

*PHY 851 – QUANTUM MECHANICS*  
*MIDTERM I*

September 25-28, 2020

- Submit your exam via GRADESCOPE. Be sure that no page contains responses for more than one problem. Having different parts 2b,2c... of a given problem on a single page is fine. Individual problems can span multiple pages.
- Once the exam is opened, you are to upload your answer within six hours.
- This exam is open-book, open-notes, open-internet but closed-mouth. You are permitted to use mathematical software, e.g. Mathematica. You are not to communicate with any other individuals regarding the exam, during the exam.

$$\int_{-\infty}^{\infty} dx e^{-x^2/(2a^2)} = a\sqrt{2\pi},$$

$$H = i\hbar\partial_t, \quad \vec{P} = -i\hbar\nabla,$$

$$\sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \quad \sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix},$$

$$U(t, -\infty) = 1 + \frac{-i}{\hbar} \int_{-\infty}^t dt' V(t')U(t', -\infty),$$

$$\langle x|x'\rangle = \delta(x - x'), \quad \langle p|p'\rangle = \frac{1}{2\pi\hbar} \delta(p - p'),$$

$$|p\rangle = \int dx |x\rangle e^{ipx/\hbar}, \quad |x\rangle = \int \frac{dp}{2\pi\hbar} |p\rangle e^{-ipx/\hbar},$$

$$H = \frac{P^2}{2m} + \frac{1}{2}m\omega^2 x^2 = \hbar\omega(a^\dagger a + 1/2),$$

$$a^\dagger = \sqrt{\frac{m\omega}{2\hbar}} X - i\sqrt{\frac{1}{2\hbar m\omega}} P,$$

$$\psi_0(x) = \frac{1}{(\pi b^2)^{1/4}} e^{-x^2/2b^2}, \quad b^2 = \frac{\hbar}{m\omega},$$

$$\rho(\vec{r}, t) = \psi^*(\vec{r}_1, t_1)\psi(\vec{r}_2, t_2)$$

$$\vec{j}(\vec{r}, t) = \frac{-i\hbar}{2m} (\psi^*(\vec{r}, t)\nabla\psi(\vec{r}, t) - (\nabla\psi^*(\vec{r}, t))\psi(\vec{r}, t))$$

$$- \frac{e\vec{A}}{mc} |\psi(\vec{r}, t)|^2.$$

$$H = \frac{(\vec{P} - e\vec{A}/c)^2}{2m} + e\Phi,$$

$$\text{For } V = \beta\delta(x - y): \quad -\frac{\hbar^2}{2m} \left( \frac{\partial}{\partial x} \psi(x)|_{y+\epsilon} - \frac{\partial}{\partial x} \psi(x)|_{y-\epsilon} \right) = -\beta\psi(y),$$

$$\vec{E} = -\nabla\Phi - \frac{1}{c} \partial_t \vec{A}, \quad \vec{B} = \nabla \times \vec{A},$$

$$\omega_{\text{cyclotron}} = \frac{eB}{mc},$$

$$e^{A+B} = e^A e^B e^{-C/2}, \quad \text{if } [A, B] = C, \text{ and } [C, A] = [C, B] = 0,$$

$$Y_{0,0} = \frac{1}{\sqrt{4\pi}}, \quad Y_{1,0} = \sqrt{\frac{3}{4\pi}} \cos\theta, \quad Y_{1,\pm 1} = \mp \sqrt{\frac{3}{8\pi}} \sin\theta e^{\pm i\phi},$$

$$Y_{2,0} = \sqrt{\frac{5}{16\pi}} (3\cos^2\theta - 1), \quad Y_{2,\pm 1} = \mp \sqrt{\frac{15}{8\pi}} \sin\theta \cos\theta e^{\pm i\phi},$$

$$Y_{2,\pm 2} = \sqrt{\frac{15}{32\pi}} \sin^2\theta e^{\pm 2i\phi}, \quad Y_{\ell-m}(\theta, \phi) = (-1)^m Y_{\ell m}^*(\theta, \phi).$$

$$|N\rangle = |n\rangle - \sum_{m \neq n} |m\rangle \frac{1}{\epsilon_m - \epsilon_n} \langle m|V|n\rangle + \dots$$

$$E_N = \epsilon_n + \langle n|V|n\rangle - \sum_{m \neq n} \frac{|\langle m|V|n\rangle|^2}{\epsilon_m - \epsilon_n}$$

