{15}{8\pi}}\sin\theta \cos\theta e^{\pm i\phi}, \end{aligned}$$

$$\begin{split} |N\rangle &= |n\rangle - \sum_{m \neq n} |m\rangle \frac{1}{\epsilon_m - \epsilon_n} \langle m|V|n\rangle + \cdots \\ & E_N = \epsilon_n + \langle n|V|n\rangle - \sum_{m \neq n} \frac{|\langle m|V|n\rangle|^2}{\epsilon_m - \epsilon_n} \\ j_0(x) &= \frac{\sin x}{x}, \ n_0(x) = -\frac{\cos x}{x}, \ j_1(x) = \frac{\sin x}{x^2} - \frac{\cos x}{x}, \ n_1(x) = -\frac{\cos x}{x^2} - \frac{\sin x}{x} \\ j_2(x) &= \left(\frac{3}{x^3} - \frac{1}{x}\right) \sin x - \frac{3}{x^2} \cos x, \ n_2(x) = -\left(\frac{3}{x^3} - \frac{1}{x}\right) \cos x - \frac{3}{x^2} \sin x, \\ \frac{d}{dt} P_{i \to n}(t) = \frac{2\pi}{h} |V_{ni}|^2 \delta(E_n - E_i), \\ \frac{d\sigma}{dt2} &= \frac{m^2}{4\pi^2 l_i t^i} \left| \int d^3 r V(r) e^{i(\vec{k}_j - \vec{k}_i) \cdot r} \right|^2, \\ \sigma &= \frac{(2S_R + 1)}{(2S_1 + 1)(2S_2 + 1)} \frac{4\pi}{k^2} \frac{(h\Gamma_R/2)^2}{(\epsilon_k - \epsilon_r)^2 + (h\Gamma_R/2)^2}, \\ \frac{d\sigma}{d\Omega} &= \left(\frac{d\sigma}{d\Omega}\right)_{\text{single}} \tilde{S}(\vec{q}), \ \tilde{S}(\vec{q}) &= \left|\sum_{\vec{a}} e^{i\vec{q} \cdot \vec{a}}\right|^2, \\ e^{i\vec{k} \cdot \vec{r}} &= \sum (2\ell + 1)i^t j_\ell(kr) P_\ell(\cos \theta), \\ P_\ell(\cos \theta) &= \sqrt{\frac{4\pi}{2\ell + 1}} Y_{\ell,m=0}(\theta, \phi), \\ P_0(x) &= 1, \ P_1(x) = x, \ P_2(x) &= (3x^2 - 1)/3, \\ f(\Omega) &= \sum_{\ell} (2\ell + 1)e^{i\delta r} \sin \delta_t \frac{1}{k} P_\ell(\cos \theta) \\ \psi_{\vec{k}}(\vec{r})|_{R \to \infty} &= e^{i\vec{k} \cdot \vec{r}} + \frac{e^{ikr}}{r} f(\Omega), \\ \frac{d\sigma}{d\Omega} &= |f(\Omega)|^2, \ \sigma &= \frac{4\pi}{k^2} \sum_{\ell} (2\ell + 1) \sin^2 \delta_\ell, \ \delta \approx -ak \\ L_{\pm}|\ell,m\rangle &= \sqrt{\ell(\ell+1) - m(m\pm 1)}|\ell,m\pm 1\rangle, \\ C_{m,m,n,M}^{\ell m} &= (\ell, s, J, M|\ell, s, m, m_s), \\ (\vec{\beta}, J, M|T_q^k|\beta, \ell, m_\ell) &= C_{mn,M}^{k\ell} (\vec{\sigma}) \int_{k < k_T} d^k k, \ d \text{ dimensions}, \\ \{\Psi_s(\vec{x}), \Psi_{s'}^{\dagger}(\vec{y})\} &= \delta^3(\vec{x} - \vec{y}) \delta_{s'}, \\ \Psi_s(\vec{r}) = \frac{1}{\sqrt{V}} \sum_{\vec{k}} e^{i\vec{k} \cdot \vec{r}} a_s^{\dagger}(\vec{k}), \ \{\Psi_s(\vec{x}), a_{\alpha}^{\dagger}\} = \phi_{\alpha,s}(\vec{x}). \end{split}$$

- 1. A particle of mass m and charge e feels a strong constant electric field  $\vec{E} = E_0 \hat{x}$ . Additionally, the particle experiences a constant weak magnetic field  $\vec{B} = B_0 \hat{z}$ . The magnetic field strength is less than the electric field strength,  $E_0 > B_0$ .
  - (a) (10 pts) Describe a reference frame where either  $\vec{E}$  or  $\vec{B}$  is zero. State the velocity (both magnitude and direction) of that frame, and describe the fields in that frame.
  - (b) Averaged over a sufficiently long time interval  $\Delta$  to avoid any oscillatory terms, answer these questions about the components of the average velocity  $(\bar{v}_x(t), \bar{v}_y(t), \bar{v}_z(t))$  at large times,  $t >> \Delta t$ . Assume the particle might has some original momentum with non-zero components in each direction. Circle the correct answers:
    - i. (5 pts) Which components average to zero?  $[v_x, v_y, v_z]$
    - ii. (5 pts) Which components average to a constant? (but not necessarily zero)  $[v_x, v_y, v_z]$
    - iii. (5 pts) Which components continue to accelerate over time?  $[v_x, v_y, v_z]$

You might choose a gauge where the vector potential  $\vec{A}$  is pointed along the  $\phi$  axis.



