

# Nuclear Equation of State & Weak Interactions around neutrino-sphere

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# Outline

- Nuclear symmetry energy
- Nuclear equation of state (EOS)
- Neutrino interactions around neutrino-sphere
- Summary & Outlook

# Nuclear Symmetry Energy in neutron rich matter

$$E(n_B, x_p) = E(n_B, x_p = 1/2) + E_{\text{sym}}(n_B)\delta^2 + \dots \quad \delta = (1 - 2x_p)$$

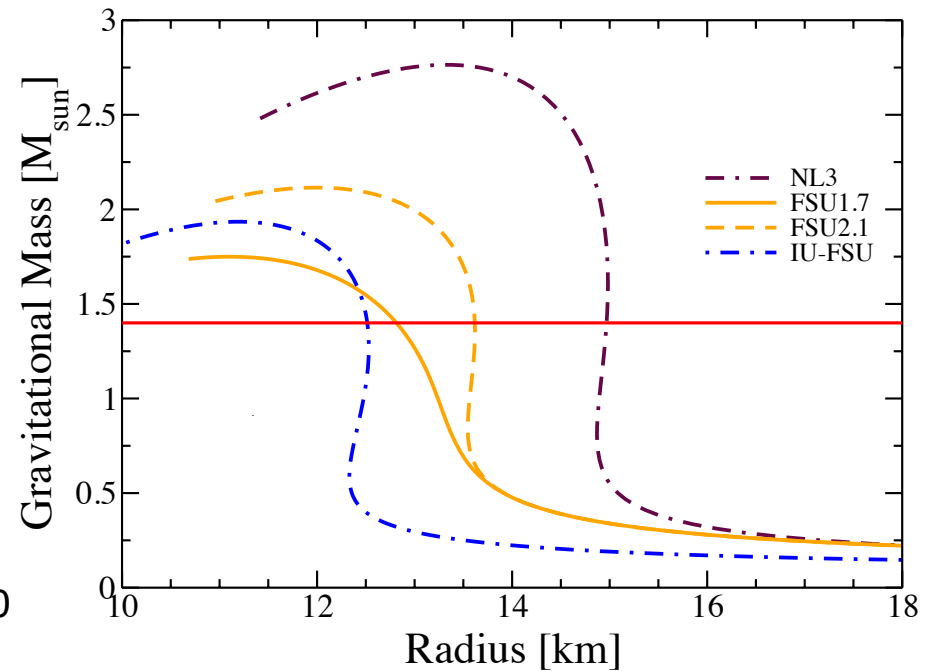
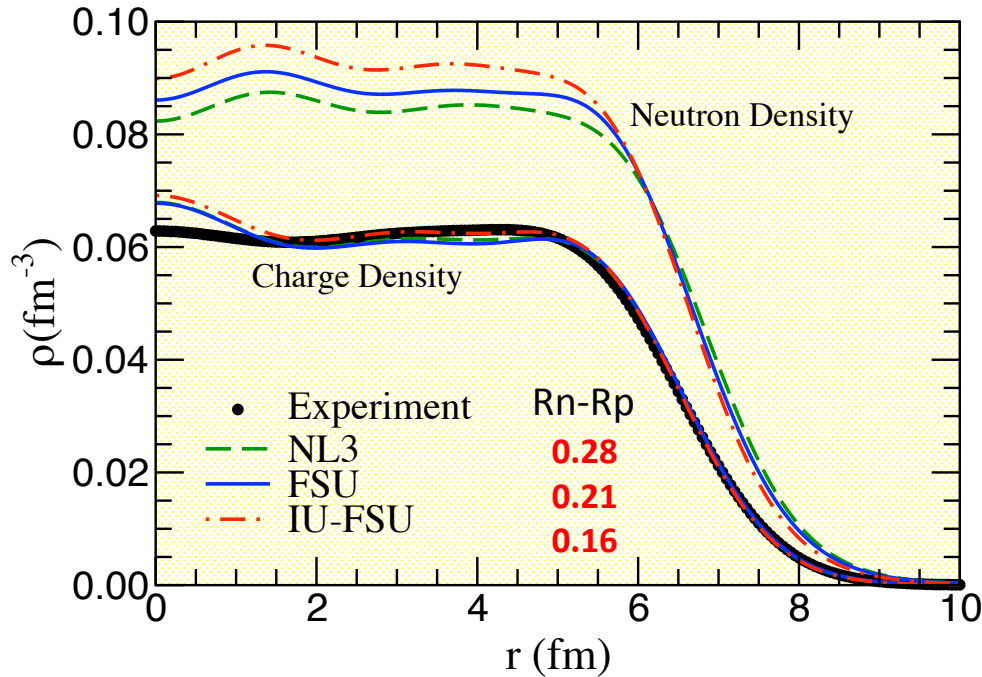
$x_p$  : proton fraction

- $E_{\text{sym}}$  describes how energy changes with p-n asymmetry
- Pressure  $\sim E'_{\text{sym}} = L/3$ , or large  $E'_{\text{sym}}(n_0)$  – stiff EOS
- Large uncertainties in  $E_{\text{sym}}$  and  $E'_{\text{sym}}$  - reflected in diff. EOS

known correlations due to efforts in last decade:

- $E'_{\text{sym}} \sim$  neutron radius of neutron rich nucleus, eg. Pb208
  - $E'_{\text{sym}} \sim$  neutron star radius
  - $E'_{\text{sym}} \sim$  Proto-Neutron Star (PNS) cooling
  - $E_{\text{sym}} \sim$  neutrino interactions
- Talks by S. Reddy, L. Roberts*  
*Talks by T. Fischer, L. Roberts*

Models	NL3	FSUGold	IU-FSU
$n_0 E'_{\text{sym}}(n_0)$ [MeV]	39	20	16



- A larger  $E'_{\text{sym}}(n_0)$  indicates a bigger radius for 1.4 solar mass neutron star and a bigger neutron radius in 208Pb. Fattoyev, Horowitz, Piekarewicz, GS (2010)
- Experiments/observations: PREX in JLab / model & data analysis on neutron star x-ray bursts/x-ray binary; (others, like isospin effects in heavy ion collision)
- Recent review: Lattimer & Lim (2012), Steiner & Gandolfi (2011)

# Nuclear EOS in a unified framework

- One flexible framework for both nuclei and dense nuclear matter
- Ideal: density functional theory, for example, UNEDF
- Our choice
  - High density: relativistic mean field model for both nuclei and dense nuclear matter
  - Low density: virial expansion of nuclei and nucleons

GS, Horowitz, O'Connor (2011).  
GS, Horowitz, Teige (2011).

