

$$\int_{-\infty}^{\infty} dx e^{-x^2/(2a^2)} = a\sqrt{2\pi},$$

$$H = i\hbar\partial_t, \quad \vec{P} = -i\hbar\nabla,$$

$$\sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \quad \sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix},$$

$$U(t, -\infty) = 1 + \frac{-i}{\hbar} \int_{-\infty}^t dt' V(t')U(t', -\infty),$$

$$\langle x|x'\rangle = \delta(x - x'), \quad \langle p|p'\rangle = \frac{1}{2\pi\hbar}\delta(p - p'),$$

$$|p\rangle = \int dx |x\rangle e^{ipx/\hbar}, \quad |x\rangle = \int \frac{dp}{2\pi\hbar} |p\rangle e^{-ipx/\hbar},$$

$$H = \frac{P^2}{2m} + \frac{1}{2}m\omega^2 x^2 = \hbar\omega(a^\dagger a + 1/2),$$

$$a^\dagger = \sqrt{\frac{m\omega}{2\hbar}}X - i\sqrt{\frac{1}{2\hbar m\omega}}P,$$

$$\psi_0(x) = \frac{1}{(\pi b^2)^{1/4}} e^{-x^2/2b^2}, \quad b^2 = \frac{\hbar}{m\omega},$$

$$\rho(\vec{r}, t) = \psi^*(\vec{r}_1, t_1)\psi(\vec{r}_2, t_2)$$

$$\vec{j}(\vec{r}, t) = \frac{-i\hbar}{2m}(\psi^*(\vec{r}, t)\nabla\psi(\vec{r}, t) - (\nabla\psi^*(\vec{r}, t))\psi(\vec{r}, t))$$

$$- \frac{e\vec{A}}{mc}|\psi(\vec{r}, t)|^2.$$

$$H = \frac{(\vec{P} - e\vec{A}/c)^2}{2m} + e\Phi,$$

For  $V = \beta\delta(x - y)$ :  $-\frac{\hbar^2}{2m}\left(\frac{\partial}{\partial x}\psi(x)|_{y+\epsilon} - \frac{\partial}{\partial x}\psi(x)|_{y-\epsilon}\right) = -\beta\psi(y),$

$$\vec{E} = -\nabla\Phi - \frac{1}{c}\partial_t\vec{A}, \quad \vec{B} = \nabla \times \vec{A},$$

$$\omega_{\text{cyclotron}} = \frac{eB}{mc},$$

$$e^{A+B} = e^A e^B e^{-C/2}, \quad \text{if } [A, B] = C, \text{ and } [C, A] = [C, B] = 0,$$

$$Y_{0,0} = \frac{1}{\sqrt{4\pi}}, \quad Y_{1,0} = \sqrt{\frac{3}{4\pi}}\cos\theta, \quad Y_{1,\pm 1} = \mp\sqrt{\frac{3}{8\pi}}\sin\theta e^{\pm i\phi},$$

$$Y_{2,0} = \sqrt{\frac{5}{16\pi}}(3\cos^2\theta - 1), \quad Y_{2,\pm 1} = \mp\sqrt{\frac{15}{8\pi}}\sin\theta\cos\theta e^{\pm i\phi},$$

$$Y_{2,\pm 2} = \sqrt{\frac{15}{32\pi}}\sin^2\theta e^{\pm 2i\phi}, \quad Y_{\ell-m}(\theta, \phi) = (-1)^m Y_{\ell m}^*(\theta, \phi).$$

$$|N\rangle = |n\rangle - \sum_{m \neq n} |m\rangle \frac{1}{\epsilon_m - \epsilon_n} \langle m|V|n\rangle + \dots$$

$$E_N = \epsilon_n + \langle n|V|n\rangle - \sum_{m \neq n} \frac{|\langle m|V|n\rangle|^2}{\epsilon_m - \epsilon_n}$$

$$j_0(x) = \frac{\sin x}{x}, \quad n_0(x) = -\frac{\cos x}{x}, \quad j_1(x) = \frac{\sin x}{x^2} - \frac{\cos x}{x}, \quad n_1(x) = -\frac{\cos x}{x^2} - \frac{\sin x}{x}$$

