



The role of nuclear radii in beta decay processes

Chien-Yeah Seng

University of Washington
and
FRIB, Michigan State University

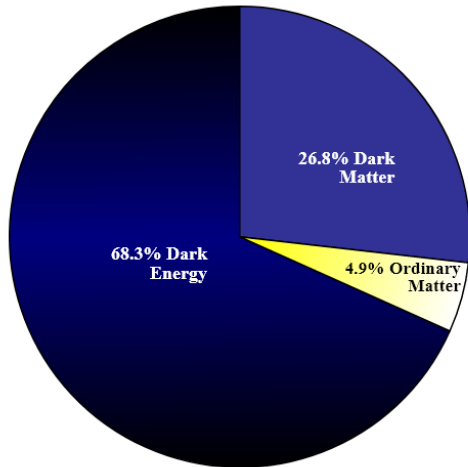
seng@frib.msu.edu

“Theoretical Justifications and Motivations for High-Profile FRIB Experiments”, FRIB

24 May, 2023

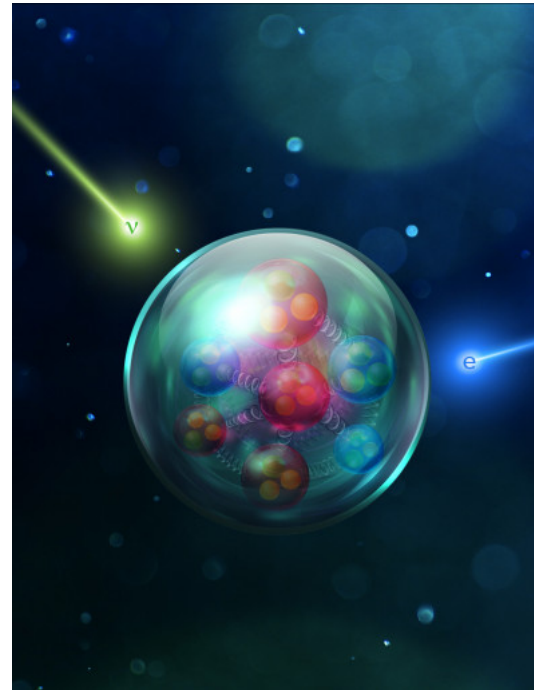
Many unresolved problems call for **physics beyond the Model (BSM) !**

Image credit: Wikipedia

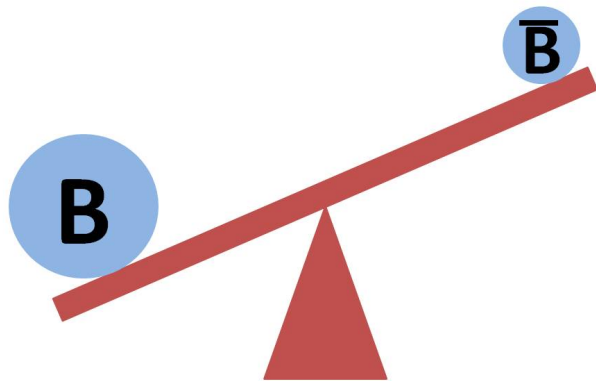


What is the origin of dark energy and dark matter?

Image credit: Jefferson Lab

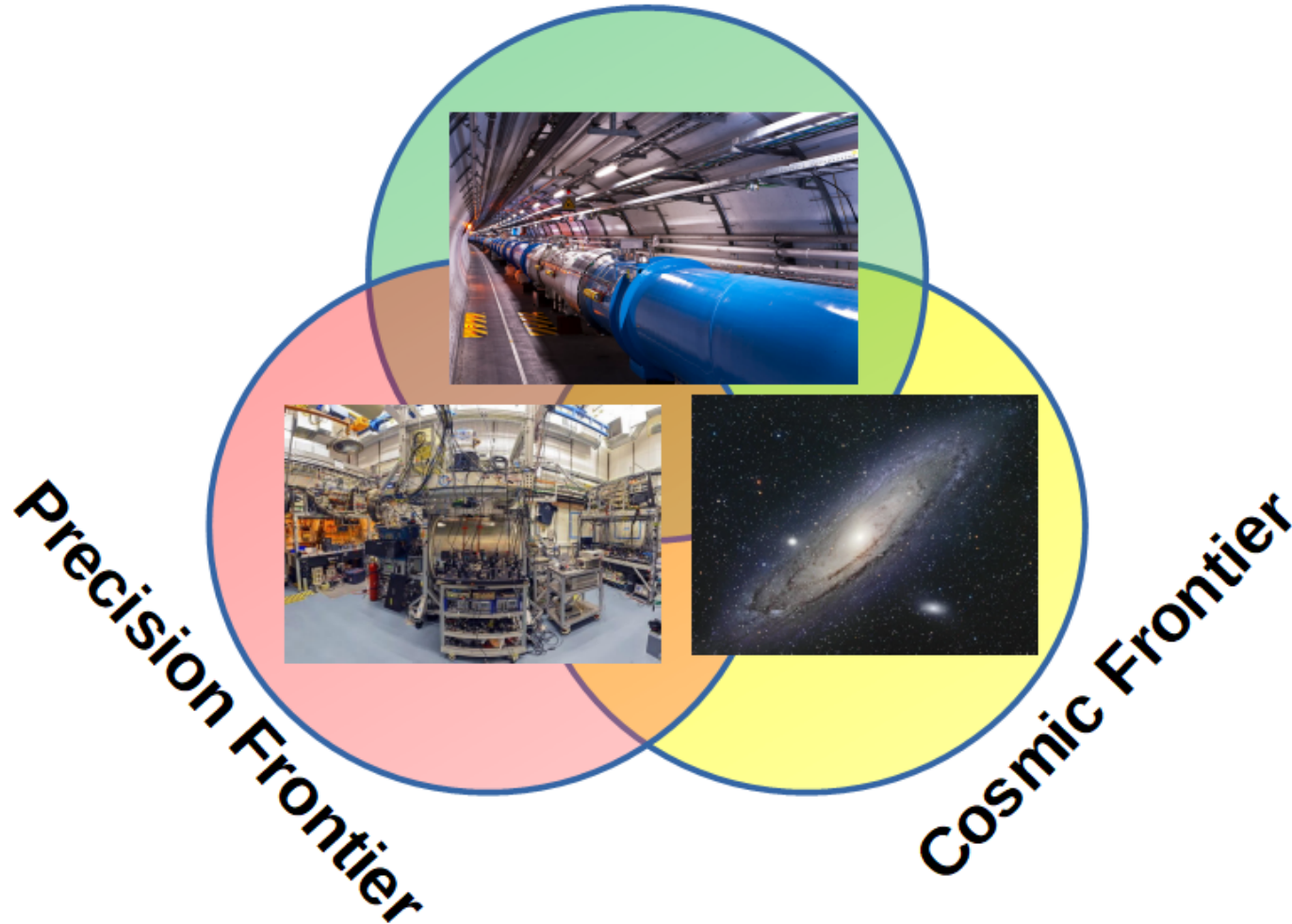


What is the nature of the neutrino mass?