$$j_0(x) = \frac{\sin x}{x}, \quad n_0(x) = -\frac{\cos x}{x}, \quad j_1(x) = \frac{\sin x}{x^2} - \frac{\cos x}{x}, \quad n_1(x) = -\frac{\cos x}{x^2} - \frac{\sin x}{x}$$

$$j_2(x) = \left(\frac{3}{x^3} - \frac{1}{x}\right) \sin x - \frac{3}{x^2} \cos x, \quad n_2(x) = -\left(\frac{3}{x^3} - \frac{1}{x}\right) \cos x - \frac{3}{x^2} \sin x,$$

$$\frac{d}{dt} P_{i \rightarrow n}(t) = \frac{2\pi}{\hbar} |V_{ni}|^2 \delta(E_n - E_i),$$

$$\frac{d\sigma}{d\Omega} = \frac{m^2}{4\pi^2 \hbar^4} \left| \int d^3r \mathcal{V}(r) e^{i(\vec{k}_f - \vec{k}_i) \cdot \vec{r}} \right|^2,$$

$$\sigma = \frac{(2S_R + 1) 4\pi (\hbar\Gamma_R/2)^2}{(2S_1 + 1)(2S_2 + 1) k^2 (\epsilon_k - \epsilon_r)^2 + (\hbar\Gamma_R/2)^2},$$

$$\frac{d\sigma}{d\Omega} = \left( \frac{d\sigma}{d\Omega} \right)_{\text{single}} \tilde{S}(\vec{q}), \quad \tilde{S}(\vec{q}) = \left| \sum_{\vec{a}} e^{i\vec{q} \cdot \vec{a}} \right|^2,$$

$$e^{i\vec{k} \cdot \vec{r}} = \sum_{\ell} (2\ell + 1) i^{\ell} j_{\ell}(kr) P_{\ell}(\cos \theta),$$

$$P_{\ell}(\cos \theta) = \sqrt{\frac{4\pi}{2\ell + 1}} Y_{\ell, m=0}(\theta, \phi),$$

$$P_0(x) = 1, \quad P_1(x) = x, \quad P_2(x) = (3x^2 - 1)/3,$$

$$f(\Omega) \equiv \sum_{\ell} (2\ell + 1) e^{i\delta_{\ell}} \sin \delta_{\ell} \frac{1}{k} P_{\ell}(\cos \theta)$$

$$\psi_{\vec{k}}(\vec{r})|_{R \rightarrow \infty} = e^{i\vec{k} \cdot \vec{r}} + \frac{e^{ikr}}{r} f(\Omega),$$

$$\frac{d\sigma}{d\Omega} = |f(\Omega)|^2, \quad \sigma = \frac{4\pi}{k^2} \sum_{\ell} (2\ell + 1) \sin^2 \delta_{\ell}, \quad \delta \approx -ak$$

$$L_{\pm} |\ell, m\rangle = \sqrt{\ell(\ell + 1) - m(m \pm 1)} |\ell, m \pm 1\rangle,$$

$$C_{m_{\ell}, m_s; JM}^{\ell, s} = \langle \ell, s, J, M | \ell, s, m_{\ell}, m_s \rangle,$$

$$\langle \tilde{\beta}, J, M | T_q^k | \beta, \ell, m_{\ell} \rangle = C_{qm_{\ell}; JM}^{k\ell} \frac{\langle \tilde{\beta}, J || T^{(k)} || \beta, \ell, J \rangle}{\sqrt{2J + 1}},$$

$$n = \frac{(2s + 1)}{(2\pi)^d} \int_{k < k_f} d^d k, \quad d \text{ dimensions,}$$

$$\{\Psi_s(\vec{x}), \Psi_{s'}^{\dagger}(\vec{y})\} = \delta^3(\vec{x} - \vec{y}) \delta_{ss'},$$

$$\Psi_s^{\dagger}(\vec{r}) = \frac{1}{\sqrt{V}} \sum_{\vec{k}} e^{i\vec{k} \cdot \vec{r}} a_s^{\dagger}(\vec{k}), \quad \{\Psi_s(\vec{x}), a_{\alpha}^{\dagger}\} = \phi_{\alpha, s}(\vec{x}).$$