# High density: Relativistic Mean Field Theory

$$\begin{aligned}
 \mathcal{L}_{\text{int}} = & \bar{\psi} \left[ \overset{\text{attractive}}{g_s \phi} - \left( \overset{\text{repulsive}}{g_v V_\mu} + \frac{g_\rho}{2} \boldsymbol{\tau} \cdot \mathbf{b}_\mu + \frac{e}{2} (1 + \tau_3) A_\mu \right) \gamma^\mu \right] \psi \\
 & - \frac{\kappa}{3!} (g_s \phi)^3 - \frac{\lambda}{4!} (g_s \phi)^4 + \frac{\zeta}{4!} g_v^4 (V_\mu V^\mu)^2 + \Lambda_v g_\rho^2 \mathbf{b}_\mu \cdot \mathbf{b}^\mu g_v^2 V_\nu V^\nu
 \end{aligned}$$

- $\Psi$ : Nucleon fields
- Meson/photon fields: classical expectation value  $\rightarrow$  mean field
- 7 adjustable parameters

## EoMs in real space:

### Nucleon:

$$\left[ \boldsymbol{\alpha} \cdot \mathbf{p} + V(\mathbf{r}) + \beta (M + S(\mathbf{r})) \right] \psi_i = \varepsilon_i \psi_i$$

### Hartree Mean fields:

$$\begin{aligned}
 V(\mathbf{r}) &= \beta \left\{ g_\omega \rho_\omega + g_\rho \vec{\tau} \cdot \vec{\rho}_\mu + e \frac{(1 + \tau_3)}{2} A_\mu + \Sigma^R \right\}, \\
 S(\mathbf{r}) &= \Gamma_\sigma \sigma,
 \end{aligned}$$

### Mesons and Photons:

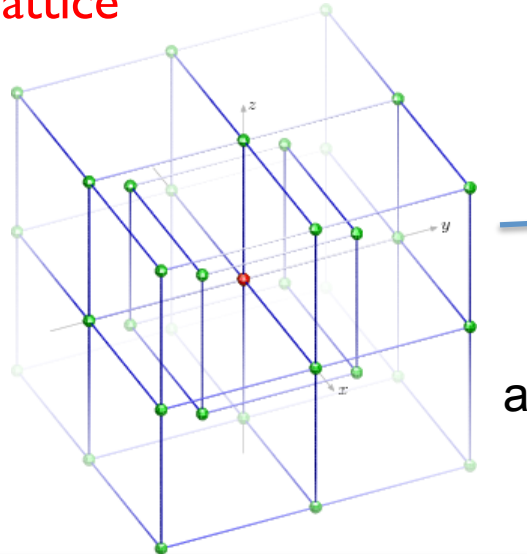
$$\left\{ \begin{aligned}
 (-\Delta + \partial_\sigma U(\sigma)) \sigma(\mathbf{r}) &= -g_\sigma \rho_s(\mathbf{r}) \\
 (-\Delta + m_\omega^2) \omega_0(\mathbf{r}) &= g_\omega \rho_v(\mathbf{r}) \\
 (-\Delta + m_\rho^2) \rho_0(\mathbf{r}) &= g_\rho \rho_3(\mathbf{r}) \\
 -\Delta A_0(\mathbf{r}) &= e(\rho_c(\mathbf{r}) - \rho_e)
 \end{aligned} \right.$$

### Finite temperature $n_i$ : Fermi-Dirac stat.

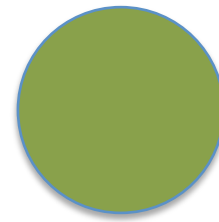
$$\left\{ \begin{aligned}
 \rho_s(\mathbf{r}) &= \sum_{i=1}^A \bar{\psi}_i(\mathbf{r}) \psi_i(\mathbf{r}) n_i, \quad \rho_3(\mathbf{r}) = \sum_{i=1}^A \psi_i^+(\mathbf{r}) \tau_3 \psi_i(\mathbf{r}) n_i \\
 \rho_v(\mathbf{r}) &= \sum_{i=1}^A \psi_i^+(\mathbf{r}) \psi_i(\mathbf{r}) n_i, \quad \rho_c(\mathbf{r}) = \sum_{i=1}^A \psi_i^+(\mathbf{r}) \frac{1 - \tau_3}{2} \psi_i(\mathbf{r}) n_i
 \end{aligned} \right.$$

