FINAL EXAM, PHYSICS 851, FALL 2021

Your Name: _____

Tuesday, December 14, 12:45 PM This exam is worth 100 points.

$$\begin{split} & \int_{-\infty}^{\infty} dx \; e^{-x^2/(2a^2)} = a\sqrt{2\pi}, \\ & H = i\hbar\partial_t, \; \vec{P} = -i\hbar\nabla, \\ \sigma_z &= \left(\begin{array}{cc} 1 & 0 \\ 0 & -1 \end{array} \right), \sigma_x = \left(\begin{array}{cc} 0 & 1 \\ 1 & 0 \end{array} \right), \; \sigma_y = \left(\begin{array}{cc} 0 & -i \\ i & 0 \end{array} \right), \\ & U(t, -\infty) = 1 + \frac{-i}{\hbar} \int_{-\infty}^{t} dt' \; V(t') U(t', -\infty), \\ & \langle x | x' \rangle = \delta(x - x'), \; \langle p | p' \rangle = \frac{1}{2\pi\hbar} \delta(p - p'), \\ & | p \rangle = \int dx \; | x \rangle e^{ipx/\hbar}, \; | x \rangle = \int \frac{dp}{2\pi\hbar} | p \rangle e^{-ipx/\hbar}, \\ & H = \frac{P^2}{2m} + \frac{1}{2} m\omega^2 x^2 = \hbar\omega (a^{\dagger}a + 1/2), \\ & a^{\dagger} = \sqrt{\frac{m\omega}{2\hbar}} X - i\sqrt{\frac{1}{2\hbarm\omega}} P, \\ & \psi_0(x) = \frac{1}{(\pi b^2)^{1/4}} e^{-x^2/2b^2}, \; b^2 = \frac{\hbar}{m\omega}, \\ & \rho(\vec{r}, t) = \psi^*(\vec{r}, t_1)\psi(\vec{r}_2, t_2) \end{split}$$
$$\vec{j}(\vec{r}, t) = \frac{-i\hbar}{2m} (\psi^*(\vec{r}, t)\nabla\psi(\vec{r}, t) - (\nabla\psi^*(\vec{r}, t))\psi(\vec{r}, t)) - \frac{c\vec{A}}{mc} |\psi(\vec{r}, t)|^2. \\ & H = \frac{(\vec{P} - e\vec{A}/c)^2}{2m} + e\Phi, \end{aligned}$$
For $V = \beta\delta(x - y) : - \frac{\hbar^2}{2m} (\partial_x\psi(x)|_{y+\epsilon} - \partial_x\psi(x)|_{y-\epsilon}) = -\beta\psi(y), \\ & \vec{E} = -\nabla\Phi - \frac{1}{c}\partial_t \vec{A}, \; \vec{B} = \nabla \times \vec{A}, \\ & \omega_{\text{cyclutron}} = \frac{eB}{mc}, \\ & e^{A+B} = e^A e^B e^{-C/2}, \; \text{if } [A, B] = C, \; \text{and } [C, A] = [C, B] = 0, \end{aligned}$
$$Y_{0,0} = \frac{1}{\sqrt{4\pi}}, \; Y_{1,0} = \sqrt{\frac{3}{4\pi}} \cos\theta, \; Y_{1,\pm 1} = \mp \sqrt{\frac{3}{8\pi}} \sin\theta e^{i\pm\phi}, \\ & Y_{2,0} = \sqrt{\frac{5}{16\pi}} (3\cos^2\theta - 1), \; Y_{2,\pm 1} = \mp \sqrt{\frac{15}{8\pi}} \sin\theta \cos\theta e^{\pm i\phi}, \\ & Y_{2,\pm 2} = \sqrt{\frac{15}{32\pi}} \sin^2 \theta e^{\pm 2i\phi}, \; Y_{\ell-m}(\theta, \phi) = (-1)^m Y_{\ell m}^*(\theta, \phi). \end{aligned}$$