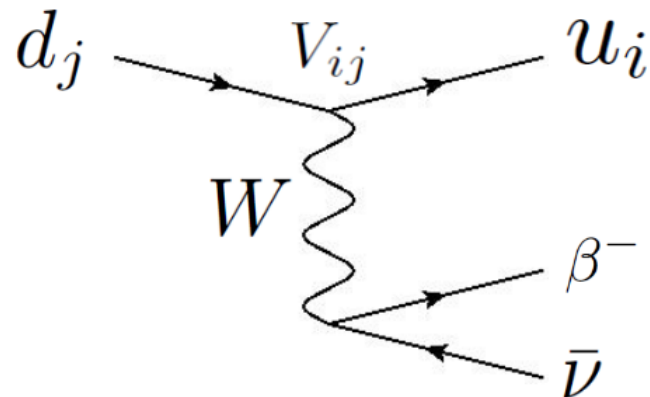
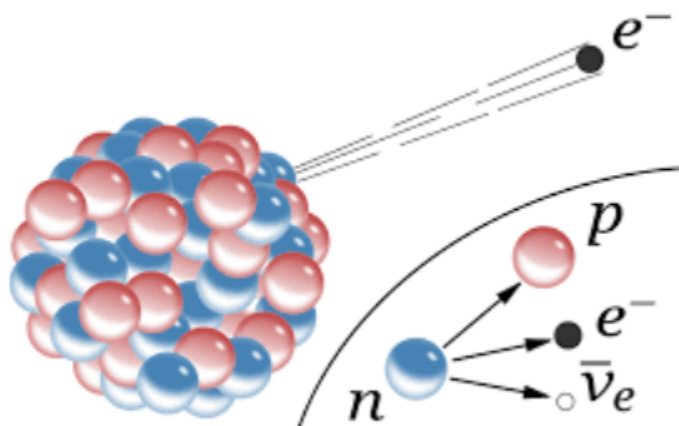
Why is there much more matter than antimatter in the observed universe?

Energy Frontier



Precision Frontier: Measure things very precisely, and look for their **deviations** from SM prediction!

Nuclear beta decay



Observables measured at **0.01%** level! Probes new physics at the scale:

$$\left(\frac{v_H}{\Lambda_{\text{BSM}}} \right)^2 \sim 0.01\% \implies \Lambda_{\text{BSM}} \sim 20 \text{ TeV}$$

Competitive to high-energy experiments at LHC!

Allowed beta decay: $\Delta J = 0, \pm 1$

$$\Delta\pi = 0$$

Lee-Yang Lagrangian

$$\begin{aligned}\mathcal{L}_{\text{LY}} = & -\bar{p}\gamma^\mu n (C_V^+ \bar{e}\gamma_\mu \nu_L + C_V^- \bar{e}\gamma_\mu \nu_R) - \bar{p}\gamma^\mu \gamma_5 n (C_A^+ \bar{e}\gamma_\mu \nu_L - C_A^- \bar{e}\gamma_\mu \nu_R) \\ & -\bar{p}n (C_S^+ \bar{e}\nu_L + C_S^- \bar{e}\nu_R) - \frac{1}{2}\bar{p}\sigma^{\mu\nu} n (C_T^+ \bar{e}\sigma_{\mu\nu} \nu_L + C_T^- \bar{e}\sigma_{\mu\nu} \nu_R) \\ & +\bar{p}\gamma_5 n (C_P^+ \bar{e}\nu_L - C_P^- \bar{e}\nu_R) + \text{h.c.}\end{aligned}$$

Scalar

Vector

Axial

Tensor



Fermi

$$\langle f | 1 | i \rangle$$



Gamow-Teller

$$\langle f | \sigma | i \rangle$$

SM: $C_V^+ = \sqrt{2}G_F V_{ud}g_V$, $C_A^+ = -\sqrt{2}G_F V_{ud}g_A$, other=0

Total decay rate:

$$\Gamma_i = (1 + \delta_i) \frac{M_F^2 m_e^5}{4\pi^3} f_V^i \hat{\xi}_i \left[1 + \gamma_i b_i \left\langle \frac{m_e}{E_e} \right\rangle_i \right]$$

Total decay rate:

$$\Gamma_i = (1 + \delta_i) \frac{M_F^2 m_e^5}{4\pi^3} f_V^i \hat{\xi}_i \left[1 + \gamma_i b_i \left\langle \frac{m_e}{E_e} \right\rangle_i \right]$$

Radiative corrections (RC) +
isospin-breaking corrections (ISB)

Phase-space factor +
nuclear/atomic
structure effects

Total decay rate:

$$\Gamma_i = (1 + \delta_i) \frac{M_F^2 m_e^5}{4\pi^3} f_V^i \hat{\xi}_i \left[1 + \underbrace{\gamma_i b_i \left\langle \frac{m_e}{E_e} \right\rangle}_i \right]$$

Radiative corrections (RC) +
isospin-breaking corrections (ISB)

SM + BSM couplings

Phase-space factor +
nuclear/atomic
structure effects

Total decay rate:

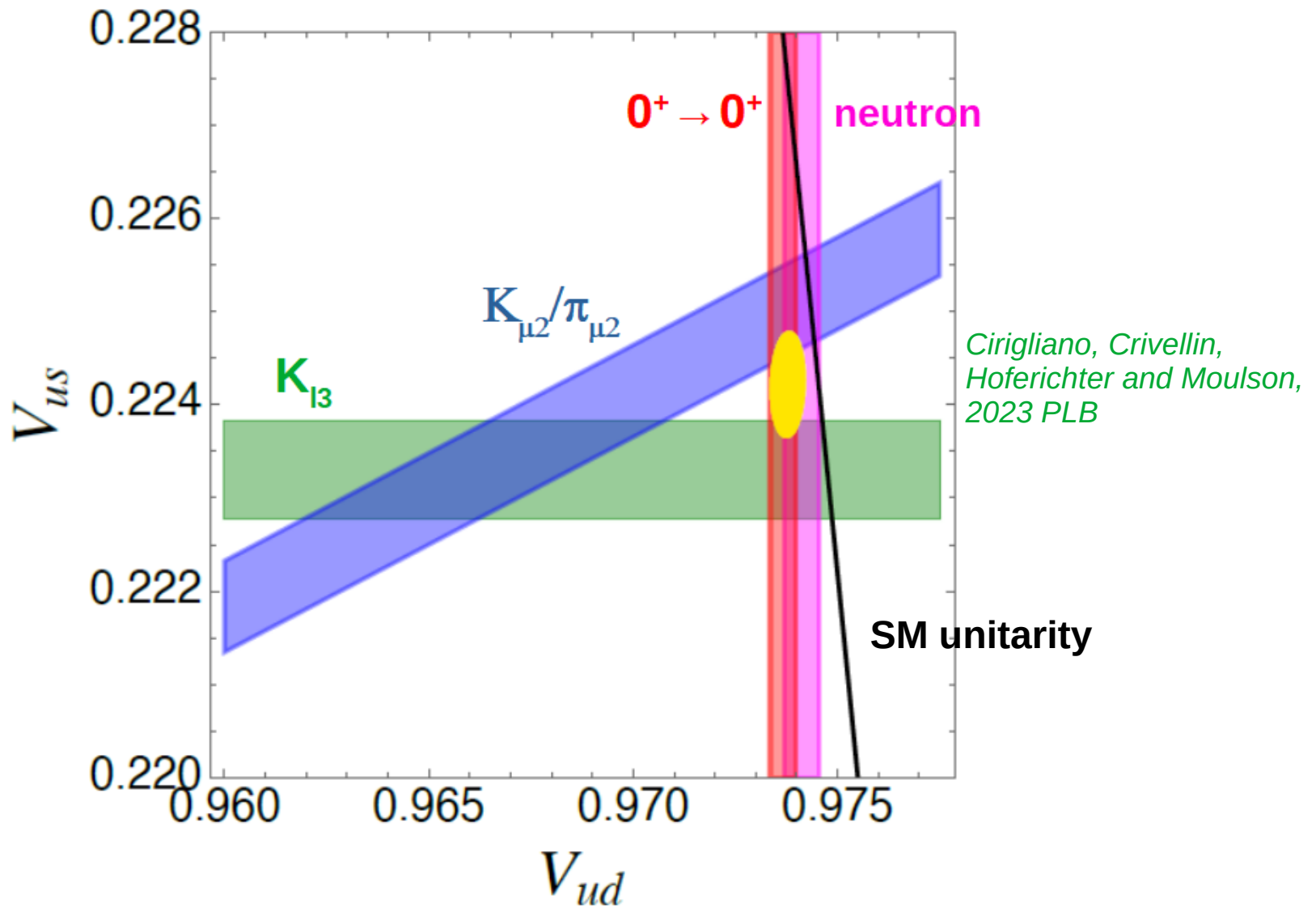
$$\Gamma_i = (1 + \delta_i) \frac{M_F^2 m_e^5}{4\pi^3} f_V^i \hat{\xi}_i \left[1 + \underbrace{\gamma_i b_i \left\langle \frac{m_e}{E_e} \right\rangle}_i \right]$$