# proto-neutron star crust ( $\Gamma \gg 1$ )

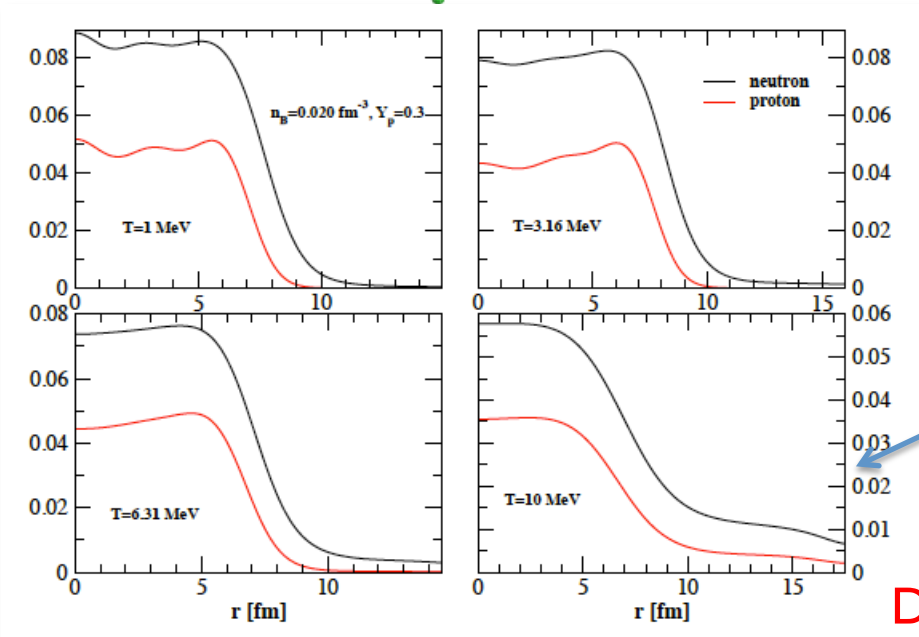
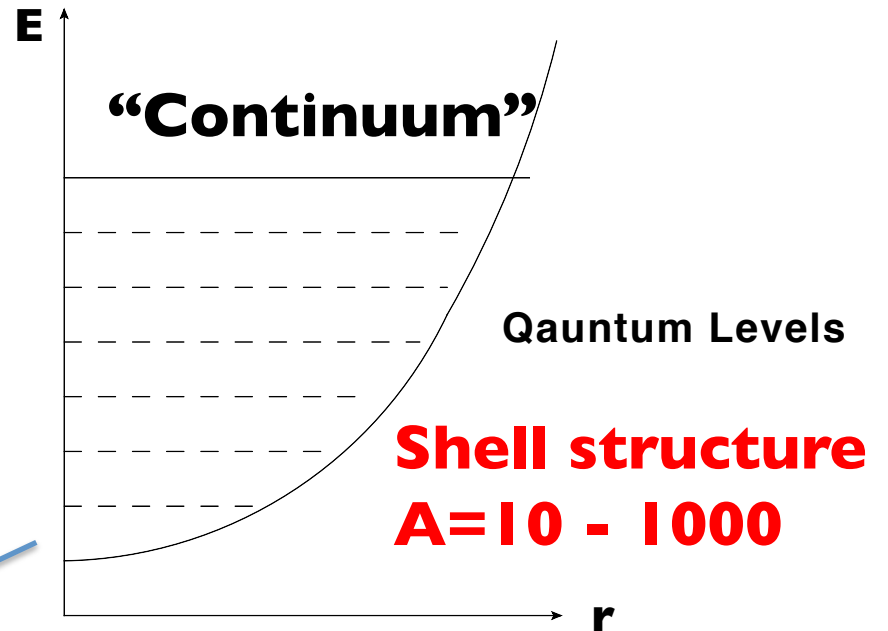
BCC lattice



$$\Gamma = (Ze)^2 / ak_B T$$



spherical Wigner-Seitz cell



Density distribution in the unit lattice

# Low density: Ensemble of nucleons and nuclei

- Grand partition function in virial expansion

$$\frac{\log Q}{V} = \frac{P}{T} = \frac{2}{\lambda_n^3} [z_n + z_p + (z_p^2 + z_n^2)b_n + 2z_p z_n b_{pn}] \quad \blacktriangleright \text{nucleon-nucleon}$$

$$+ \frac{1}{\lambda_\alpha^3} [z_\alpha + z_\alpha^2 b_\alpha + 2z_\alpha (z_n + z_p) b_{\alpha n}] \quad \blacktriangleright \text{nucleon-alpha}$$

$$+ \sum_i \frac{1}{\lambda_i^3} z_i \Omega_i \quad \blacktriangleright \text{nuclei}$$

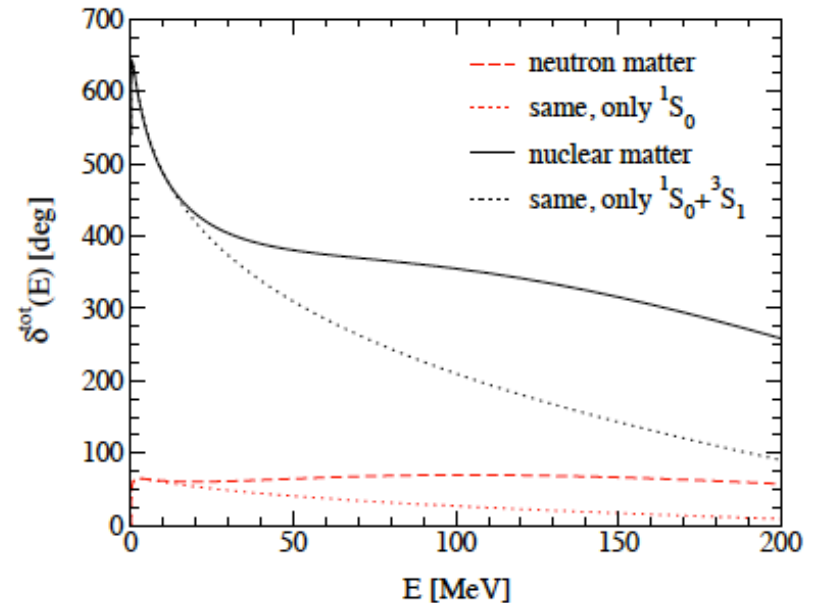
I. Nucleon and alpha: mod. ind.

