Friday, March 5, 1:50-2:40 PM
$$\int_{-\infty}^{\infty} dx \ e^{-x^2/(2a^2)} = a\sqrt{2\pi},$$

$$H = i\hbar\partial_t, \ \vec{P} = -i\hbar\nabla,$$

$$\sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix},$$

$$U(t, -\infty) = 1 + \frac{-i}{\hbar} \int_{-\infty}^{t} dt' \ V(t')U(t', -\infty),$$

$$\langle x|x'\rangle = \delta(x - x'), \ \langle p|p'\rangle = \frac{1}{2\pi\hbar} \delta(p - p'),$$

$$|p\rangle = \int dx \ |x\rangle e^{ipx/\hbar}, \quad |x\rangle = \int \frac{dp}{2\pi\hbar} |p\rangle e^{-ipx/\hbar},$$

$$H = \frac{P^2}{2m} + \frac{1}{2} m\omega^2 x^2 = \hbar\omega (a^{\dagger}a + 1/2),$$

$$a^{\dagger} = \sqrt{\frac{m\omega}{2\hbar}} X - i\sqrt{\frac{1}{2\hbar m\omega}} P,$$

$$\psi_0(x) = \frac{1}{(\pi b^2)^{1/4}} e^{-x^2/2b^2}, \quad b^2 = \frac{\hbar}{m\omega},$$

$$\rho(\vec{r}, t) = \psi^*(\vec{r}_1, t_1) \psi(\vec{r}_2, t_2)$$

$$\vec{J}(\vec{r}, t) = \frac{-i\hbar}{2m} (\psi^*(\vec{r}, t) \nabla \psi(\vec{r}, t) - (\nabla \psi^*(\vec{r}, t)) \psi(\vec{r}, t))$$

$$- \frac{e\vec{A}}{mc} |\psi(\vec{r}, t)|^2.$$

$$H = \frac{(\vec{P} - e\vec{A}/c)^2}{2m} + e\Phi,$$
For $V = \beta\delta(x - y) : -\frac{\hbar^2}{2m} \left(\frac{\partial}{\partial x} \psi(x)|_{y + \epsilon} - \frac{\partial}{\partial x} \psi(x)|_{y - \epsilon} \right) = -\beta\psi(y),$

$$\vec{E} = -\nabla \Phi - \frac{1}{c} \partial_t \vec{A}, \quad \vec{B} = \nabla \times \vec{A},$$

$$egin{aligned} \omega_{ ext{cyclotron}} &= rac{eB}{mc}, \ e^{A+B} &= e^A e^B e^{-C/2}, & ext{if } [A,B] = C, ext{ and } [C,A] = [C,B] = 0, \ Y_{0,0} &= rac{1}{\sqrt{4\pi}}, & Y_{1,0} &= \sqrt{rac{3}{4\pi}}\cos\theta, & Y_{1,\pm 1} &= \mp\sqrt{rac{3}{8\pi}}\sin\theta e^{i\pm\phi}, \ Y_{2,0} &= \sqrt{rac{5}{16\pi}}(3\cos^2\theta - 1), & Y_{2,\pm 1} &= \mp\sqrt{rac{15}{8\pi}}\sin\theta\cos\theta e^{\pm i\phi}, \ Y_{2,\pm 2} &= \sqrt{rac{15}{32\pi}}\sin^2\theta e^{\pm 2i\phi}, & Y_{\ell-m}(\theta,\phi) &= (-1)^mY_{\ell m}^*(\theta,\phi). \end{aligned}$$

$$|N\rangle = |n\rangle - \sum_{m \neq n} |m\rangle \frac{1}{\epsilon_m - \epsilon_n} \langle m|V|n\rangle + \cdots$$

$$E_N = \epsilon_n + \langle n|V|n\rangle - \sum_{m \neq n} \frac{|\langle m|V|n\rangle|^2}{\epsilon_m - \epsilon_n}$$

$$j_0(x) = \frac{\sin x}{x}, \ n_0(x) = -\frac{\cos x}{x}, \ j_1(x) = \frac{\sin x}{x^2} - \frac{\cos x}{x}, \ n_1(x) = -\frac{\cos x}{x^2} - \frac{\sin x}{x}$$

$$j_2(x) = \left(\frac{3}{x^3} - \frac{1}{x}\right) \sin x - \frac{3}{x^2} \cos x, \ n_2(x) = -\left(\frac{3}{x^3} - \frac{1}{x}\right) \cos x - \frac{3}{x^2} \sin x,$$