Radiative corrections (RC) +
isospin-breaking corrections (ISB)

SM + BSM couplings

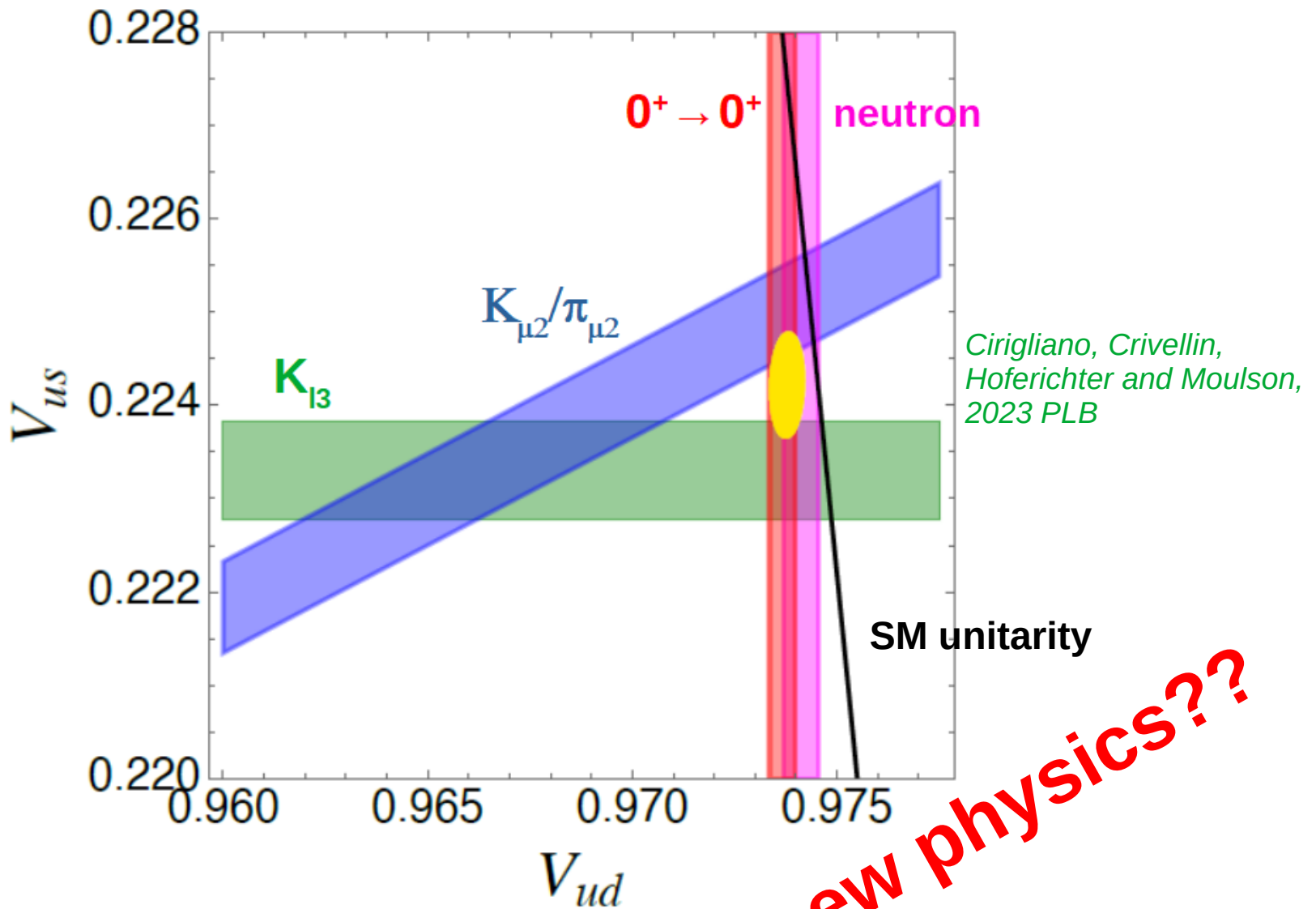
Phase-space factor +
nuclear/atomic
structure effects

- Assuming SM couplings: Deduce V_{ud}
- Variation of lifetime over nuclei: Constrain BSM couplings



$$|V_{ud}|_{0^+} = 0.97367(11)_{\text{exp}}(30)_{\text{th}}$$

$$|V_{ud}|_{0^+}^2 + |V_{us}|_{K_{\ell 3}}^2 - 1 = -0.0021(6)_{V_{ud}}(2)_{V_{us}}$$

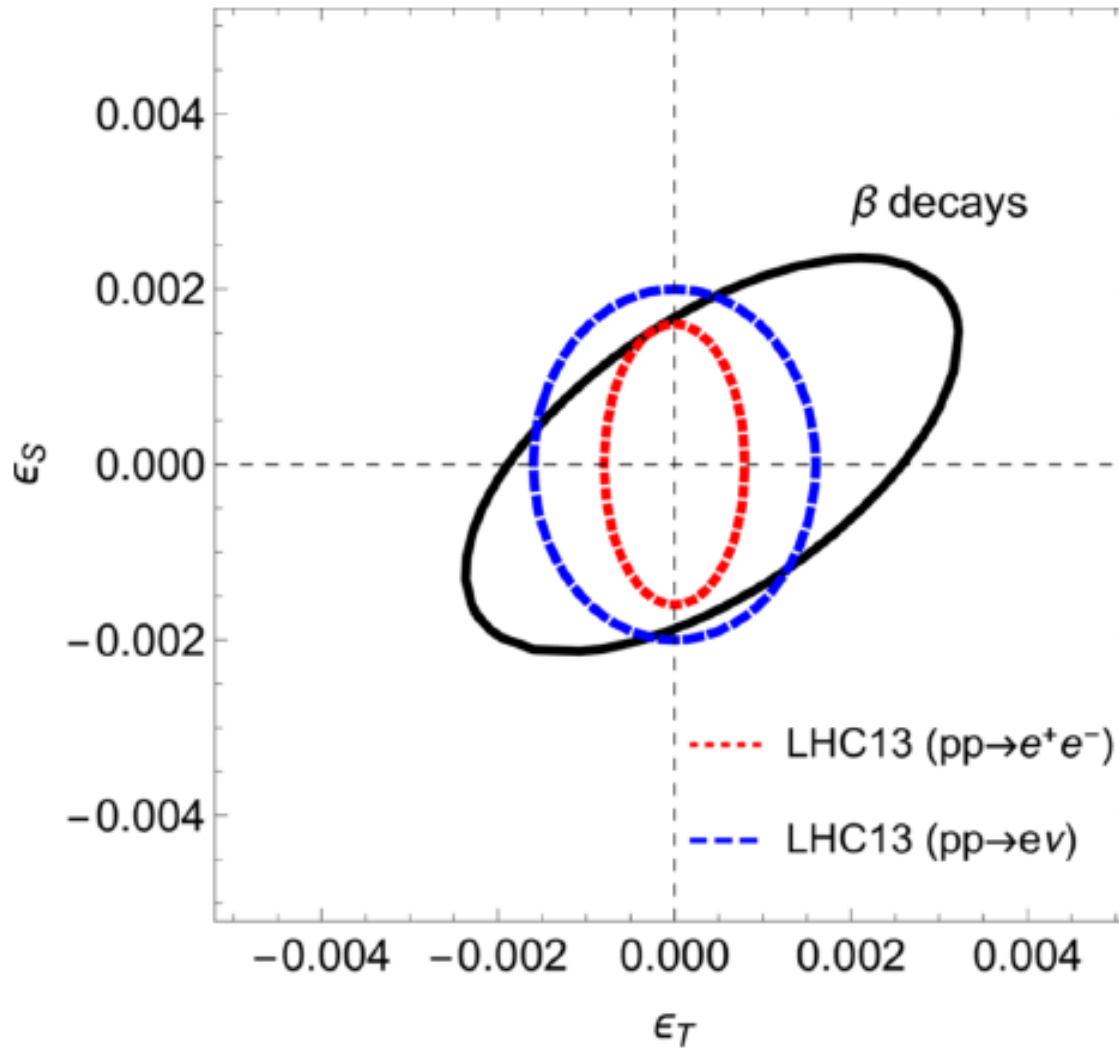


$$|V_{ud}|_{0^+} = 0.97367(11)_{\text{exp}}(30)_{\text{th}}$$

$$|V_{ud}|_{0^+}^2 + |V_{us}|_{K_{l3}}^2 - 1 = 0.0021(6)_{V_{ud}}(2)_{V_{us}}$$

Sign of new physics??