Horowitz, Schwenk '05

2<sup>nd</sup> virial coef.  $b_2$  determined from elastic scattering phase shifts: model-independent

- $b_n$  for n-n (p-p), –  $b_{pn}$  for n-p,
- $b_\alpha$  for alpha-alpha, –  $b_{\alpha-n}$  for alpha-n

$$b_n(T) = \frac{1}{2^{1/2} \pi T} \int_0^\infty dE e^{-E/2T} \delta_n^{\text{tot}}(E) - 2^{-5/2}.$$





# Low density: Ensemble of nucleons and nuclei

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Horowitz, Schwenk '05

2. Heavy species: 8980 nuclei

FRDM mass table: Moller et al '97.  $\Delta_{\text{rms}} \sim 0.6 \text{ MeV}$

Chemical equilibrium -----

$$\mu_i = Z\mu_p + N\mu_n \quad z_i = z_p^Z z_n^N e^{(E_i - E_i^C)/T}$$

Coulomb correction -----

$$E_i^C = \frac{3}{5} \frac{Z_i^2 \alpha}{r_A} \left[ -\frac{3}{2} \frac{r_A}{r_i} + \frac{1}{2} \left( \frac{r_A}{r_i} \right)^3 \right]$$

Nuclear partition function  $\Omega_i$  ----

eg, Fowler, Engelbrecht, Woosley, '78

Nuclei spacing

$$\frac{4}{3} \pi r_i^3 \left( \sum_j Z_j n_j \right) = Z_i$$

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3. Solving:  $z_n, z_p, r_i$ :

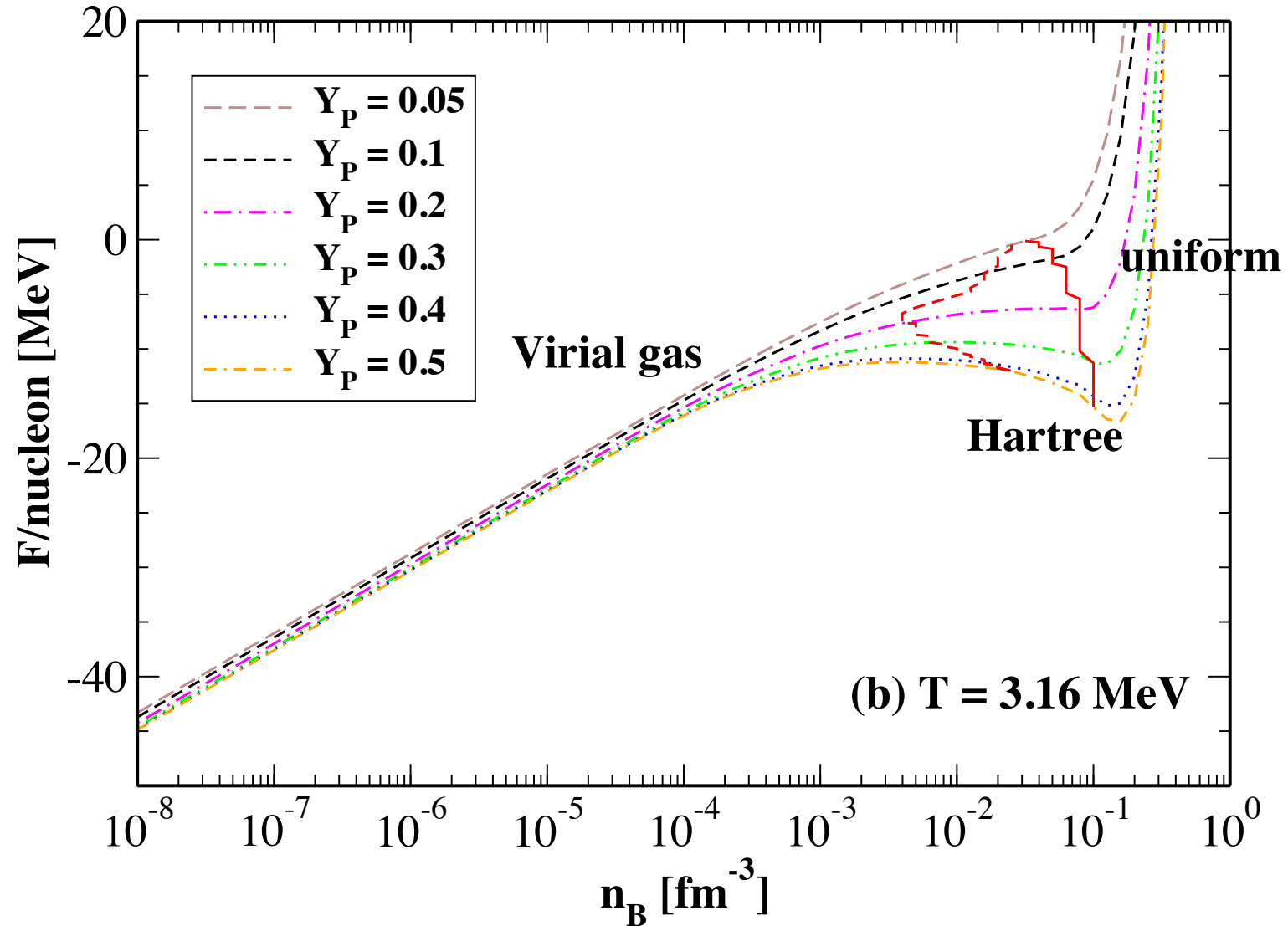
$$n_B = n_n + n_p + 4n_\alpha + \sum_i A_i n_i = \text{Baryon}$$