$$j_2(x) = \left(\frac{3}{x^3} - \frac{1}{x}\right) \sin x - \frac{3}{x^2} \cos x, \quad n_2(x) = -\left(\frac{3}{x^3} - \frac{1}{x}\right) \cos x - \frac{3}{x^2} \sin x,$$

$$\frac{d}{dt} P_{i \rightarrow n}(t) = \frac{2\pi}{\hbar} |V_{ni}|^2 \delta(E_n - E_i),$$

$$\frac{d\sigma}{d\Omega} = \frac{m^2}{4\pi^2 \hbar^4} \left| \int d^3r \mathcal{V}(r) e^{i(\vec{k}_f - \vec{k}_i) \cdot \vec{r}} \right|^2,$$

$$\sigma = \frac{(2S_R + 1) 4\pi (\hbar\Gamma_R/2)^2}{(2S_1 + 1)(2S_2 + 1) k^2 (\epsilon_k - \epsilon_r)^2 + (\hbar\Gamma_R/2)^2},$$

$$\frac{d\sigma}{d\Omega} = \left( \frac{d\sigma}{d\Omega} \right)_{\text{single}} \tilde{S}(\vec{q}), \quad \tilde{S}(\vec{q}) = \left| \sum_{\vec{a}} e^{i\vec{q} \cdot \vec{a}} \right|^2,$$

$$e^{i\vec{k} \cdot \vec{r}} = \sum_{\ell} (2\ell + 1) i^{\ell} j_{\ell}(kr) P_{\ell}(\cos \theta),$$

$$P_{\ell}(\cos \theta) = \sqrt{\frac{4\pi}{2\ell + 1}} Y_{\ell, m=0}(\theta, \phi),$$

$$P_0(x) = 1, \quad P_1(x) = x, \quad P_2(x) = (3x^2 - 1)/3,$$

$$f(\Omega) \equiv \sum_{\ell} (2\ell + 1) e^{i\delta_{\ell}} \sin \delta_{\ell} \frac{1}{k} P_{\ell}(\cos \theta)$$

$$\psi_{\vec{k}}(\vec{r})|_{R \rightarrow \infty} = e^{i\vec{k} \cdot \vec{r}} + \frac{e^{ikr}}{r} f(\Omega),$$

$$\frac{d\sigma}{d\Omega} = |f(\Omega)|^2, \quad \sigma = \frac{4\pi}{k^2} \sum_{\ell} (2\ell + 1) \sin^2 \delta_{\ell}, \quad \delta \approx -ak$$

$$L_{\pm} |\ell, m\rangle = \sqrt{\ell(\ell + 1) - m(m \pm 1)} |\ell, m \pm 1\rangle,$$

$$C_{m_{\ell}, m_s; JM}^{\ell, s} = \langle \ell, s, J, M | \ell, s, m_{\ell}, m_s \rangle,$$

$$\langle \tilde{\beta}, J, M | T_q^k | \beta, \ell, m_{\ell} \rangle = C_{qm_{\ell}; JM}^{k\ell} \frac{\langle \tilde{\beta}, J || T^{(k)} || \beta, \ell, J \rangle}{\sqrt{2J + 1}},$$

$$n = \frac{(2s + 1)}{(2\pi)^d} \int_{k < k_f} d^d k, \quad d \text{ dimensions,}$$

$$\{\Psi_s(\vec{x}), \Psi_{s'}^{\dagger}(\vec{y})\} = \delta^3(\vec{x} - \vec{y}) \delta_{ss'},$$

$$\Psi_s^{\dagger}(\vec{r}) = \frac{1}{\sqrt{V}} \sum_{\vec{k}} e^{i\vec{k} \cdot \vec{r}} a_s^{\dagger}(\vec{k}), \quad \{\Psi_s(\vec{x}), a_{\alpha}^{\dagger}\} = \phi_{\alpha, s}(\vec{x}).$$

1. (5 pts) Consider three spin operators  $S_x, S_y$  and  $S_z$ . Circle the operators that commute with  $S_z$ .