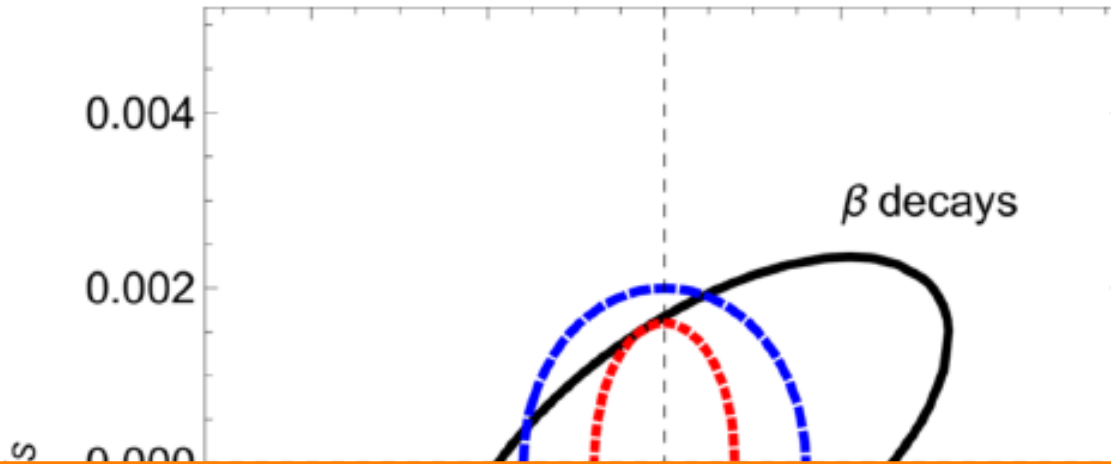
Interplay with LHC experiments:



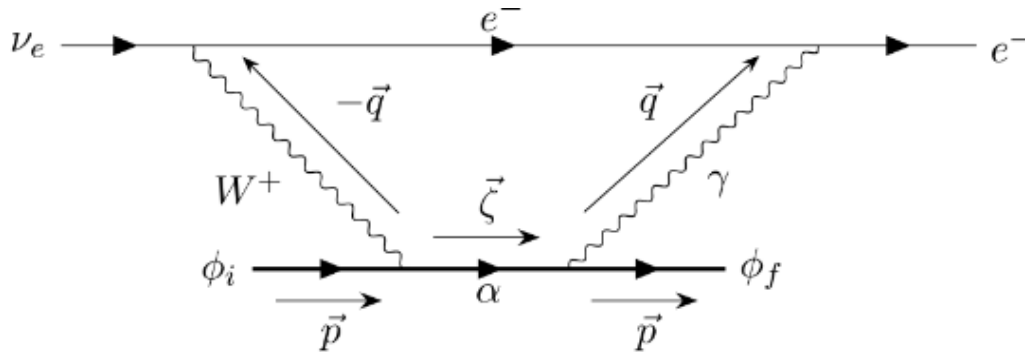
Falkowski, Gonzalez-Alonso and Naviliat-Cuncic, 2021 JHEP

$$C_{S,T}^+ = \sqrt{2}G_F V_{ud}g_S\epsilon_{S,T}$$

Interplay with LHC experiments:



Falkowski, Gonzalez-Alonso and Naviliat-Cuncic, 2021 JHEP



New radiative correction calculation could change the game!

See Petr Navratil's talk

ϵ_T

$$C_{S,T}^+ = \sqrt{2}G_F V_{ud} g_S \epsilon_{S,T}$$

Role of charge radii measurements:

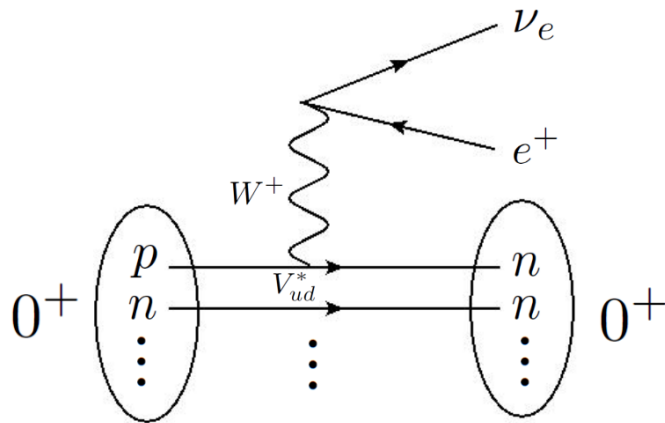
- 1. Beta decay form factors**
- 2. Isospin-symmetry-breaking corrections**

Role of charge radii measurements:

- 1. Beta decay form factors**
2. Isospin-symmetry-breaking corrections

“Superallowed” beta decays of $T=1, J^p=0^+$ nuclei

$$i(0^+) \rightarrow f(0^+) + e^+ + \nu_e$$



Provides the **best measurement of V_{ud}** :

- **23** measured transitions
- **15** with lifetime precision better than **0.23%**

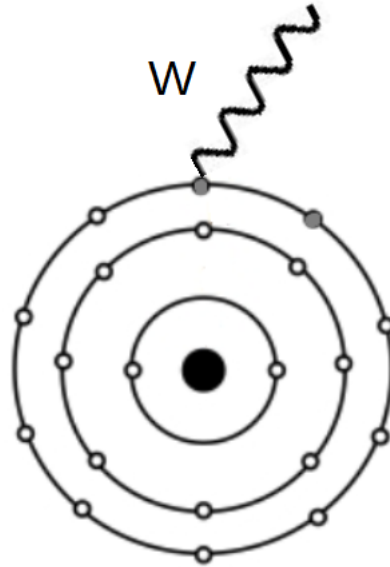
Hardy and Towner, 2020 PRC

$T_z = -1$
${}_{6}^{10}\text{C} \rightarrow {}_{5}^{10}\text{B}$
${}_{8}^{14}\text{O} \rightarrow {}_{7}^{14}\text{N}$
${}_{10}^{18}\text{Ne} \rightarrow {}_{9}^{18}\text{F}$
${}_{12}^{22}\text{Mg} \rightarrow {}_{11}^{22}\text{Na}$
${}_{14}^{26}\text{Si} \rightarrow {}_{13}^{26}\text{Al}$
${}_{16}^{30}\text{S} \rightarrow {}_{15}^{30}\text{P}$
${}_{18}^{34}\text{Ar} \rightarrow {}_{17}^{34}\text{Cl}$
${}_{20}^{38}\text{Ca} \rightarrow {}_{19}^{38}\text{K}$
${}_{22}^{42}\text{Ti} \rightarrow {}_{21}^{42}\text{Sc}$
${}_{24}^{46}\text{Cr} \rightarrow {}_{23}^{46}\text{V}$
${}_{26}^{50}\text{Fe} \rightarrow {}_{25}^{50}\text{Mn}$
${}_{28}^{54}\text{Ni} \rightarrow {}_{27}^{54}\text{Co}$