$$Y_P = (n_p + 2n_\alpha + \sum_i Z_i n_i) / n_B = \text{Charge}$$

4. Mass fraction:

$$X_a = A_a n_a / n_B$$

# Matching EOS from low to high densities



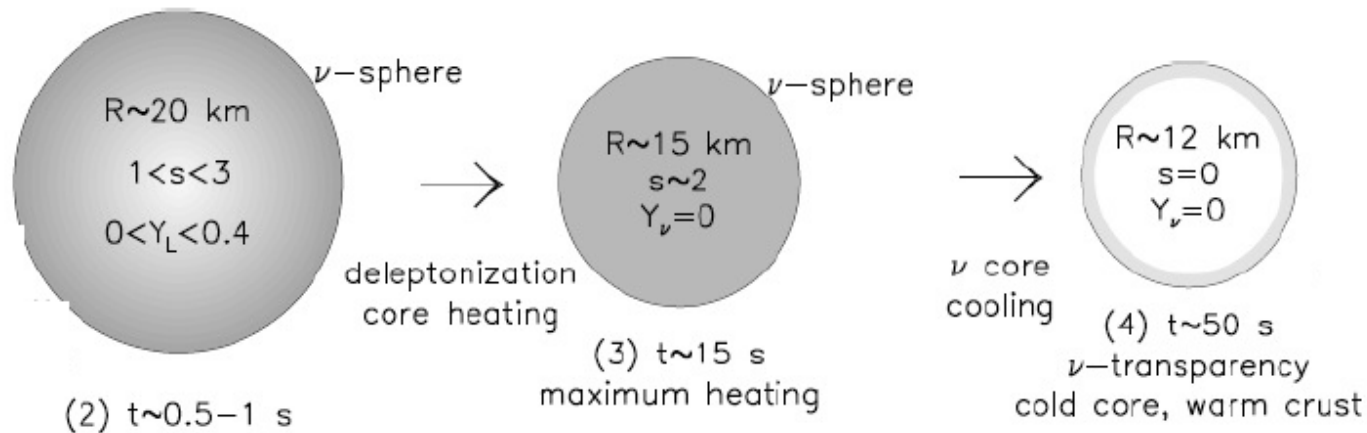
### 3-D Parameter spaces ( $T, \rho, Y_p$ ) for gas-solid-liquid

	Virial Gas	Hartree	Uniform matter
Temperature [MeV]	0.1~20	0.1~12	0.1~80
Density [fm <sup>-3</sup> ]	1E-8~1E-1	1E-4~1E-1	1E-8~1.6
Proton fraction	0.05~0.56	0.05~0.56	0~0.56
# points in phase space	73,840	17,021	90,478
CPU time (hr)	10,000	100,000	100

- Matching them and interpolate to get thermodynamically consistent table.
- EOS tables: [http://cecelia.physics.indiana.edu/gang\\_shen\\_eos/](http://cecelia.physics.indiana.edu/gang_shen_eos/)

# Proto-neutron star (PNS) neutrinos and EOS

- Neutrino transport drives the evolution of a PNS (hot, lepton rich, SN remnant)  $\rightarrow$  NS (cold, deleptonized, compact)



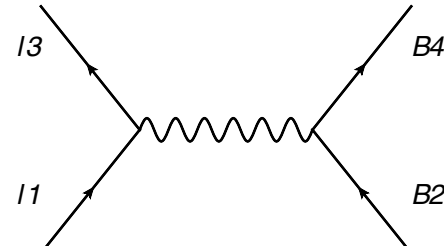
- Neutrino sphere: “effective boundary” for neutrinos to decouple; dynamical.
- Matter around neutrino-sphere important for spectra of “emitted” neutrino in neutrino driven wind

# Neutrino-nucleon scattering

- Neutrinos couple to nuclear (isospin) density and (isospin) spin density:

V-A theory:  $j^\mu(x) = \bar{\psi}(x)\gamma^\mu(C_V - C_A\gamma_5)(\tau_j)\psi(x)$

Non Relativistic limit:  $\rightarrow C_V\psi^\dagger(\tau_j)\psi\delta^{\mu 0} - C_A\psi^\dagger\sigma^i(\tau_j)\psi\delta^{\mu i}$

$$\frac{1}{V} \frac{d^2\sigma}{d\cos\theta dE_3} = \frac{G_F^2}{4\pi^2} E_3^2 (1 - f_3(E_3))$$


$$\times [C_V^2 (1 + \cos\theta) S_\rho(q_0, q) + C_A^2 (3 - \cos\theta) S_\sigma(q_0, q)]$$

- Free Fermi gas response functions: peaked around  $q_0 \sim 0$

$$S_{\rho,\sigma}(q_0, q) = \frac{1}{2\pi^2} \int d^3p_2 \delta(q_0 + E_2 - E_4) f_2(1 - f_4),$$

Dispersion  $E_i(p) = M_i + p^2/2M_i$

# Mean field (MF) shifts in charged current reactions

Roberts (2012), Roberts, Reddy, GS (2012)

Martinez-Pinedo et al. (2012)

- Dispersion relation of nucleon quasi-particle from MF EOS:

$$E_i(k) = \sqrt{k^2 + M^{*2}} + U_i,$$
$$\Delta U = U_n - U_p = \alpha * (n_n - n_p) \quad \sim \text{symmetry energy}$$

- Response function is peaked around  $q_0 \approx -\Delta U$ , for  $\nu_e + n \rightarrow p + e$

- Around  $\nu$ -sphere, mean field shift is comparable to temperature  $\rightarrow$  increases phase space in final state of electron exponentially !