1. An electron is in the  $|\uparrow\rangle$  (up along the  $z$  axis) state.

- (a) (10 pts) What is the probability of finding it in the  $|\uparrow\rangle$  state if observed after rotating the coordinate system by 90 degrees about the  $z$  axis?
- (b) (10 pts) If the rotation were instead 180 degrees about the  $x$  axis, what would the probability of observing the  $|\uparrow\rangle$  state be?

a) 100%

b)  $|\psi\rangle = e^{i\sigma_x \phi/2} \begin{pmatrix} 1 \\ 0 \end{pmatrix}$

$$= \cos \frac{\phi}{2} + i \sin \frac{\phi}{2} \sigma_x \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

$$= \begin{pmatrix} \cos \frac{\phi}{2} \\ i \sin \frac{\phi}{2} \end{pmatrix}$$

$$P_{\text{prob}} = \cos^2 \frac{\phi}{2}, \quad \phi = 180^\circ$$

$$P_{\text{prob}} = 0$$

2. (20 pts) Consider a two-state system, with the Hamiltonian,

$$H = A + B\vec{\sigma} \cdot \hat{n},$$

where  $\hat{n}$  is a unit vector in the  $x - z$  plane pointing 30 degrees from the  $z$  axis.

At  $t = 0$ , the state

$$|\uparrow\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

is occupied. Find the probability the state  $|\downarrow\rangle$  is occupied as a function of time.

$$e^{-iHt} = e^{-iAt/\hbar} e^{-iB\vec{\sigma} \cdot \hat{n} t/\hbar}$$

$$= e^{-iAt/\hbar} \left\{ \cos(Bt/\hbar) - i(\vec{\sigma} \cdot \hat{n}) \sin(Bt/\hbar) \right\}$$

$$\vec{\sigma} \cdot \hat{n} = \sqrt{\frac{3}{2}} \sigma_z + \frac{1}{2} \sigma_x$$

$$\begin{pmatrix} 0 \\ 1 \end{pmatrix}^\dagger e^{-iHt} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = e^{-iAt/\hbar} \left( \frac{-i \sin Bt/\hbar}{2} \right)$$

$$P_{nb} = \frac{1}{4} \sin^2 Bt/\hbar$$

3. (20 pts) A particle of mass  $m$  is in the first excited state of a harmonic oscillator with characteristic frequency  $\omega$ . The well suddenly dissolves. What is the differential probability,  $dN/dp$ , for observing the particle with momentum  $p$ ?

$$P_p = |\langle p | \psi_1 \rangle|^2$$

$$= \left| \int dx e^{-ipx/\hbar} \psi_1(x) \right|^2$$

$$\psi_0(x) = \frac{1}{(\pi b^2)^{1/4}} e^{-x^2/2b^2}, \quad b^2 = \frac{\hbar}{m\omega}$$

$$\psi_1(x) = \left( \sqrt{\frac{m\omega}{2\hbar}} x - i \sqrt{\frac{1}{2m\hbar\omega}} (-i\hbar \partial_x) \right) \psi_0(x)$$

$$= \sqrt{2} \frac{x}{b} e^{-x^2/2b^2} \frac{1}{(\pi b^2)^{1/4}}$$

$$P_p = \frac{2}{\sqrt{\pi} b^3} \left| \frac{1}{\sqrt{L}} \int dx e^{-ipx/\hbar} x e^{-x^2/2b^2} \right|^2$$

$$= \frac{2}{\sqrt{\pi} b^3 L} \left| i\hbar \partial_p \int dx e^{-ipx/\hbar} e^{-x^2/2b^2} \right|^2$$

$$= \frac{2}{\sqrt{\pi} b^3 L} \left| i\hbar \partial_p \int dx e^{-\frac{(x + \frac{ipb^2}{\hbar})^2}{2b^2}} e^{-\frac{p^2 b^2}{2\hbar^2}} \right|^2$$

$$\rightarrow \frac{2\hbar^2}{\sqrt{\pi} b^3 L} \left| \partial_p \left( \sqrt{2\pi} b^2 e^{-\frac{p^2 b^2}{2\hbar^2}} \right) \right|^2$$

$$= \frac{2\hbar^2}{\sqrt{\pi} b^3 L} \left| \sqrt{2\pi} b \frac{p b^2}{\hbar^2} e^{-\frac{p^2 b^2}{2\hbar^2}} \right|^2 = \frac{4\sqrt{\pi} b^3}{L \hbar^2} p^2 e^{-\frac{p^2 b^2}{\hbar^2}}$$