$$\begin{split} |N\rangle &= |n\rangle - \sum_{m \neq n} |m\rangle \frac{1}{\epsilon_m - \epsilon_n} \langle m|V|n\rangle + \cdots \\ E_N &= \epsilon_n + \langle n|V|n\rangle - \sum_{m \neq n} \frac{|\langle m|V|n\rangle|^2}{\epsilon_m - \epsilon_n} \\ j_0(x) &= \frac{\sin x}{x}, \ n_0(x) = -\frac{\cos x}{x}, \ j_1(x) = \frac{\sin x}{x^2} - \frac{\cos x}{x}, \ n_1(x) = -\frac{\cos x}{x^2} - \frac{\sin x}{x} \\ j_2(x) &= \left(\frac{3}{x^3} - \frac{1}{x}\right) \sin x - \frac{3}{x^2} \cos x, \ n_2(x) = -\left(\frac{3}{x^3} - \frac{1}{x}\right) \cos x - \frac{3}{x^2} \sin x, \\ \frac{d}{dt} P_{1\rightarrow n}(t) &= \frac{2\pi}{h} |V_{ni}|^2 \delta(E_n - E_i), \\ \frac{d\sigma}{d\Omega} &= \frac{d\pi^2}{4\pi^2 h^4} \left| \int d^3 r \mathcal{V}(r) e^{i(\vec{k}_f - \vec{k}_i) \cdot \vec{r}} \right|^2, \\ \sigma &= \frac{(2S_R + 1)}{(2S_1 + 1)(2S_2 + 1)} \frac{4\pi}{k^2} \frac{(h\Gamma_R/2)^2}{(\epsilon_k - \epsilon_r)^2 + (h\Gamma_R/2)^2}, \\ \frac{d\sigma}{d\Omega} &= \left(\frac{d\sigma}{d\Omega}\right)_{\text{single}} \delta(\vec{q}), \ \vec{S}(\vec{q}) &= \frac{1}{N} \left| \sum_{\vec{a}} e^{i\vec{q} \cdot \vec{a}} \right|^2 = \sum_{\vec{b}\vec{a}} e^{i\vec{q} \cdot \vec{b}\vec{a}}, \\ \frac{d\sigma}{d\Omega} &= \left(\frac{d\sigma}{d\Omega}\right)_{\text{Rutherford}} \left| \frac{1}{e} \int d^3 r \rho(\vec{r}) e^{i\vec{k} \cdot \vec{r}} \right|^2 \\ e^{i\vec{k} \cdot \vec{r}} &= \sum (2\ell + 1)i^4 j_\ell(kr) P_\ell(\cos \theta), \\ P_\ell(\cos \theta) &= \sqrt{\frac{4\pi}{2\ell + 1}} Y_{\ell,m=0}(\theta, \phi), \\ P_0(x) &= 1, \ P_1(x) = x, \ P_2(x) = (3x^2 - 1)/3, \\ f(\Omega) &= \sum_{\ell} (2\ell + 1)e^{i\delta_\ell} \sin \delta_\ell \frac{1}{k} P_\ell(\cos \theta) \\ \psi_{\vec{k}}(\vec{r})|_{R\rightarrow\infty} &= e^{i\vec{k} \cdot \vec{r}} + \frac{e^{ikr}}{r} f(\Omega), \\ \frac{d\sigma}{d\Omega} &= |f(\Omega)|^2, \ \sigma &= \frac{4\pi}{k^2} \sum_{\ell} (2\ell + 1)\sin^2 \delta_\ell, \ \delta \approx -ak \\ L_{\pm}|\ell,m\rangle &= \sqrt{\ell(\ell+1) - m(m\pm 1)}|\ell,m\pm 1\rangle, \\ \langle \vec{\beta}, J, M|T_q^k|\beta, \ell, m_\ell\rangle &= \langle JM|k, q, \ell, m_\ell\rangle \frac{\langle \vec{\beta}, J||T^{(k)}||\beta, \ell, J)}{\sqrt{2J+1}}, \\ n &= \frac{(2s+1)}{(2\pi)^d} \int_{k < k_f} d^d k, \ d \text{ dimensions}, \\ \{\Psi_s(\vec{x}), \Psi_s^k(\vec{y})\} &= \delta^3(\vec{x} - \vec{y})\delta_s', \ \{\Psi_s(\vec{x}), a_\alpha^{\perp}\} = \phi_{\alpha,s}(\vec{x}). \end{cases}$$