$T_z = 0$
${}_{13}^{26m}\text{Al} \rightarrow {}_{12}^{26}\text{Mg}$
${}_{17}^{34}\text{Cl} \rightarrow {}_{16}^{34}\text{S}$
${}_{19}^{38m}\text{K} \rightarrow {}_{18}^{38}\text{Ar}$
${}_{21}^{42}\text{Sc} \rightarrow {}_{20}^{42}\text{Ca}$
${}_{23}^{46}\text{V} \rightarrow {}_{22}^{46}\text{Ti}$
${}_{25}^{50}\text{Mn} \rightarrow {}_{24}^{50}\text{Cr}$
${}_{27}^{54}\text{Co} \rightarrow {}_{26}^{54}\text{Fe}$
${}_{31}^{62}\text{Ga} \rightarrow {}_{30}^{62}\text{Zn}$
${}_{33}^{66}\text{As} \rightarrow {}_{32}^{66}\text{Ge}$
${}_{35}^{70}\text{Br} \rightarrow {}_{34}^{70}\text{Se}$
${}_{37}^{74}\text{Rb} \rightarrow {}_{36}^{74}\text{Kr}$

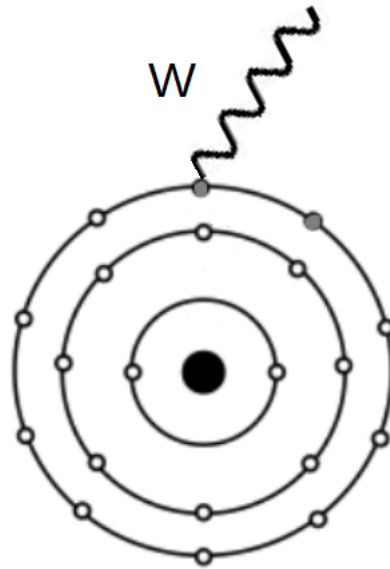
Nuclear (charged) weak distribution

$\rho_w(r)$: Distribution of “active” nucleons eligible for weak transitions in a nucleus



Nuclear (charged) weak distribution

$\rho_w(r)$: Distribution of “active” nucleons eligible for weak transitions in a nucleus



Old : Impulse approximation in **non-relativistic (NR) nuclear models**.

$$\langle f|O|i\rangle = \sum_{\alpha\beta} \langle \alpha|O|\beta\rangle \langle f|a_{\alpha}^{\dagger}a_{\beta}|i\rangle$$

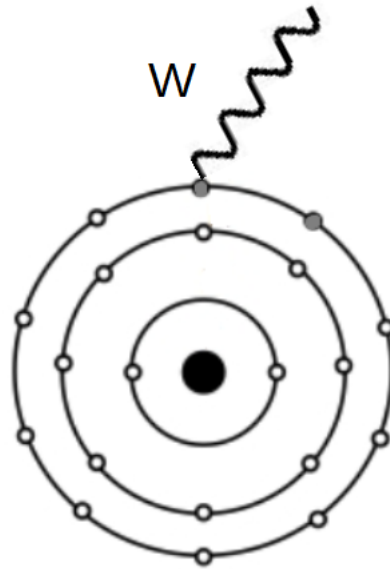
α, β : single-nucleon states

Hardy and Towner, 2005 PRC

How reliable is such approximation at **0.01% level**?

Nuclear (charged) weak distribution

$\rho_w(r)$: Distribution of “active” nucleons eligible for weak transitions in a nucleus



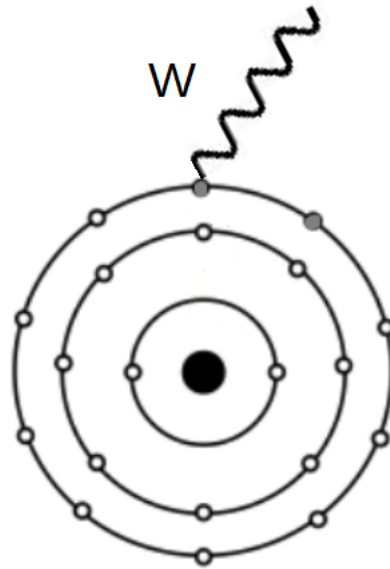
New: Utilize **CVC** to relate (charged) weak distribution to nuclear **charge distributions**

Hostein, 1974 RMP; CYS, 2023 PRL

$$\begin{aligned}\rho_w(r) &= \rho_{\text{Ch},1}(r) + Z_0 (\rho_{\text{Ch},0}(r) - \rho_{\text{Ch},1}(r)) \\ &= \rho_{\text{Ch},1}(r) + \frac{Z_{-1}}{2} (\rho_{\text{Ch},-1}(r) - \rho_{\text{Ch},1}(r))\end{aligned}$$

Nuclear (charged) weak distribution

$\rho_w(r)$: Distribution of “active” nucleons eligible for weak transitions in a nucleus



enhanced by large atomic number!

New: Utilize **CVC** to relate (charged) weak distribution to nuclear **charge distributions**

Hostein, 1974 RMP; CYS, 2023 PRL

$$\begin{aligned}\rho_w(r) &= \rho_{\text{Ch},1}(r) + Z_0 (\rho_{\text{Ch},0}(r) - \rho_{\text{Ch},1}(r)) \\ &= \rho_{\text{Ch},1}(r) + \frac{Z_{-1}}{2} (\rho_{\text{Ch},-1}(r) - \rho_{\text{Ch},1}(r))\end{aligned}$$

Form factor expansion: $F(q^2) = 1 + \frac{q^2}{6} \langle r^2 \rangle + \dots$

M.S. radius

Form factor expansion: $F(q^2) = 1 + \frac{q^2}{6} \langle r^2 \rangle + \dots$
M.S. radius

Existing charge radii compilation:

- Fricke and Heilig, “Nuclear charge radii”, 2004
- Angeli and Marinova, Atom.Data Nucl.Data Tabl. 99 (2013) 69
- Li, Luo and Wang, Atom.Data Nucl.Data Tabl. 140 (2021) 101440
- New measurements, e.g. ^{38}Ca , $^{38\text{m}}\text{K}$, ^{54}Ni ...

Form factor expansion: $F(q^2) = 1 + \frac{q^2}{6} \langle r^2 \rangle + \dots$
M.S. radius

Existing charge radii compilation:

- Fricke and Heilig, “Nuclear charge radii”, 2004
- Angeli and Marinova, Atom.Data Nucl.Data Tabl. 99 (2013) 69
- Li, Luo and Wang, Atom.Data Nucl.Data Tabl. 140 (2021) 101440
- New measurements, e.g. ^{38}Ca , $^{38\text{m}}\text{K}$, ^{54}Ni ...

Data + isospin predicts **unexpectedly large (charged) weak radii!**

$$\begin{array}{l} \text{E.g. } \langle r_{\text{Ch}}^2 \rangle_{^{38}\text{Ar}}^{1/2} = 3.4028(19) \text{ fm} \\ \langle r_{\text{Ch}}^2 \rangle_{^{38}\text{Ca}}^{1/2} = 3.467(1) \text{ fm} \end{array} \quad \longrightarrow \quad \langle r_w^2 \rangle^{1/2} = 4.00(4) \text{ fm}$$

Simplified shell model prediction: $\langle r_w^2 \rangle^{1/2} \approx 3.62 \text{ fm}$

