

$$\int_{-\infty}^{\infty} dx e^{-x^2/(2a^2)} = a\sqrt{2\pi},$$

$$H = i\hbar\partial_t, \quad \vec{P} = -i\hbar\nabla,$$

$$\sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \quad \sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix},$$

$$U(t, -\infty) = 1 + \frac{-i}{\hbar} \int_{-\infty}^t dt' V(t')U(t', -\infty),$$

$$\langle x|x'\rangle = \delta(x - x'), \quad \langle p|p'\rangle = \frac{1}{2\pi\hbar}\delta(p - p'),$$

$$|p\rangle = \int dx |x\rangle e^{ipx/\hbar}, \quad |x\rangle = \int \frac{dp}{2\pi\hbar} |p\rangle e^{-ipx/\hbar},$$

$$H = \frac{P^2}{2m} + \frac{1}{2}m\omega^2 x^2 = \hbar\omega(a^\dagger a + 1/2),$$

$$a^\dagger = \sqrt{\frac{m\omega}{2\hbar}}X - i\sqrt{\frac{1}{2\hbar m\omega}}P,$$

$$\psi_0(x) = \frac{1}{(\pi b^2)^{1/4}} e^{-x^2/2b^2}, \quad b^2 = \frac{\hbar}{m\omega},$$

$$\rho(\vec{r}, t) = \psi^*(\vec{r}_1, t_1)\psi(\vec{r}_2, t_2)$$

$$\vec{j}(\vec{r}, t) = \frac{-i\hbar}{2m}(\psi^*(\vec{r}, t)\nabla\psi(\vec{r}, t) - (\nabla\psi^*(\vec{r}, t))\psi(\vec{r}, t))$$

$$- \frac{e\vec{A}}{mc}|\psi(\vec{r}, t)|^2.$$

$$H = \frac{(\vec{P} - e\vec{A}/c)^2}{2m} + e\Phi,$$

For $V = \beta\delta(x - y)$: $-\frac{\hbar^2}{2m}\left(\frac{\partial}{\partial x}\psi(x)|_{y+\epsilon} - \frac{\partial}{\partial x}\psi(x)|_{y-\epsilon}\right) = -\beta\psi(y),$

$$\vec{E} = -\nabla\Phi - \frac{1}{c}\partial_t\vec{A}, \quad \vec{B} = \nabla \times \vec{A},$$

$$\omega_{\text{cyclotron}} = \frac{eB}{mc},$$

$$e^{A+B} = e^A e^B e^{-C/2}, \quad \text{if } [A, B] = C, \text{ and } [C, A] = [C, B] = 0,$$

$$Y_{0,0} = \frac{1}{\sqrt{4\pi}}, \quad Y_{1,0} = \sqrt{\frac{3}{4\pi}}\cos\theta, \quad Y_{1,\pm 1} = \mp\sqrt{\frac{3}{8\pi}}\sin\theta e^{\pm i\phi},$$

$$Y_{2,0} = \sqrt{\frac{5}{16\pi}}(3\cos^2\theta - 1), \quad Y_{2,\pm 1} = \mp\sqrt{\frac{15}{8\pi}}\sin\theta\cos\theta e^{\pm i\phi},$$

$$Y_{2,\pm 2} = \sqrt{\frac{15}{32\pi}}\sin^2\theta e^{\pm 2i\phi}, \quad Y_{\ell-m}(\theta, \phi) = (-1)^m Y_{\ell m}^*(\theta, \phi).$$

$$|N\rangle = |n\rangle - \sum_{m \neq n} |m\rangle \frac{1}{\epsilon_m - \epsilon_n} \langle m|V|n\rangle + \dots$$

$$E_N = \epsilon_n + \langle n|V|n\rangle - \sum_{m \neq n} \frac{|\langle m|V|n\rangle|^2}{\epsilon_m - \epsilon_n}$$