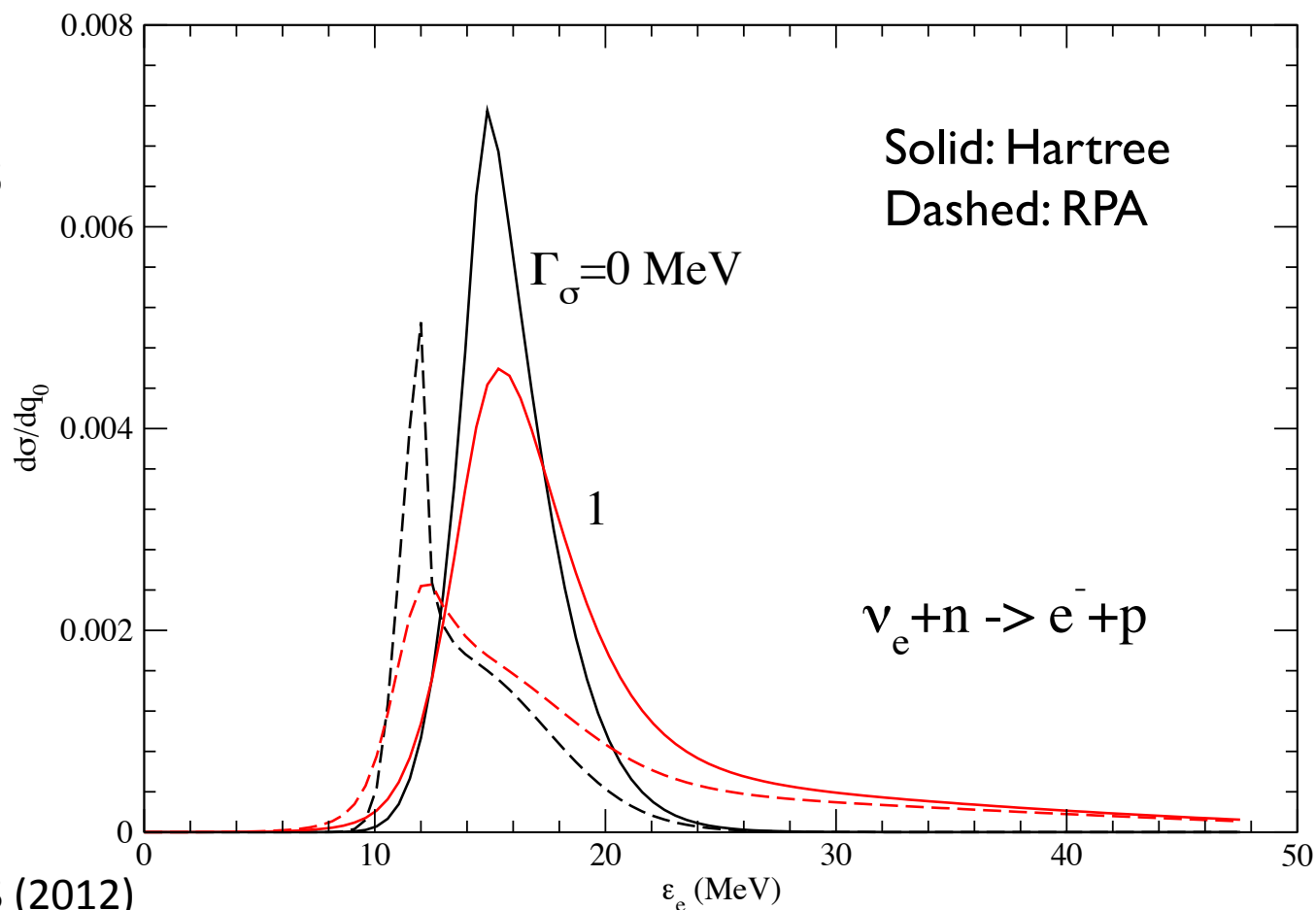
$$E_e = E_\nu + \Delta U \quad d\sigma \sim E_e^2 (1 - f_e(E_e))$$

- $\nu_e$  decouple at lower temp., while  $\bar{\nu}_e$  at higher temp. May reduce electron fraction in neutrino driven wind – help r-p nucleosynthesis

$$Y_{e,\text{NDW}} \approx \left[ 1 + \frac{\dot{N}_{\bar{\nu}_e} \langle \sigma(\epsilon)_{p,\bar{\nu}_e} \rangle}{\dot{N}_{\nu_e} \langle \sigma(\epsilon)_{n,\nu_e} \rangle} \right]^{-1}$$

# Charged current reactions in RPA and beyond

T=8 MeV  
N=0.006 fm<sup>-3</sup>  
Y<sub>e</sub>=0.03  
 $\mu_e=30$  MeV  
 $E_\nu=12$  MeV

