

# Final Exam Review

\* Wed. 5/1 3:00-5:00 pm (Nonprogrammable calculator)

\* Exam is comprehensive

\* I'll provide some formulas (see example exam)

\* Here I list the review topics of material from exam 3 on  
(for topics on exams 1-3, see earlier review notes)  
posted on web page

1.) Leftover Central Force material not on exam 3

\* material from §8.7 (Planetary Motion/Kepler's Laws)

$$\theta(r) = \int \frac{l/r^2 dr}{\sqrt{2\mu(E + \frac{K}{r} - \frac{l^2}{2\mu r^2})}} + \text{const}$$

↓ integrated

$$\frac{\alpha}{r} = 1 + \varepsilon \cos\theta$$

\* normalized so  $r(\theta=0) = r_{\min}$

$$\alpha \equiv \frac{l^2}{\mu K}$$
$$\varepsilon = \sqrt{1 + \frac{2El^2}{\mu K^2}}$$

## Types of conic sections for orbits

①  $\epsilon > 1$  ( $E > 0$ )  $\Rightarrow$  Hyperbola

②  $\epsilon = 1$  ( $E = 0$ )  $\Rightarrow$  Parabola

③  $0 < \epsilon < 1$  ( $V_{\min}^{\text{orb}} < E < 0$ )  $\Rightarrow$  Ellipse

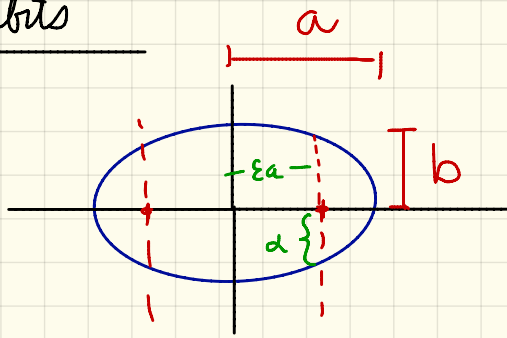
④  $\epsilon = 0$  ( $E = V_{\min}^{\text{orb}}$ )  $\Rightarrow$  Circle

} unbound

} bound

\* Understand figs. 8-8 + 8-9 as to how these look

# Elliptical Orbits



\* Sun sits at focus

$$a = \frac{\alpha}{1 - \epsilon^2} = \frac{K}{2|E|}$$

$$b = \frac{\alpha}{\sqrt{1 - \epsilon^2}} = \frac{l}{\sqrt{2m|E|}}$$

$$r_{\min} = a(1 - \epsilon) = \frac{\alpha}{1 + \epsilon}$$

$$r_{\max} = a(1 + \epsilon) = \frac{\alpha}{1 - \epsilon}$$

$$\text{Period } \tau = \pi K \sqrt{\frac{m}{2}} |E|^{-3/2} = \sqrt{\frac{4\pi^2 m}{K} a^3} \quad \left( \begin{array}{l} \text{derive from} \\ \frac{dA}{dt} = \frac{l}{2m} \end{array} \right)$$

## Kepler's Laws

- ① Planets move in ellipses w/ Sun at one of the foci
- ②  $\frac{dA}{dt} = \text{const} = \frac{h}{2\mu}$
- ③ Period  $T^2 \propto a^3$

## ② Lagrangian Mechanics

\* generalized coordinates (e.g., N-bodies  $\Rightarrow S = 3N$  degrees of freedom)

$$L = T - U$$

Can write  $L = L(\mathbf{r}_{\alpha,i}, \dot{\mathbf{r}}_{\alpha,i})$   $\begin{matrix} \alpha = 1, \dots, N \\ i = x, y, z \end{matrix}$

or more generally,

$$L = L(q_j, \dot{q}_j) \quad j = 1, \dots, S \quad (S = 3N \text{ here})$$

where  $q_j = q_j(\mathbf{r}_{\alpha,i})$  and  $\mathbf{r}_{\alpha,i} = \mathbf{r}_{\alpha,i}(q_1, \dots, q_S)$

"generalized coords"

Key Point: Lagrange eqns. of motion look the same  
whether using  $r_x, r_y, r_z$  or  $q_1, q_2, q_3$

eg: 1 particle in 3d

$$\text{could use } L = L(x, y, z, \dot{x}, \dot{y}, \dot{z})$$

$$\frac{\partial L}{\partial x} - \frac{d}{dt} \frac{\partial L}{\partial \dot{x}} = 0 \quad (\text{+ ditto for } y, z)$$

or

$$\text{could use } L = L(q_1, q_2, q_3, \dot{q}_1, \dot{q}_2, \dot{q}_3)$$

$$\Rightarrow \frac{\partial L}{\partial q_1} - \frac{d}{dt} \frac{\partial L}{\partial \dot{q}_1} = 0 \quad \text{etc}$$

## \* Hamilton's Principle

The physical trajectories  $q_1(t), \dots, q_s(t)$  of a system

Minimize the integral

$$J[q_1, q_2, \dots, q_s] = \int_{t_1}^{t_2} L(q_1, \dots, q_s; \dot{q}_1, \dot{q}_2, \dots, \dot{q}_s; t) dt$$

Euler eqns. follow from

$$\frac{dJ}{d\alpha} = 0$$

$$\text{where } q_i(t; \alpha) = q_i(t) + \alpha \eta_i(t)$$

$$\eta_i(t_1) = \eta_i(t_2) = 0$$



\* See book & class examples for setting up  $L$  in generalized coords

## Cyclic Coordinates

$$\text{If } \frac{\partial L}{\partial q_j} = 0 \Rightarrow \frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}_j} \right) = 0$$

$$\Rightarrow \frac{\partial L}{\partial \dot{q}_j} = \text{conserved (constant)}$$

" $q_j$  is a cyclic coordinate"

Generalized Force ;  $Q_j = \frac{\partial L}{\partial q_j}$

Generalized Momentum :  $\tilde{p}_j = \frac{\partial L}{\partial \dot{q}_j}$

## Probs. w/ constraints

eg, particle on some surface

$$L = L(q_i, \dot{q}_i) \quad i = 1 \dots S$$

AND

$$f_k(q_i, t) = 0 \quad k = 1 \dots m$$

\* 2 ways to deal w/ these problems

① Use the  $m$  constraint eqns to eliminate  $m$  of the  $q_i$

$$L(q_1, q_2, \dots, q_s) \rightarrow L(q_1, q_2, \dots, q_{s-m})$$

Solve:  $\frac{\partial L}{\partial q_i} - \frac{d}{dt} \frac{\partial L}{\partial \dot{q}_i} \quad i = 1 \dots (s-m)$

Easy, but no way to find force of constraint

## ② Lagrange Multipliers

\* treat all  $S$  DOF as independent but solve modified Euler-Lagrange eqns +  $f_k = 0$

$$\frac{\partial L}{\partial q_i} - \frac{d}{dt} \frac{\partial L}{\partial \dot{q}_i} + \sum_{k=1}^m \lambda_k(t) \frac{\partial f_k}{\partial q_i} = 0 \quad i=1, \dots, S$$

AND

$$f_k(q_1, \dots, q_s) = 0 \quad k=1, \dots, m$$

$(S+m)$  eqns for  $S+m$  unknowns  $(q_1, \dots, q_s) + (\lambda_1, \dots, \lambda_m)$

$\lambda_k(t) = \text{"Lagrange Multiplier"}$

# Generalized Forces of Constraint

$$Q_j^{\text{constraint}} = \sum_{k=1}^m \lambda_k \frac{\partial f_k}{\partial q_j}$$