A	$R_{\text{Ch},-1}$ (fm)	$R_{\text{Ch},0}$ (fm)	$R_{\text{Ch},1}$ (fm)
10	^{10}C	$^{10}\text{B}(\text{ex})$	^{10}Be : 2.3550(170) ^a
14	^{14}O	$^{14}\text{N}(\text{ex})$	^{14}C : 2.50 25(87) ^a
18	^{18}Ne : 2.9714(76) ^a	$^{18}\text{F}(\text{ex})$	^{18}O : 2.77 26(56) ^a
22	^{22}Mg : 3.0691(89) ^b	$^{22}\text{Na}(\text{ex})$	^{22}Ne : 2.9525(40) ^a
26	^{26}Si	^{26m}Al	^{26}Mg : 3.0337(18) ^a
30	^{30}S	$^{30}\text{P}(\text{ex})$	^{30}Si : 3.1336(40) ^a
34	^{34}Ar : 3.3654(40) ^a	^{34}Cl	^{34}S : 3.2847(21) ^a
38	^{38}Ca : 3.467(1) ^c	^{38m}K : 3.437(4) ^d	^{38}Ar : 3.4028(19) ^a
42	^{42}Ti	^{42}Sc : 3.5702(238) ^a	^{42}Ca : 3.5081(21) ^a
46	^{46}Cr	^{46}V	^{46}Ti : 3.6070(22) ^a
50	^{50}Fe	^{50}Mn : 3.7120(196) ^a	^{50}Cr : 3.6588(65) ^a
54	^{54}Ni : 3.738(4) ^e	^{54}Co	^{54}Fe : 3.6933(19) ^a
62	^{62}Ge	^{62}Ga	^{62}Zn : 3.9031(69) ^b
66	^{66}Se	^{66}As	^{66}Ge
70	^{70}Kr	^{70}Br	^{70}Se
74	^{74}Sr	^{74}Rb : 4.1935(172) ^b	^{74}Kr : 4.1870(41) ^a

A	$R_{\text{Ch},-1}$ (fm)	$R_{\text{Ch},0}$ (fm)	$R_{\text{Ch},1}$ (fm)
10	^{10}C	$^{10}\text{B}(\text{ex})$	^{10}Be : 2.3550(170) ^a
14	^{14}O	$^{14}\text{N}(\text{ex})$	^{14}C : 2.50 25(87) ^a
18	^{18}Ne : 2.9714(76) ^a	$^{18}\text{F}(\text{ex})$	^{18}O : 2.77 26(56) ^a
22	^{22}Mg : 3.0691(89) ^b	$^{22}\text{Na}(\text{ex})$	^{22}Ne : 2.9525(40) ^a
26	^{26}Si	^{26m}Al	^{26}Mg : 3.0337(18) ^a
30	^{30}S	$^{30}\text{P}(\text{ex})$	^{30}Si : 3.1336(40) ^a
34	^{34}Ar : 3.3654(40) ^a	^{34}Cl	^{34}S : 3.2847(21) ^a
38	^{38}Ca : 3.467(1) ^c	^{38m}K : 3.437(4) ^d	^{38}Ar : 3.4028(19) ^a
42	^{42}Ti	^{42}Sc : 3.5702(238) ^a	^{42}Ca : 3.5081(21) ^a
46	^{46}Cr	^{46}V	^{46}Ti : 3.6070(22) ^a
50	^{50}Fe	^{50}Mn : 3.7120(196) ^a	^{50}Cr : 3.6588(65) ^a
54	^{54}Ni : 3.738(4) ^e	^{54}Co	^{54}Fe : 3.6933(19) ^a
62	^{62}Ge	^{62}Ga	^{62}Zn : 3.9031(69) ^b
66	^{66}Se	^{66}As	^{66}Ge
70	^{70}Kr	^{70}Br	^{70}Se
74	^{74}Sr	^{74}Rb : 4.1935(172) ^b	^{74}Kr : 4.1870(41) ^a

Sufficient data available

A	$R_{\text{Ch},-1}$ (fm)	$R_{\text{Ch},0}$ (fm)	$R_{\text{Ch},1}$ (fm)
10	^{10}C	$^{10}\text{B}(\text{ex})$	$^{10}\text{Be}: 2.3550(170)^{\text{a}}$
14	^{14}O	$^{14}\text{N}(\text{ex})$	$^{14}\text{C}: 2.5025(87)^{\text{a}}$
18	$^{18}\text{Ne}: 2.9714(76)^{\text{a}}$	$^{18}\text{F}(\text{ex})$	$^{18}\text{O}: 2.7726(56)^{\text{a}}$
22	$^{22}\text{Mg}: 3.0691(89)^{\text{b}}$	$^{22}\text{Na}(\text{ex})$	$^{22}\text{Ne}: 2.9525(40)^{\text{a}}$
26	^{26}Si	$^{26\text{m}}\text{Al}$	$^{26}\text{Mg}: 3.0337(18)^{\text{a}}$
30	^{30}S	$^{30}\text{P}(\text{ex})$	$^{30}\text{Si}: 3.1336(40)^{\text{a}}$
34	$^{34}\text{Ar}: 3.3654(40)^{\text{a}}$	^{34}Cl	$^{34}\text{S}: 3.2847(21)^{\text{a}}$
38	$^{38}\text{Ca}: 3.467(1)^{\text{c}}$	$^{38\text{m}}\text{K}: 3.437(4)^{\text{d}}$	$^{38}\text{Ar}: 3.4028(19)^{\text{a}}$
42	^{42}Ti	$^{42}\text{Sc}: 3.5702(238)^{\text{a}}$	$^{42}\text{Ca}: 3.5081(21)^{\text{a}}$
46	^{46}Cr	^{46}V	$^{46}\text{Ti}: 3.6070(22)^{\text{a}}$
50	^{50}Fe	$^{50}\text{Mn}: 3.7120(196)^{\text{a}}$	$^{50}\text{Cr}: 3.6588(65)^{\text{a}}$
54	$^{54}\text{Ni}: 3.738(4)^{\text{e}}$	^{54}Co	$^{54}\text{Fe}: 3.6933(19)^{\text{a}}$
62	^{62}Ge	^{62}Ga	$^{62}\text{Zn}: 3.9031(69)^{\text{b}}$
66	^{66}Se	^{66}As	^{66}Ge
70	^{70}Kr	^{70}Br	^{70}Se
74	^{74}Rb	$^{74}\text{Rb}: 4.1935(172)^{\text{b}}$	$^{74}\text{Kr}: 4.1870(41)^{\text{a}}$

One more measurement needed

Preliminary re-analysis of statistical rate functions:

Transition	f_{new}	f_{HT}	$\frac{f_{\text{new}} - f_{\text{HT}}}{f_{\text{new}}} (\%)$
$^{18}\text{Ne} \rightarrow ^{18}\text{F}$	134.62(0) _{rad} (0) _{shape} (17) _{Q_{EC}}	134.64(17) _{Q_{EC}}	-0.01(0) _{rad} (0) _{shape}
$^{22}\text{Mg} \rightarrow ^{22}\text{Na}$	418.27(1) _{rad} (1) _{shape} (13) _{Q_{EC}}	418.35(13) _{Q_{EC}}	-0.02(0) _{rad} (0) _{shape}
$^{34}\text{Ar} \rightarrow ^{34}\text{Cl}$	3409.89(16) _{rad} (18) _{shape} (25) _{Q_{EC}}	3410.85(25) _{Q_{EC}}	-0.03(0) _{rad} (1) _{shape}
$^{38}\text{Ca} \rightarrow ^{38m}\text{K}$	5327.49(14) _{rad} (36) _{shape} (31) _{Q_{EC}}	5328.88(31) _{Q_{EC}}	-0.03(0) _{rad} (1) _{shape}
$^{42}\text{Ti} \rightarrow ^{42}\text{Sc}$	7124.3(5.7) _{rad} (0.8) _{shape} (1.4) _{Q_{EC}}	7130.1(1.4) _{Q_{EC}}	-0.08(8) _{rad} (1) _{shape}
$^{50}\text{Fe} \rightarrow ^{50}\text{Mn}$	15053(18) _{rad} (3) _{shape} (60) _{Q_{EC}}	15060(60) _{Q_{EC}}	-0.04(12) _{rad} (2) _{shape}
$^{54}\text{Ni} \rightarrow ^{54}\text{Co}$	21137(3) _{rad} (1) _{shape} (52) _{Q_{EC}}	21137(57) _{Q_{EC}}	+0.00(2) _{rad} (0) _{shape}
$^{34}\text{Cl} \rightarrow ^{34}\text{S}$	1995.076(81) _{rad} (103) _{shape} (94) _{Q_{EC}}	1996.003(96) _{Q_{EC}}	-0.05(0) _{rad} (1) _{shape}
$^{38m}\text{K} \rightarrow ^{38}\text{Ar}$	3296.32(8) _{rad} (21) _{shape} (15) _{Q_{EC}}	3297.39(15) _{Q_{EC}}	-0.03(0) _{rad} (1) _{shape}
$^{42}\text{Sc} \rightarrow ^{42}\text{Ca}$	4468.53(3.36) _{rad} (0.52) _{shape} (0.46) _{Q_{EC}}	4472.46(46) _{Q_{EC}}	-0.09(8) _{rad} (1) _{shape}
$^{50}\text{Mn} \rightarrow ^{50}\text{Cr}$	10737.93(11.50) _{rad} (2.02) _{shape} (0.50) _{Q_{EC}}	10745.99(49) _{Q_{EC}}	-0.08(11) _{rad} (2) _{shape}
$^{54}\text{Co} \rightarrow ^{54}\text{Fe}$	15769.4(2.3) _{rad} (0.7) _{shape} (2.7) _{Q_{EC}}	15766.8(2.7) _{Q_{EC}}	+0.02(1) _{rad} (0) _{shape}
$^{74}\text{Rb} \rightarrow ^{74}\text{Kr}$	47326(127) _{rad} (18) _{shape} (94) _{Q_{EC}}	47281(93) _{Q_{EC}}	+0.10(27) _{rad} (4) _{shape}