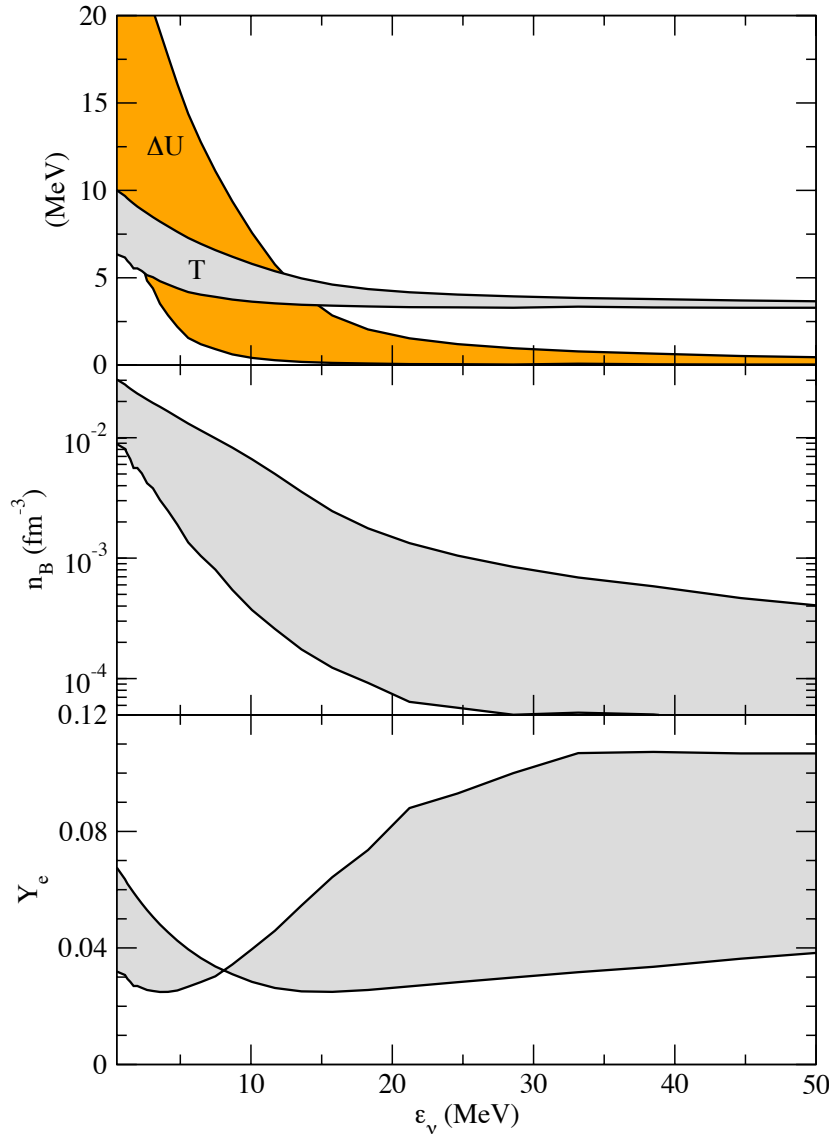


Roberts, Reddy, GS (2012)

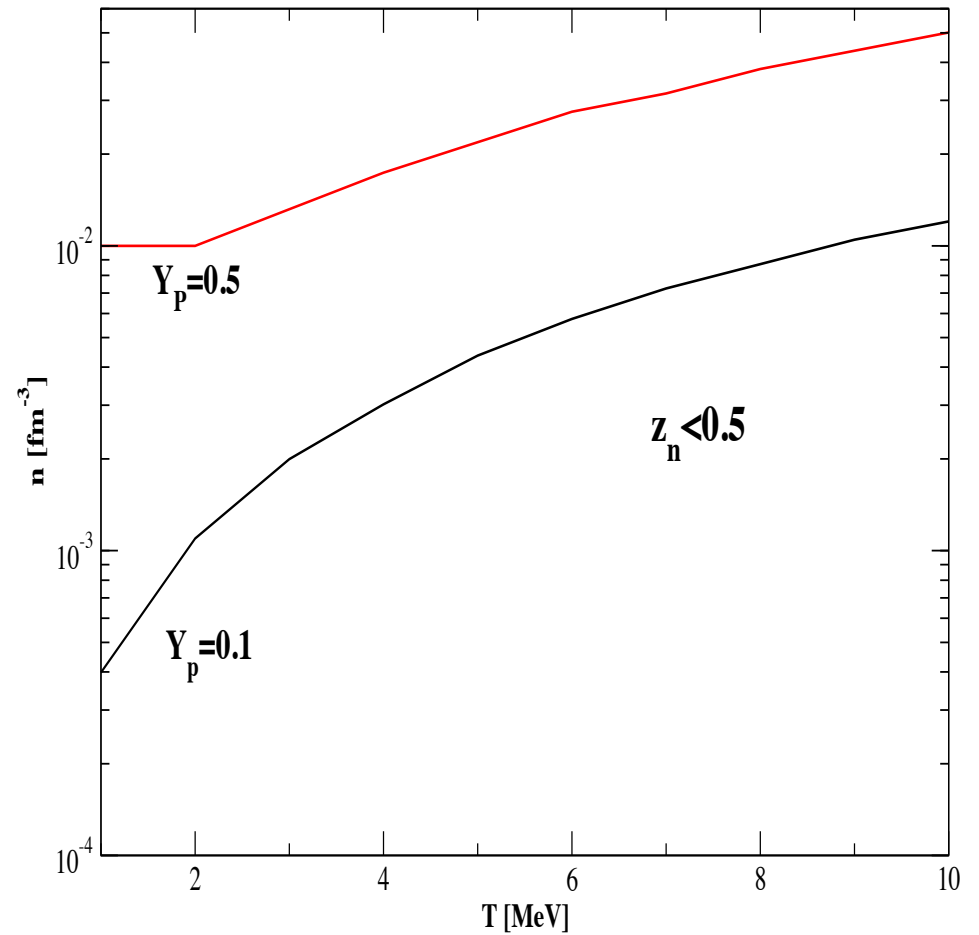
- RPA suppresses response and shifts its strength via collective mode. Multi-particle dynamics enhances it.
- Net effect mild suppression. Follow SN trajectory to study details.



# Typical conditions around nu-sphere



Courtesy: Luke Roberts



Virial expansion of  $n, p$  to 2<sup>nd</sup> order

# Virial expansion for MF shifts

- Virial EOS of n, p: 
$$P = \frac{2T}{\lambda^3} \{ z_n + z_p + (z_n^2 + z_p^2) b_n + 2z_p z_n b_{pn} \},$$
  - 2<sup>nd</sup> virial coeff.  $b_n, b_{pn}$  from mod. ind. scattering phase shifts **Horowitz, Schwenk '05**
  - Suitable for matter around neutrino-sphere
- Single particle energy (MF) shift I, II:

