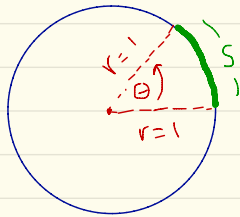


* A quick review on Solid Angles

Solid angle Ω - used to quantify a range of directions in 3D

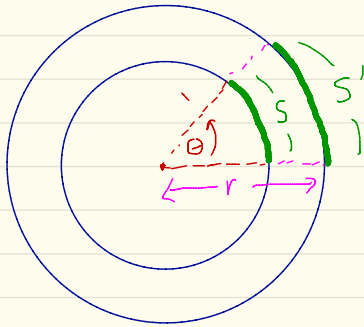
* Take a step back & recall what an angle θ means in 2D.

Consider a unit ($r=1$) circle



$$\theta = s$$

* more generally, for a circle of radius $r \neq 1$



$$\theta = s = \frac{s'}{r}$$

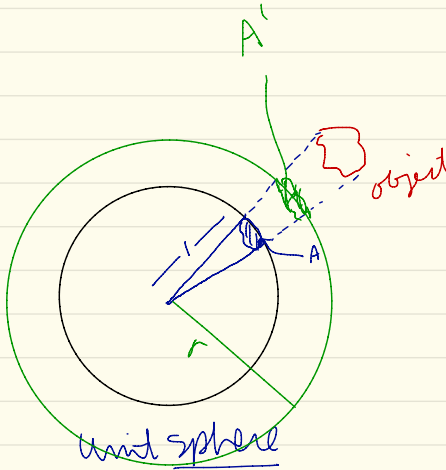
(radius)

\Downarrow

$$\theta_{\max} = 2\pi$$

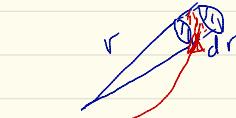
Solid Angle:

Quantifies range of directions in 3D, e.g., for seeing an object



$$\begin{aligned}\text{Solid angle } \Omega &= A \quad (\text{area on unit sphere}) \\ &= \frac{A'}{r^2} \quad (\text{area on any sphere})\end{aligned}$$

Differential Solid angle $d\Omega$



$$\begin{aligned}d^3r &= r^2 \sin\theta d\theta d\phi dr \\ &= dA dr\end{aligned}$$

dA = differential of area on sphere of radius r

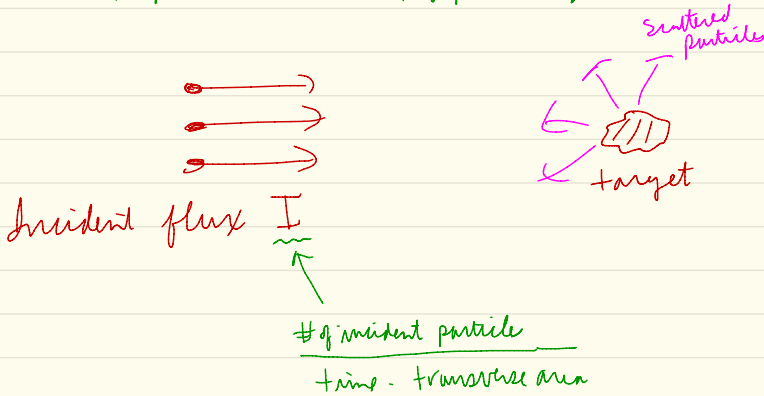
$$\therefore d\Omega = \frac{dA}{r^2} = \sin\theta d\theta d\phi$$

= range of directions cut out by $d\theta$ + $d\phi$ on unit sphere

$$\Omega_{\max} = \int_0^\pi \sin\theta d\theta \int_0^{2\pi} d\phi = (2\pi)(-\cos\theta)|_0^\pi = 4\pi \quad (\text{area of unit sphere})$$

Scattering Cross Section (statistical description of collisions)

* Describes capability of a target to subject to incident particles to some fate (e.g., scatter into some direction. In QM, more exotic "fates" possible, e.g., incident particles can be absorbed & "new" particles emerge.)



* Let $dN \equiv \#$ of projectile particles meeting some fate due to interaction w/target (e.g., getting deflected into some range of directions)

* Physically, expect $dN \propto I$

$$dN = I d\sigma$$

$\#$ of particles/time meeting some fate (e.g., getting deflected in some range of directions)

Coefficient of proportionality "cross section" ($[d\sigma] = \text{area}$)

$$dN = I d\sigma = I \frac{d\sigma}{d\Omega} d\Omega$$



if the "rate" is to be scattered into some range of directions specified by $d\Omega$.

$$\Rightarrow \frac{dN}{d\Omega} = I \frac{d\sigma}{d\Omega}$$

deflected into $d\Omega$ about some direction per + time

"Differential cross section"

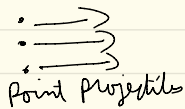
* Total scattering cross section

$$\sigma_{\text{t}} = \int d\sigma = \int \frac{d\sigma}{d\Omega} d\Omega = \frac{1}{I} \int dN = \frac{N}{I}$$

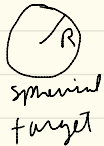
$$\Rightarrow N = I \sigma_{\text{t}} = \# \text{ incident particles per unit time affected (scattered) by the target}$$

* If we are considering contact forces between target/projectile, then σ_{t} just common sense geometry

ex 1:



Point Projectiles



spherical target

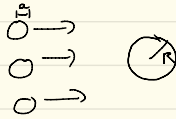
$$\Rightarrow \sigma_{\text{t}} = \pi R^2$$

Side Remark: If we look at the quantum mechanical σ_x for the previous system, we find A) $\sigma_x^{qm} \neq \sigma_x^{cm}$, and moreover $\sigma_x^{qm} = \sigma_x^{qm}(\bar{v})$.

e.g., low E $\sigma_x^{qm} \rightarrow 4 \times \sigma_x^{cm}$

(de Broglie wavelength $\lambda = \frac{h}{p}$
 QM effects make projectiles "effectively" bigger)

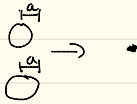
ex 2: finite projectile



$$\sigma_x \neq \pi R^2$$

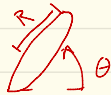
$$\sigma_x = \pi (R+a)^2$$

ex 3!