- $S_x$
- $S_z$
- $S_x^2$
- $S_z^2$
- $S_x^2 + S_y^2 + S_z^2$

2. (5 pts) Consider two sets of spin operators,  $S_x, S_y, S_z$  and  $L_x, L_y, L_z$ . You can assume  $\vec{S}$  operates on intrinsic spin and that  $\vec{L}$  describes orbital angular momentum. Circle the operators that commute with  $S_z$ .

- $L_x$
- $L_z$
- $L_x^2$
- $L_z^2$
- $L_x^2 + L_y^2 + L_z^2$

3. (5 pts) Now consider the operators  $\vec{J} \equiv \vec{L} + \vec{S}$ . Circle the operators that commute with  $S_z$ .

- $J_x$
- $J_z$
- $J_x^2$
- $J_z^2$
- $J_x^2 + J_y^2 + J_z^2$

4. (A proton and a neutron are in the ground state of a harmonic oscillator. An interaction is added,

$$V_{s.s.} = -\alpha \vec{S}_p \cdot \vec{S}_n$$

At  $t = 0$  the proton is in a  $|\uparrow\rangle$  state and the neutron is in a  $|\downarrow\rangle$  state, which we label as  $|\uparrow, \downarrow\rangle$ . With this labeling the first spin refers to the proton and the second to the neutron.

- (a) (15 pts) In the basis above, express  $V_{s.s.}$  as a  $4 \times 4$  matrix. Use a basis where the states are expressed as

$$|\uparrow, \uparrow\rangle = \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix}, \quad |\uparrow, \downarrow\rangle = \begin{pmatrix} 0 \\ 1 \\ 0 \\ 0 \end{pmatrix}, \quad |\downarrow, \uparrow\rangle = \begin{pmatrix} 0 \\ 0 \\ 1 \\ 0 \end{pmatrix}, \quad |\downarrow, \downarrow\rangle = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 1 \end{pmatrix}.$$

- (b) (15 pts) Find the probability that the pair is each of the following states as a function of time for  $t > 0$ .
- $|\uparrow, \uparrow\rangle$
  - $|\uparrow, \downarrow\rangle$  (this is the state at  $t = 0$ )
  - $|\downarrow, \uparrow\rangle$
  - $|\downarrow, \downarrow\rangle$

(Extra work space for #4)

5. A beam of spinless particles of mass  $m$  and kinetic energy  $E$  is aimed at a spherically symmetric repulsive potential

$$V(r) = \begin{cases} V_0, & r < a \\ 0, & r > a \end{cases}$$

Assume  $E < V_0$ .

- (a) (10 pts) Find the  $\ell = 0$  phase shift as a function of the incoming wave number  $k$ .
- (b) (5 pts) What is the cross section as  $k \rightarrow 0$ ?
- (c) (10 pts) What is the relative probability density for a particle in the wave packet to be at the origin compared to the probability with no potential? I.e. If  $\rho_0$  is the probability density at  $r = 0$  in the absence of the potential and  $\rho$  is the density with the potential, find  $\rho/\rho_0$ .

(Extra work space for #5)

6. A particle of mass  $m$  moves in a one-dimensional attractive potential

$$V(x) = -V_0 \exp(-x^2/2a^2).$$

Use a gaussian form for a trial wave function,

$$\langle x|\mathbf{b}\rangle = \psi_{\mathbf{b}}(x) = \frac{1}{(\pi b^2)^{1/4}} e^{-x^2/2b^2},$$

where  $\mathbf{b}$  is the variational parameter.

- (a) (10 pts) What is  $\langle \mathbf{b}|\mathbf{KE}|\mathbf{b}\rangle$ ? –the expectation of the kinetic energy.
- (b) (10 pts) What is  $\langle \mathbf{b}|\mathbf{V}|\mathbf{b}\rangle$ ? –the expectation of the potential energy.
- (c) (10 pts) Find an expression that when solved for  $\mathbf{b}$  and then plugged into (a) and (b) provides an estimate of the energy. This expression can be a polynomial that needs to be solved for  $\mathbf{b}$ . (No credit will be given for expressions that are dimensionally inconsistent)

(Extra work space for #6)