CYS, PRELIMINARY RESULTS

Shifts at 0.01% level are observed!

Upward shift of V_{ud} ?

Role of charge radii measurements:

1. Beta decay form factors
2. **Isospin-symmetry-breaking corrections**

Isospin symmetry breaking (**ISB**) correction alters the Fermi matrix element:

$$|M_F|^2 = |\langle f | \hat{\tau}_+ | i \rangle|^2 = |M_F^0|^2 (1 - \delta_C)$$

Caused by **isospin mixing** of nuclear states, predominantly due to Coulomb repulsion between protons

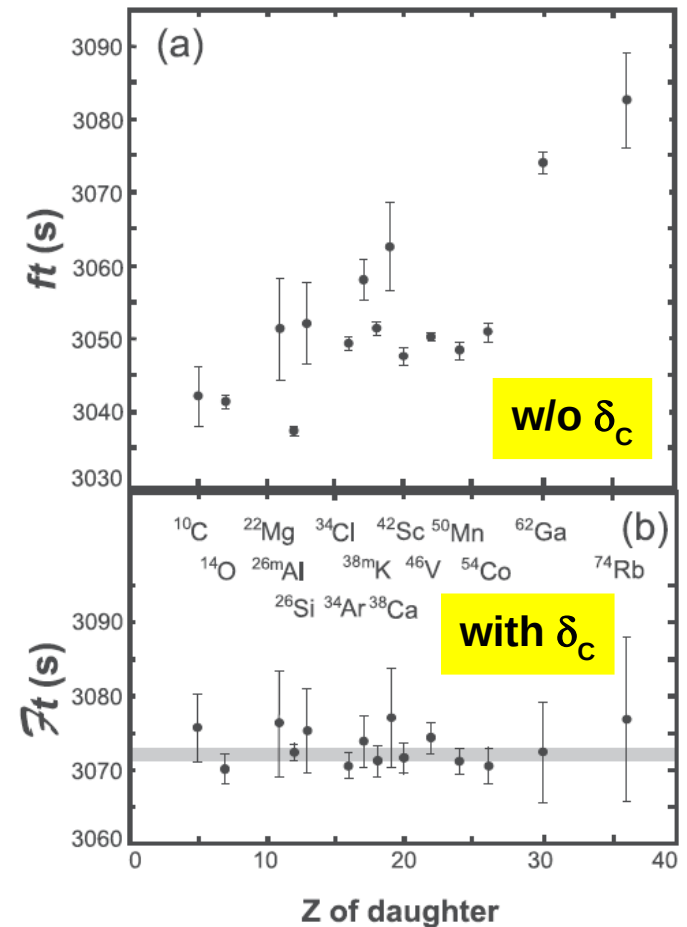
$$H = H_0 + \textcircled{V} \text{ ISB}$$

$$\delta_C = \mathcal{O}(V^2) \quad \mathbf{0.1\% \sim 1\%}$$

Crucial in obtaining a nucleus-independent Ft-value from the nucleus-dependent ft-values:

$$|V_{ud}|^2 \textcircled{\mathcal{F}t} (1 + \Delta_R^V) = 2984.43 \text{ s}$$

$$\mathcal{F}t = ft (1 + \delta'_R) (1 + \delta_{NS} - \delta_C)$$



Hardy and Towner, 2020 PRC

- Computing δ_c : Classic problem over **6 decades!** *MacDonald, 1958 Phys.Rev*
- Current input adopted in global analysis: **Shell model + Woods-Saxon (WS) potential**
- Successful in aligning Ft values of different superallowed transitions

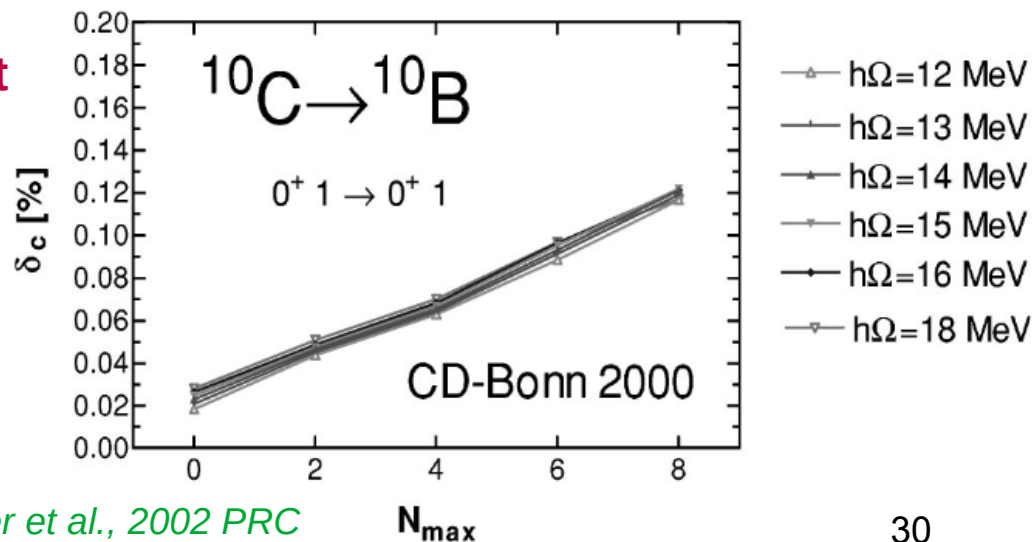
Hardy and Towner, 2020 PRC

Transitions	δ_c (%)				
	WS	DFT	HF	RPA	Micro
$^{26m}\text{Al} \rightarrow ^{26}\text{Mg}$	0.310	0.329	0.30	0.139	0.08
$^{34}\text{Cl} \rightarrow ^{34}\text{S}$	0.613	0.75	0.57	0.234	0.13
$^{38m}\text{K} \rightarrow ^{38}\text{Ar}$	0.628	1.7	0.59	0.278	0.15
$^{42}\text{Sc} \rightarrow ^{42}\text{Ca}$	0.690	0.77	0.42	0.333	0.18
$^{46}\text{V} \rightarrow ^{46}\text{Ti}$	0.620	0.563	0.38	/	0.21
$^{50}\text{Mn} \rightarrow ^{50}\text{Cr}$	0.660	0.476	0.35	/	0.24
$^{54}\text{Co} \rightarrow ^{54}\text{Fe}$	0.770	0.586	0.44	0.319	0.28

(Selected results)

Caveats:

- Significant **model dependence**.
- **No direct experimental constraint**
- Ab-initio calculations still in preliminary stages



Caurier et al., 2002 PRC

Probing isospin mixing in experiment?