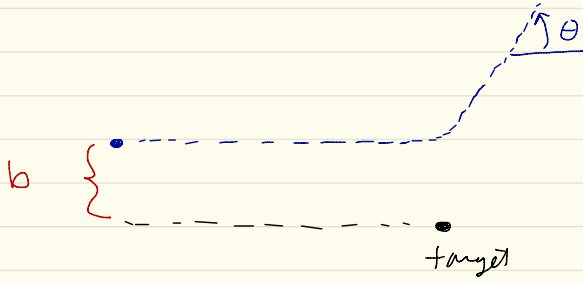
$$\sigma_x = \pi a^2$$

ex 4!



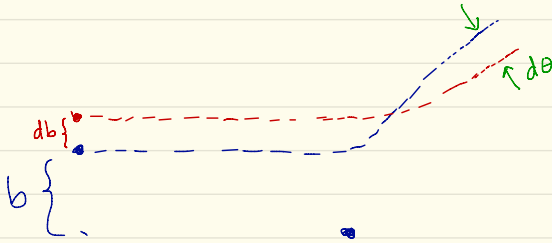
$$\sigma_x = \pi R^2 \sin^2 \theta$$

* Scattering Cross Sections for Central Forces (no ϕ -dependence)



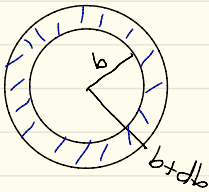
b = "impact parameter"
 (asymptotic \perp distance
 between projectile & target)

* Clearly, we expect $\theta = \theta(b)$. Consider some db about b



NOTE: Here we make the reasonable assumption that the central force is such that deflection decreases w/ increasing b

i.e., $\frac{db}{d\theta} < 0$



View looking down the
 beam of projectiles

$$\text{let } dN = \frac{\# \text{ projectiles scattered } \theta \text{ thru } d\theta}{\text{Sec}}$$

$$= I \cdot (2\pi b db) \quad \text{— area of shaded region}$$

$$= I d\sigma(\theta)$$

$$\text{so, } d\sigma = 2\pi b db = 2\pi b \left| \frac{db}{d\theta} \right| d\theta$$

← ab. value since by assumption
 $d\sigma > 0$ and $\frac{db}{d\theta} < 0$ and $d\theta > 0$.

$$\Rightarrow \frac{d\sigma(\theta)}{d\theta} = 2\pi b \left| \frac{db}{d\theta} \right|$$

* relating back to $\frac{d\sigma}{d\Omega}$ we use

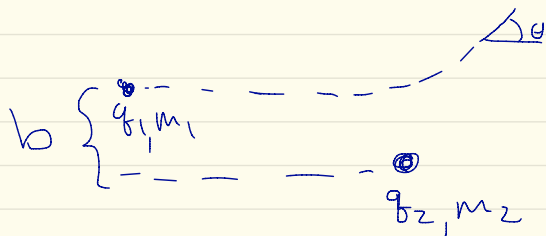
$$d\Omega = d\phi \sin\theta d\theta = 2\pi \sin\theta d\theta$$

(axially symmetric)

$$\therefore \frac{d\sigma(\theta)}{d\Omega} = \frac{1}{2\pi \sin\theta} \frac{d\sigma}{d\theta} = \frac{b}{\sin\theta} \left| \frac{db}{d\theta} \right|$$

↑
 general result for
 differential cross section
 per solid angle for
 central force (i.e., $d\sigma(\theta, \phi) = d\sigma(\theta)$)
 & axially symmetric problems.

Example: Coulomb repulsion scattering



$$\vec{F} = k \frac{q_1 q_2}{r^3} \vec{r}$$

(assume $m_2 \gg m_1$)

Claim : $b(\theta) = \frac{kq_1q_2}{2T} \cot \frac{\theta}{2}$ $T = \frac{1}{2} m_1 v_i^2$

$$\frac{db}{d\theta} = \frac{kq_1q_2}{4T} \frac{1}{\sin^2 \frac{\theta}{2}}$$

$$\therefore \frac{d\sigma}{d\Omega} = \frac{b}{\sin\theta} \left| \frac{db}{d\theta} \right| = \frac{kq_1q_2}{2T} \frac{\cos \frac{\theta}{2}}{\sin \frac{\theta}{2}} \times \frac{1}{\sin\theta} \frac{kq_1q_2}{4T} \frac{1}{\sin^2 \frac{\theta}{2}}$$

$$= \frac{kq_1q_2}{2T} \frac{\cancel{\cos \frac{\theta}{2}}}{\sin \frac{\theta}{2}} \times \frac{1}{2 \sin \frac{\theta}{2} \cancel{\cos \frac{\theta}{2}}} \times \frac{kq_1q_2}{4T} \frac{1}{\sin^2 \frac{\theta}{2}}$$

$$= \frac{1}{16} \left(\frac{kq_1q_2}{T} \right)^2 \frac{1}{\sin^4 \frac{\theta}{2}} \quad \text{"Rutherford cross section"}$$

Remarkably, in a lucky accident due to the $\frac{1}{r^2}$ form of the Coulomb force, the CM + GM expressions agree. This allowed Rutherford to deduce the presence of the atomic nucleus in α -Gold scattering.