$$j_0(x) = \frac{\sin x}{x}, \quad n_0(x) = -\frac{\cos x}{x}, \quad j_1(x) = \frac{\sin x}{x^2} - \frac{\cos x}{x}, \quad n_1(x) = -\frac{\cos x}{x^2} - \frac{\sin x}{x}$$

$$j_2(x) = \left(\frac{3}{x^3} - \frac{1}{x}\right) \sin x - \frac{3}{x^2} \cos x, \quad n_2(x) = -\left(\frac{3}{x^3} - \frac{1}{x}\right) \cos x - \frac{3}{x^2} \sin x,$$

$$\frac{d}{dt} P_{i \rightarrow n}(t) = \frac{2\pi}{\hbar} |V_{ni}|^2 \delta(E_n - E_i),$$

$$\frac{d\sigma}{d\Omega} = \frac{m^2}{4\pi^2 \hbar^4} \left| \int d^3r \mathcal{V}(r) e^{i(\vec{k}_f - \vec{k}_i) \cdot \vec{r}} \right|^2,$$

$$\sigma = \frac{(2S_R + 1) 4\pi (\hbar\Gamma_R/2)^2}{(2S_1 + 1)(2S_2 + 1) k^2 (\epsilon_k - \epsilon_r)^2 + (\hbar\Gamma_R/2)^2},$$

$$\frac{d\sigma}{d\Omega} = \left(\frac{d\sigma}{d\Omega} \right)_{\text{single}} \tilde{S}(\vec{q}), \quad \tilde{S}(\vec{q}) = \frac{1}{N} \left| \sum_{\vec{a}} e^{i\vec{q} \cdot \vec{a}} \right|^2,$$

$$e^{i\vec{k} \cdot \vec{r}} = \sum_{\ell} (2\ell + 1) i^{\ell} j_{\ell}(kr) P_{\ell}(\cos \theta),$$

$$P_{\ell}(\cos \theta) = \sqrt{\frac{4\pi}{2\ell + 1}} Y_{\ell, m=0}(\theta, \phi),$$

$$P_0(x) = 1, \quad P_1(x) = x, \quad P_2(x) = (3x^2 - 1)/3,$$

$$f(\Omega) \equiv \sum_{\ell} (2\ell + 1) e^{i\delta_{\ell}} \sin \delta_{\ell} \frac{1}{k} P_{\ell}(\cos \theta)$$

$$\psi_{\vec{k}}(\vec{r})|_{R \rightarrow \infty} = e^{i\vec{k} \cdot \vec{r}} + \frac{e^{ikr}}{r} f(\Omega),$$

$$\frac{d\sigma}{d\Omega} = |f(\Omega)|^2, \quad \sigma = \frac{4\pi}{k^2} \sum_{\ell} (2\ell + 1) \sin^2 \delta_{\ell}, \quad \delta \approx -ak$$

$$L_{\pm} |\ell, m\rangle = \sqrt{\ell(\ell + 1) - m(m \pm 1)} |\ell, m \pm 1\rangle,$$

$$C_{m_{\ell}, m_s; JM}^{\ell, s} = \langle \ell, s, J, M | \ell, s, m_{\ell}, m_s \rangle,$$

$$\langle \tilde{\beta}, J, M | T_q^k | \beta, \ell, m_{\ell} \rangle = C_{qm_{\ell}; JM}^{k\ell} \frac{\langle \tilde{\beta}, J || T^{(k)} || \beta, \ell, J \rangle}{\sqrt{2J + 1}},$$

$$n = \frac{(2s + 1)}{(2\pi)^d} \int_{k < k_f} d^d k, \quad d \text{ dimensions,}$$

$$\{\Psi_s(\vec{x}), \Psi_s^{\dagger}(\vec{y})\} = \delta^3(\vec{x} - \vec{y}) \delta_{ss'},$$

$$\Psi_s^{\dagger}(\vec{r}) = \frac{1}{\sqrt{V}} \sum_{\vec{k}} e^{i\vec{k} \cdot \vec{r}} a_s^{\dagger}(\vec{k}), \quad \{\Psi_s(\vec{x}), a_{\alpha}^{\dagger}\} = \phi_{\alpha, s}(\vec{x}).$$