2. Consider a particle of mass  $m_0$  confined to a two-dimensional harmonic oscillator,

$$V(x,y)=rac{1}{2}m\omega^2(r^2), \hspace{1em} r=\sqrt{x^2+y^2}.$$

Consider solutions of the form,

$$\psi_{nm}(x,y)=\phi_{nm}(r)e^{im\phi}, \ \ m= \ {
m an \ integer}.$$

- (a) (10 pts) Write a one-dimensional differential equation (Schrödinger's equation) in terms of r for  $\phi_{nm}(r)$  assuming  $\phi_{nm}$  is an eigenstate with energy  $E_{nm}$ . Here, n denotes different eigenstates with the same m.
- (b) (10 pts) What are the eigen-energies? Also state the degeneracies of each level. (easier to work out in Cartesian basis)
- (c) (5 pts) Consider the state (un-normalized)

$$egin{aligned} \psi_{NM}&=(a^{\dagger}_x+ia^{\dagger}_y)^M(a^{\dagger}_xa^{\dagger}_x+a^{\dagger}_ya^{\dagger}_y)^{(N-M)/2}|0
angle, \quad M\geq 0\ \psi_{NM}&=(a^{\dagger}_x-ia^{\dagger}_y)^{|M|}(a^{\dagger}_xa^{\dagger}_x+a^{\dagger}_ya^{\dagger}_y)^{(N-M)/2}|0
angle, \quad M< 0 \end{aligned}$$

with energy  $(N+1)\hbar\omega$ . Here,  $a_x^{\dagger}$  and  $a_y^{\dagger}$  are the raising operators in the Cartesian basis. Show that under rotations that  $\psi_{NM}$  behaves as  $e^{iM\phi}$ . I.e., show that

$$R(\Delta\phi)\psi_{NM}(r,\phi)=e^{iM\Delta\phi}\psi_{NM}(r,\phi).$$

You may use the fact that  $a_x^{\dagger}$  and  $a_y^{\dagger}$  rotate amongst each other like x and y. (5 pts) For a given eigen energy  $(N \pm 1)\hbar c_y$  what values of M are possible?

(d) (5 pts) For a given eigen-energy 
$$(N + 1)\hbar\omega$$
, what values of  $M$  are possible?  

$$E_{2m} \left(\partial_{r}^{3} + \frac{1}{r} + \frac{1}{r}\partial_{r}^{3}\right) + v(r) = F \mathcal{F}$$

$$E_{2m} \left(\partial_{r}^{3} + \frac{1}{r}\partial_{r}\right) + \frac{1}{2}m\omega'r' + \frac{1}{2mr'} = F \mathcal{F}$$

$$E_{m} = \left(N_{+} + N_{2} + 1\right) \hbar \omega = \left(M + 1\right) \hbar \omega = \left(M + 1\right) \hbar \omega, N = 0, 123$$

$$degen \left(M\right) = N + 1$$

$$e^{(N + 1)} \left(a_{+}^{+} \cos \theta + a_{5} \sin \theta\right) + i\left(a_{5}^{+} \cos \theta - a_{x}^{+} \sin \theta\right)$$

$$= \left(a_{+}^{+} + i a_{5}^{-1}\right) - \left(a_{x}^{+} \cos \theta + a_{5} \sin \theta\right) + i\left(a_{5}^{+} \cos \theta - a_{x}^{+} \sin \theta\right)$$

$$= \left(a_{+}^{+} + i a_{5}^{-1}\right) - \left(a_{x}^{+} \cos \theta + a_{5} \sin \theta\right) + i\left(a_{5}^{+} \cos \theta - a_{x}^{+} \sin \theta\right)$$

$$= \left(a_{+}^{+} + i a_{5}^{-1}\right) - \left(a_{+}^{+} \cos \theta + a_{5} \sin \theta\right) + i\left(a_{5}^{+} \cos \theta - a_{x}^{+} \sin \theta\right)$$

$$= \left(a_{+}^{+} + i a_{5}^{-1}\right) - \left(a_{+}^{+} \cos \theta + a_{5} \sin \theta + i \left(a_{5}^{+} \cos \theta - a_{x}^{+} \sin \theta\right)\right)$$

$$= \left(a_{+}^{+} + i a_{5}^{-1}\right) - \left(a_{+}^{+} \cos \theta - a_{+}^{-1} \sin \theta\right)$$

$$M = -N, -N + 2, -N + 4 - N - 3, N$$

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$$M = N + i n + 2, -N + 4 - N - 3, N$$

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3. (20 pts) Two distinguishable spinless particles of mass m occupy the same  $\ell = 1$  eigenstate of a spherically symmetric potential. An spin-spin interaction is then applied,

$$V_{
m ss} = -eta ec{L}_1 \cdot ec{L}_2.$$

- (a) (5 pts) What values of the total angular momentum might the system attain?
- (b) (15 pts) What are the shifts of the energy due to  $V_{ss}$ ? Assume you needn't worry about mixing with other levels.



4. (20 pts) Two spinless particles ("1" and "2") of mass m occupy an  $\ell = 1$  eigenstate and an  $\ell = 2$  eigenstate of a spherically symmetric potential. A magnetic field with the following interaction is then applied,

$$V_{
m b}=-\mu_1ec{B}\cdotec{L}_1-\mu_2ec{B}\cdotec{L}_2.$$

- (a) (5 pts) What values of the total angular momentum might the system attain?
- (b) (15 pts) What are the shifts of the energy due to  $V_b$ ? Assume you needn't worry about mixing with other levels.

a) J = 1, 2, 3b)  $\Delta E = -Bt \{ \mathcal{L}, \mathcal{M}, \mathcal{H}, \mathcal{H}, \mathcal{M}, \mathcal{H} \}$   $ranger (row - \mathcal{M}, Bt - 2\mathcal{M}, Bt$  $tw + \mathcal{M}, Bt + 2\mathcal{M}, Bt$