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- 1. (10 pts) Consider spin-vector operators, \vec{L} , \vec{S} and $\vec{J} = \vec{L} + \vec{S}$. You can assume \vec{S} operates on intrinsic spin and that \vec{L} describes orbital angular momentum. Circle the operators that commute with J_z ?
 - L_x
 - S_z
 - J_x
 - $J_x^2 + J_y^2 + J_z^2$
 - $S_x^2 + S_y^2 + S_z^2$
 - $\vec{L} \cdot \vec{S}$
- 2. (5 pts) A particle of mass m moving in one-dimension feels the potential

$$V(x) = \left\{egin{array}{ccc} \infty, & x < 0 \ -V_0, & 0 < x < a \ 0, & x > a \end{array}
ight.$$

Here V_0 is a positive constant. Assume that m, V_0 and a are initially chosen such that there exists one, but only one, bound state.

Circle the actions that could increase the number of bound states

- Increasing the magnitude of V_0
- Increasing the distance \boldsymbol{a}
- Increasing the mass m.
- 3. (5 pts) A particle of mass m moving in one-dimension feels the potential

$$V(x) = \left\{egin{array}{cc} \infty, & x < 0 \ -V_0 \delta(x-a), & x > 0 \end{array}
ight.$$

Here V_0 is a positive constant. Assume that m, V_0 and a are initially chosen such that there exists one, but only one, bound state.

Circle the actions that could increase the number of bound states

- Increasing the magnitude of V_0
- Increasing the distance \boldsymbol{a}
- Increasing the mass m.

4. A proton and a neutron are in the ground state of a harmonic oscillator. An interaction is added,

$$V_{
m s.s.} = -\mu_n B S_{nz} - \mu_p B S_{pz},$$

where S_{nz} and S_{pz} are the spin projection operators for the proton and neutron. We will use $|JM\rangle$ to denote a state of total angular momentum J and projection M. At t = 0 one is in the state $|J = 0, M = 0\rangle$. Find the probability you are in the following states as a function of time.

- (a) (5 pts) $|J = 0, M = 0\rangle$
- (b) (5 pts) $|J = 1, M = 1\rangle$
- (c) (5 pts) $|J = 1, M = 0\rangle$
- (d) (5 pts) $|J = 1, M = -1\rangle$

Advice: To simplify the algebra you might wish to express your work in terms of $\beta_n \equiv \mu_n B/2$ and $\beta_p \equiv \mu_p B/2$, which have units of frequency. (Extra work space for #4)

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5. A beam of spinless particles of mass m and kinetic energy E is aimed at a spherically symmetric repulsive potential

$$V(r) = \left\{egin{array}{cc} V_0, & r < a \ 0, & r > a \end{array}
ight.$$

Assume $E < V_0$.

- (a) (15 pts) Find the $\ell = 0$ phase shift as a function of the incoming wave number k.
- (b) (10 pts) What is the cross section as $k \to 0$?

(Extra work space for #5)

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6. A particle of mass m moves in a one-dimensional potential

 $V(x) = \alpha |x|,$

where α is a positive constant. NOTE THE ABSOLUTE VALUE!!!! Perform a variational estimate using a Gaussian form for a trial wave function,

$$\langle x|b
angle = \psi_b(x) = rac{1}{(\pi b^2)^{1/4}}e^{-x^2/2b^2},$$

where **b** is the variational parameter. In terms of m, b and α ,

- (a) (10 pts) What is $\langle b | KE | b \rangle$? –the expectation of the kinetic energy.
- (b) (10 pts) What is $\langle \boldsymbol{b} | \boldsymbol{V} | \boldsymbol{b} \rangle$? –the expectation of the potential energy.
- (c) (15 pts) What is the estimate of the ground state energy?

Advice: The dimensions of α are energy/length if you are checking dimensions.

(Extra work space for #6)

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