Your Name: _____

1. (15 pts) At $t = 0$ an electron is in the $|\uparrow\rangle$ (up along the z axis) state, which is represented by

$$|\uparrow\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}.$$

The evolution is determined by the Hamiltonian,

$$\mathbf{H} = A\sigma_z + B\sigma_y.$$

What is the probability the electron will be found in the $|\downarrow\rangle$ state as a function of time?

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Extra work space for #1.

2. Consider a **TWO-DIMENSIONAL** world with two types of particles, an *Aaron* particle and a *Barbara* particle. The Aaron particle in state \mathbf{a} can decay into a Barbara particle in state \mathbf{b} via the interaction,

$$\langle \text{Barbara, } \mathbf{b} | V | \text{Aaron, } \mathbf{a} \rangle = \int dx dy \psi_b^*(x, y) v_0 \psi_a(x, y).$$

The Aaron and Barbara particles have the same mass m , but Aaron particles feel a harmonic oscillator potential,

$$V_A(x, y) = \frac{1}{2} m \omega^2 (x^2 + y^2),$$

while the Barbara particles feel no such potential. An Aaron particle in the ground state of the harmonic oscillator decays into a Barbara particle. Assume that v_0 is sufficiently small that Fermi's golden rule can be applied.

- (a) (5 pts) What is the magnitude of \mathbf{k} , the outgoing momentum wave vector of the Barbara particle?
- (b) (15 pts) What is the matrix element $\langle \text{Barbara, } \mathbf{b} | V | \text{Aaron, } \mathbf{a} \rangle$? Barbara's state \mathbf{b} is a plane wave of wave number \vec{k} ,

$$\psi_b = \frac{e^{i\vec{k}\cdot\vec{r}}}{\sqrt{A}}, \quad A = \text{area} \rightarrow \infty.$$

Aaron particle's state \mathbf{a} refers to the ground state of the harmonic oscillator,

$$\psi_a = \frac{1}{\pi^{1/2} b} e^{-(x^2+y^2)/2b^2}, \quad b = \sqrt{\frac{\hbar}{m\omega}}.$$

- (c) (25 pts) What is the decay rate of an Aaron? (You will be penalized if your answer is dimensionally inconsistent – note that v_0 has dimensions of energy)

Your Name: _____

Extra work space for #2.

3. A point charge Ze is placed at a point $\vec{r} = \mathbf{0}$, and the differential cross section is measured with a beam of electrons of wave number \mathbf{k} moving along the z axis. The Rutherford differential cross section is

$$\left(\frac{d\sigma}{d\Omega}\right)_{\text{point}} = \alpha = \frac{m^2 Z^2 e^4}{(\hbar k)^4 (1 - \cos \theta)^2}.$$

Now, that same charge Ze is spread out uniformly along a line from $z = -a$ to $z = a$. I.e. the charge density is

$$\rho(x, y, z) = \begin{cases} 0, & z < -a \\ \delta(x)\delta(y)\frac{Ze}{2a}, & -a < z < a \\ 0, & a < z \end{cases}.$$

The cross section is measured again.

- (a) (15 pts) What is the differential cross section? (Express answer in terms of α , \mathbf{k} , and \mathbf{a})

$$\left(\frac{d\sigma}{d\Omega}\right)_{\text{line}} = ???$$

- (b) (5 pts) What is $(d\sigma/d\Omega)_{\text{line}}$ in the limit that $\mathbf{k} \ll 1/a$?
 (c) (10 pts) At what scattering angles does the differential cross section disappear?

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Extra work space for #3.