$$P_p = \frac{4\sqrt{\pi} b^3}{L h^3} p^2 e^{-\frac{p^2 b^2}{h^2}} = P_{\text{prob}} \text{ to go into a single state}$$

$$\frac{dN}{dp} = P_p \cdot \frac{dN_{\text{states}}}{dp} = P_p \cdot \frac{L}{2\pi h}$$

$$= \frac{2 b^3 p^2}{\sqrt{\pi} h^3} e^{-p^2 b^2 / h^2}$$

where  
 $b = \sqrt{h/mv}$

$$\int p^2 dp e^{-p^2 b^2 / h^2} = \left(\frac{h}{\sqrt{2} b}\right)^3 (2\pi)^{1/2} = \frac{\sqrt{\pi} h^3}{2b^3}$$

4. (15 pts) Consider a particle of mass  $m$  confined to two-dimensions which feels the potential,

$$V(x, y) = \frac{1}{2}m\omega^2(x^2 + y^2).$$

What are the energy levels and their degeneracies?

$$E = (n_x + n_y + 1)\hbar\omega$$

$$= N\hbar\omega, \quad N = 1, 2, 3, \dots$$

$$d(N) = N$$

1 way for  $N=1$

2 ways for  $N=2$

3 " "  $N=3$

,

!



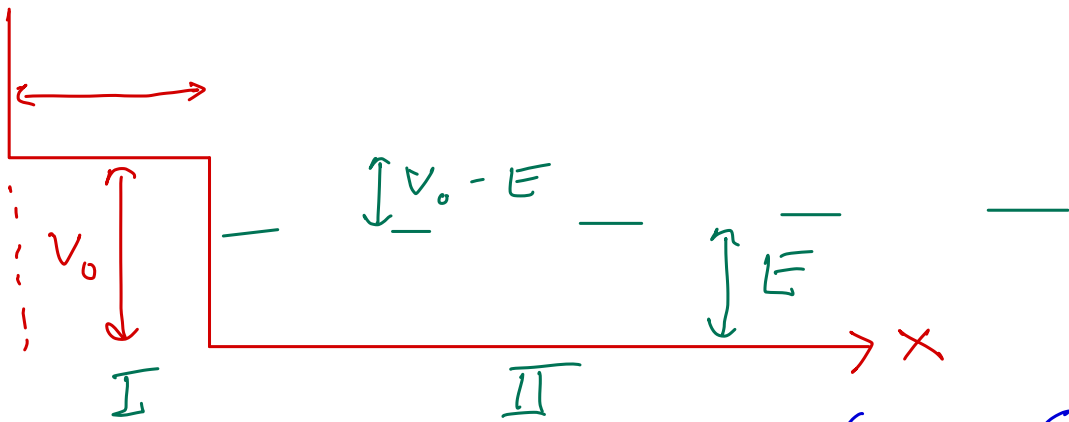
5. Consider a particle of mass  $m$  approaching a positive potential step from the  $x = +\infty$  direction with wave number  $k$ . The potential step is larger than the kinetic energy of the approaching particle, and has the form,

$$V(x) = \begin{cases} \infty, & x < 0 \\ V_0, & 0 < x < a \\ 0, & x > a \end{cases} .$$

For  $x > a$ , the wave function has the form

$$\psi(x) = e^{-ikx} - e^{2i\delta(k)} e^{ikx} .$$

- (a) (20 pts) Find the phase shift  $\delta(k)$ .  
 (b) (5 pts) What is  $\delta(k)$  in the limit  $V_0 \rightarrow \infty$ ?



(a)  $\psi_{\text{I}} = A \sinh q x$ ,  $\psi_{\text{II}} = \sin(kx + \delta)$   
 $q = \sqrt{2m(V_0 - E)/\hbar^2}$ ,  $k = \sqrt{2mE/\hbar^2}$   
 $A \sinh q a = \sin(ka + \delta)$   
 $q A \cosh q a = k \cos(ka + \delta)$   
 $\frac{1}{q} \tanh q a = \frac{1}{k} \tan(ka + \delta)$   
 $\delta = -ka + \tan^{-1} \frac{k}{q} \tanh q a$

(b) As  $V_0 \rightarrow \infty$ ,  $q \rightarrow \infty$   
 $\delta = -ka$  ( $\psi_{\text{I}} = \sin(kx - a)$ )