$$\frac{d}{dt} P_{i \rightarrow n}(t) = \frac{2\pi}{h} |V_{ni}|^2 \delta(E_n - E_i),$$

$$\frac{d\sigma}{d\Omega} = \frac{m^2}{4\pi^2 h^4} \left| \int d^3 r V(r) e^{i(\vec{k}_f - \vec{k}_i) \cdot \vec{r}} \right|^2,$$

$$\sigma = \frac{(2S_R + 1)}{(2S_1 + 1)(2S_2 + 1)} \frac{4\pi}{k^2} \frac{(h\Gamma_R/2)^2}{(\epsilon_n - \epsilon_r)^2 + (h\Gamma_R/2)^2},$$

$$\frac{d\sigma}{d\Omega} = \left(\frac{d\sigma}{d\Omega}\right)_{\text{single}} \tilde{S}(\vec{q}), \ \tilde{S}(\vec{q}) = \left| \sum_{\vec{a}} e^{i\vec{q} \cdot \vec{a}} \right|^2,$$

$$e^{i\vec{k} \cdot \vec{r}} = \sum_{\ell} (2\ell + 1)i^{\ell}j_{\ell}(kr)P_{\ell}(\cos\theta),$$

$$P_{\ell}(\cos\theta) = \sqrt{\frac{4\pi}{2\ell + 1}} Y_{\ell,m=0}(\theta, \phi),$$

$$P_0(x) = 1, \ P_1(x) = x, \ P_2(x) = (3x^2 - 1)/3,$$

$$f(\Omega) \equiv \sum_{\ell} (2\ell + 1)e^{i\delta_\ell}\sin\delta_\ell \frac{1}{k}P_{\ell}(\cos\theta)$$

$$\psi_{\vec{k}}(\vec{r})|_{R \rightarrow \infty} = e^{i\vec{k} \cdot \vec{r}} + \frac{e^{ikr}}{r} f(\Omega),$$

$$\frac{d\sigma}{d\Omega} = |f(\Omega)|^2, \quad \sigma = \frac{4\pi}{k^2} \sum_{\ell} (2\ell + 1)\sin^2\delta_\ell, \quad \delta \approx -ak$$

$$L_{\pm}|\ell, m\rangle = \sqrt{\ell(\ell + 1) - m(m \pm 1)}|\ell, m \pm 1\rangle,$$

$$C_{m,m,n,i,JM}^{\ell,s} = \langle \ell, s, J, M|\ell, s, m_t, m_s\rangle,$$

$$\langle \vec{\beta}, J, M|T_q^k|\beta, \ell, m_\ell\rangle = C_{qm,\epsilon;JM}^{k\ell} \frac{\langle \vec{\beta}, J||T^{(k)}||\beta, \ell, J\rangle}{\sqrt{2J + 1}},$$

$$n = \frac{(2s + 1)}{(2\pi)^d} \int_{k < k_f} d^d k, \quad d \text{ dimensions},$$

$$\{\Psi_s(\vec{x}), \Psi_s^{\ell}(\vec{y})\} = \delta^3(\vec{x} - \vec{y})\delta_{ss'},$$

$$\Psi_s^{\ell}(\vec{r}) = \frac{1}{\sqrt{V}} \sum_{r} e^{i\vec{k} \cdot \vec{r}_a^{\dagger}(\vec{k}), \quad \{\Psi_s(\vec{x}), a_\alpha^{\dagger}\} = \phi_{\alpha,s}(\vec{x}).$$

1. Type α, β and γ particles exist in a TWO-DIMENSIONAL world. The α particle has mass M_{α} and is described by the two-dimensional field operator within a large area A,

$$egin{align} \Phi_lpha(ec r,t) &= rac{1}{\sqrt{A}} \sum_{ec k} e^{-iE_kt/\hbar + iec k\cdot ec r} a_{ec k}, \ \Phi_lpha^\dagger(ec r,t) &= rac{1}{\sqrt{A}} \sum_{ec k} e^{iE_kt/\hbar - iec k\cdot ec r} a_{ec k}^\dagger, \ E_k &= M_lpha c^2 + rac{\hbar^2 k^2}{2M_lpha}. \end{align}$$