Nuclear mass splitting: $\delta E_n = {}_0\langle n|V|n\rangle_0$ **X**

$$|n\rangle = |n\rangle_0 + \underbrace{\sum_{k \neq n} |k\rangle_0 \frac{{}_0\langle k|V|n\rangle_0}{E_n^{(0)} - E_k^{(0)}}}_{\text{We need this}} + \dots$$

We need this

Second-class current: doubly-suppressed (kinematics*ISB), null results so far *Minamisono et al., 2011 PRC*

$$V_\mu = \bar{u}(p_2) \left(g_V \gamma_\mu - i \frac{g_M - g_V}{2M} \sigma_{\mu\nu} q^\nu + \frac{g_S}{2M} q_\mu \right) u(p_1),$$

$$A_\mu = -\bar{u}(p_2) \gamma_5 \left(g_A \gamma_\mu - i \frac{g_T}{2M} \sigma_{\mu\nu} q^\nu + \frac{g_P}{2M} q_\mu \right) u(p_1).$$

Isospin mixing can be probed through the **variation of charge radii within the same nuclear isomultiplet**

Consider T=1 system. The non-zerosness of

$$\Delta M_B^{(1)} \equiv \frac{1}{2} (Z_{+1} r_{\text{ch},+1}^2 + Z_{-1} r_{\text{ch},-1}^2) - Z_0 r_{\text{ch},0}^2$$

signifies ISB. No double suppression! *Seng and Gorchtein, 2023 PLB*

Simple isovector-monopole-dominance picture: $\delta_C \propto \Delta M_B^{(1)}$
Auerbach, 2009 PRC

Therefore, the ability to reproduce $\Delta M_B^{(1)}$ serves as a test of the theory accuracy of δ_C

Challenge: **THREE** charge radii need to be measured!

A	$R_{\text{Ch},-1}$ (fm)	$R_{\text{Ch},0}$ (fm)	$R_{\text{Ch},1}$ (fm)
10	^{10}C	$^{10}\text{B}(\text{ex})$	^{10}Be : 2.3550(170) ^a
14	^{14}O	$^{14}\text{N}(\text{ex})$	^{14}C : 2.50 25(87) ^a
18	^{18}Ne : 2.9714(76) ^a	$^{18}\text{F}(\text{ex})$	^{18}O : 2.77 26(56) ^a
22	^{22}Mg : 3.0691(89) ^b	$^{22}\text{Na}(\text{ex})$	^{22}Ne : 2.9525(40) ^a
26	^{26}Si	^{26m}Al	^{26}Mg : 3.0337(18) ^a
30	^{30}S	$^{30}\text{P}(\text{ex})$	^{30}Si : 3.1336(40) ^a
34	^{34}Ar : 3.3654(40) ^a	^{34}Cl	^{34}S : 3.2847(21) ^a
38	^{38}Ca : 3.467(1) ^c	^{38m}K : 3.437(4) ^d	^{38}Ar : 3.4028(19) ^a
42	^{42}Ti	^{42}Sc : 3.5702(238) ^a	^{42}Ca : 3.5081(21) ^a
46	^{46}Cr	^{46}V	^{46}Ti : 3.6070(22) ^a
50	^{50}Fe	^{50}Mn : 3.7120(196) ^a	^{50}Cr : 3.6588(65) ^a
54	^{54}Ni : 3.738(4) ^e	^{54}Co	^{54}Fe : 3.6933(19) ^a
62	^{62}Ge	^{62}Ga	^{62}Zn : 3.9031(69) ^b
66	^{66}Se	^{66}As	^{66}Ge
70	^{70}Kr	^{70}Br	^{70}Se
74	^{74}Sr	^{74}Rb : 4.1935(172) ^b	^{74}Kr : 4.1870(41) ^a

A	$R_{\text{Ch},-1}$ (fm)	$R_{\text{Ch},0}$ (fm)	$R_{\text{Ch},1}$ (fm)
10	^{10}C	$^{10}\text{B}(\text{ex})$	^{10}Be : 2.3550(170) ^a
14	^{14}O	$^{14}\text{N}(\text{ex})$	^{14}C : 2.50 25(87) ^a
18	^{18}Ne : 2.9714(76) ^a	$^{18}\text{F}(\text{ex})$	^{18}O : 2.77 26(56) ^a
22	^{22}Mg : 3.0691(89) ^b	$^{22}\text{Na}(\text{ex})$	^{22}Ne : 2.9525(40) ^a
26	^{26}Si	^{26m}Al	^{26}Mg : 3.0337(18) ^a
30	^{30}S	$^{30}\text{P}(\text{ex})$	^{30}Si : 3.1336(40) ^a
34	^{34}Ar : 3.3654(40) ^a	^{34}Cl	^{34}S : 3.2847(21) ^a
38	^{38}Ca : 3.467(1) ^c	^{38m}K : 3.437(4) ^d	^{38}Ar : 3.4028(19) ^a
42	^{42}Ti	^{42}Sc : 3.5702(238) ^a	^{42}Ca : 3.5081(21) ^a

$$A = 38 \quad : \quad \Delta M_B^{(1)} = -0.03(12)_+(52)_0(7)_- \text{fm}^2$$

Possible future measurement of $r(^{38m}\text{K})$ at TRIUMF with ~4 times better in precision

A	$R_{\text{Ch},-1}$ (fm)	$R_{\text{Ch},0}$ (fm)	$R_{\text{Ch},1}$ (fm)
10	^{10}C	$^{10}\text{B}(\text{ex})$	^{10}Be : 2.3550(170) ^a
14	^{14}O	$^{14}\text{N}(\text{ex})$	^{14}C : 2.50 25(87) ^a
18	^{18}Ne : 2.9714(76) ^a	$^{18}\text{F}(\text{ex})$	^{18}O : 2.77 26(56) ^a
22	^{22}Mg : 3.0691(89) ^b	$^{22}\text{Na}(\text{ex})$	^{22}Ne : 2.9525(40) ^a
26	^{26}Si	^{26m}Al	^{26}Mg : 3.0337(18) ^a
30	^{30}S	$^{30}\text{P}(\text{ex})$	^{30}Si : 3.1336(40) ^a
34	^{34}Ar : 3.3654(40) ^a	^{34}Cl	^{34}S : 3.2847(21) ^a
38	^{38}Ca : 3.467(1) ^c	^{38m}K : 3.437(4) ^d	^{38}Ar : 3.4028(19) ^a
42	^{42}Ti	^{42}Sc : 3.5702(238) ^a	^{42}Ca : 3.5081(21) ^a
46	^{46}Cr	^{46}V	^{46}Ti : 3.6070(22) ^a
50	^{50}Fe	^{50}Mn : 3.7120(196) ^a	^{50}Cr : 3.6588(65) ^a
54	^{54}Ni : 3.738(4) ^e	^{54}Co	^{54}Fe : 3.6933(19) ^a
62	^{62}Ge	^{62}Ga	^{62}Zn : 3.9031(69) ^b
66	^{66}Se	^{66}As	^{66}Ge
70	^{70}Kr	^{70}Br	^{70}Se
74	^{74}Sr	^{74}Rb : 4.1935(172) ^b	^{74}Kr : 4.1870(41) ^a

One more measurement needed

Summary

- **Nuclear beta decays** offer precision tests of SM; **various anomalies** suggest the existence of new physics
- Measurements of **nuclear charge radii** provide important nuclear structure information relevant to beta decays:
 - (1) The (vector) decay form factor, determined by **TWO measured radii** through CVC;
 - (2) ISB correction in Fermi transitions, probed by **THREE measured radii**
- New data will critically impact the measurement of V_{ud} and the **search of new physics** (BSM couplings)