The β and γ particles are massless and described by the operators,

$$egin{align} \Psi_eta(ec{r},t) &= rac{1}{\sqrt{A}} \sum_{ec{q}} e^{-iE_qt/\hbar + iec{q}\cdotec{r}} b_q, \ \Psi_eta^\dagger(ec{r},t) &= rac{1}{\sqrt{A}} \sum_{ec{q}} e^{iE_qt/\hbar - iec{q}\cdotec{r}} b_q^\dagger, \ \Psi_\gamma(ec{r},t) &= rac{1}{\sqrt{A}} \sum_{ec{q}} e^{-iE_qt/\hbar + iec{q}\cdotec{r}} c_q, \ \Psi_\gamma^\dagger(ec{r},t) &= rac{1}{\sqrt{A}} \sum_{ec{q}} e^{iE_qt/\hbar - iec{q}\cdotec{r}} c_q^\dagger, \ E_q &= \hbar c q. \end{align}$$

The massive α particle can decay to an β and a γ particle via the interaction

$$H_{
m int} = g \int dx dy \; \left[\Phi_lpha(ec{r},t) \Psi_eta^\dagger(ec{r},t) \Psi_\gamma^\dagger(ec{r},t) + \Phi_lpha^\dagger(ec{r},t) \Psi_eta(ec{r},t) \Psi_\gamma(ec{r},t)
ight],$$

where the coupling constant g is small. The creation and destruction operators obey the commutation rules $[a_{\vec{k}}, a_{\vec{k}'}^{\dagger}] = \delta_{\vec{k}\vec{k}'}, [b_{\vec{q}}, b_{\vec{q}'}^{\dagger}] = \delta_{\vec{q}\vec{q}'}$ and $[c_{\vec{q}}, c_{\vec{q}'}^{\dagger}] = \delta_{\vec{q}\vec{q}'}$. FYI: In two dimensions $\delta_{\vec{k}\vec{k}'} = \delta_{k_x k'_x} \delta_{k_y k'_y}$.

- (a) (5 pts) Evaluate the commutator $[\Phi_{\alpha}(\vec{r},t),\Phi_{\alpha}^{\dagger}(\vec{r}',t)]$.
- (b) (5 pts) What is the dimensionality of g?
- (c) (20 pts) Calculate the rate at which an α particle at rest decays into a β and a γ particle.
- (d) (5 pts) How would your answer change if the $\boldsymbol{\beta}$ and $\boldsymbol{\gamma}$ particles were identical?

(Extra work space for #1)

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2. Imagine you had calculated the following matrix element,

$$\mathcal{M} = \langle \alpha, J_f = 1, M_f = 0 | z^2 | \beta, J_i = 3, M_i = 0 \rangle.$$

For the following matrix elements, first state whether they are zero, and if not, express them in terms of \mathcal{M} . Evaluate any Clebsch-Gordan coefficients in your answers. You can use tables or online resources to get the values. (5 pts. each)

•
$$\langle \alpha, J_f = 1, M_f = 0 | x^2 + y^2 | \beta, J_i = 3, M_i = 0 \rangle$$

•
$$\langle \alpha, J_f = 1, M_f = 0 | z^2 | \beta, J_i = 3, M_i = 2 \rangle$$

$$ullet$$
 $\langle lpha, J_f=1, M_f=0|xz|eta, J_i=3, M_i=0
angle$

$$ullet$$
 $\langle lpha, J_f=1, M_f=0|xz|eta, J_i=3, M_i=2
angle$

$$\bullet \ \langle \alpha, J_f = 1, M_f = 0 | xy | \beta, J_i = 3, M_i = 2 \rangle$$

(Extra work space for #2)

- 3. Electrons of mass m are placed in a three-dimensional spherically symmetric harmonic oscillator with potential, $V = m\omega^2 r^2/2$.
 - (a) (5 pts) What is the ground state energy for a single electron?
 - (b) (5 pts) What is the ground state energy if 20 electrons are in the well? (sum over all particles, $\epsilon_1 + \epsilon_2 + \cdots + \epsilon_{20}$)

Now add an interaction with an external magnetic field,

$$V_{eta} = -\mu ec{B} \cdot (ec{L} + 2ec{S}),$$

where the strength of the magnetic field is adjusted so that $\mu B = \omega$.

- (c) (10 pts) What is the new single-particle ground state energy?
- (d) (10 pts) What is the ground state energy if 20 electrons are in the well?

(Extra work space for #3)