I. 
$$E_n = \left( \frac{\partial f}{\partial n_n} \right)_{n_p} = T \ln \frac{n_n \lambda^3}{2} - \lambda^3 T (n_n b_n + n_p b_{pn}), \quad (16)$$

$$U_n = E_n - E_n^0 = -\lambda^3 T (n_n \hat{b}_n + n_p b_{pn}), \quad (19)$$

$$E_p = \left( \frac{\partial f}{\partial n_p} \right)_{n_n} = T \ln \frac{n_p \lambda^3}{2} - \lambda^3 T (n_p b_n + n_n b_{pn}). \quad (17)$$

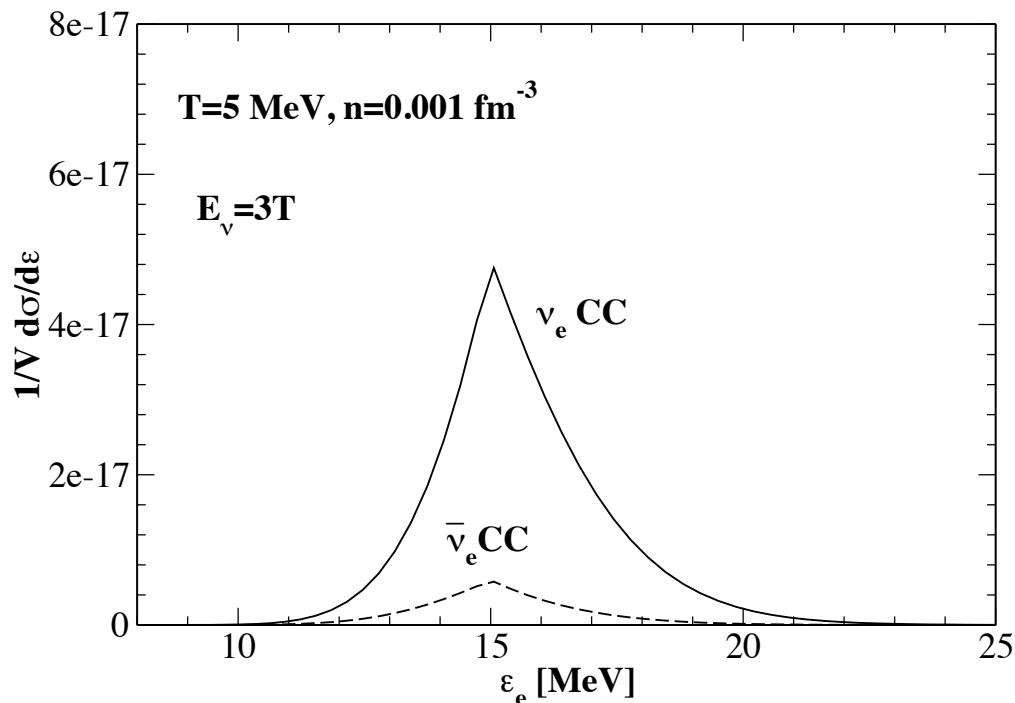
$$U_p = E_p - E_p^0 = -\lambda^3 T (n_p \hat{b}_n + n_n b_{pn}), \quad (20)$$

$$\Delta U = U_n - U_p = \lambda^3 T (n_n - n_p) (b_{pn} - \hat{b}_n)$$

II. 
$$U_n = \mu_n - \mu_n^f = -\lambda^3 T (n_n \hat{b}_n + n_p b_{pn}) + O(n_i^2).$$

# Virial expansion for MF shifts

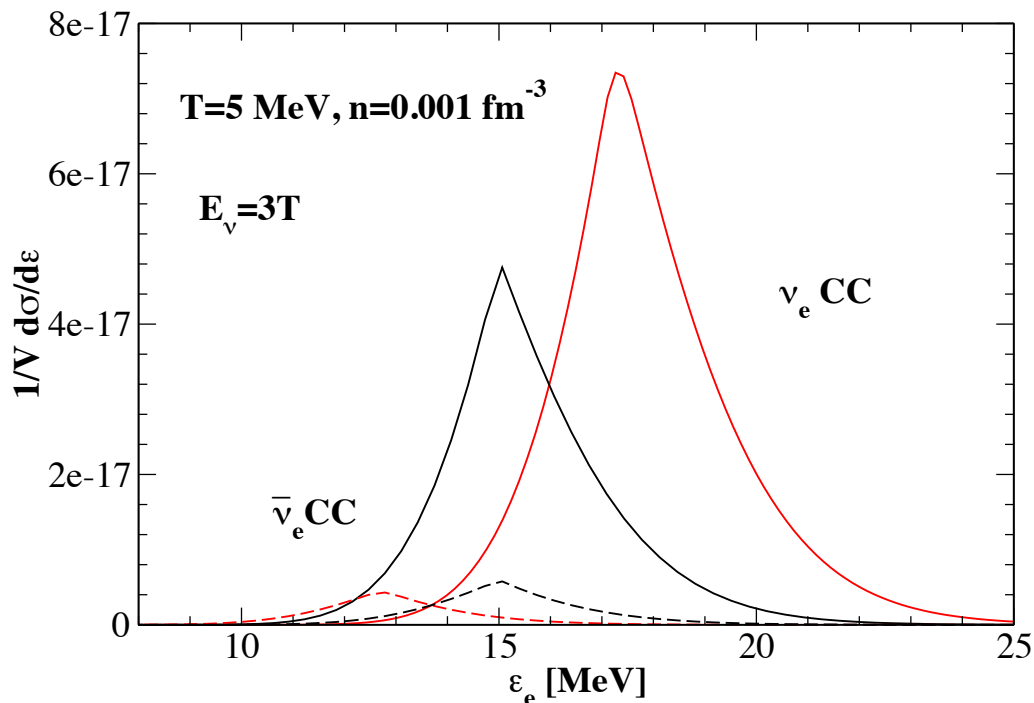
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- W/O MF shift:  
 $d\sigma \sim$  density of n or p

# Virial expansion for MF shifts

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- W MF shift

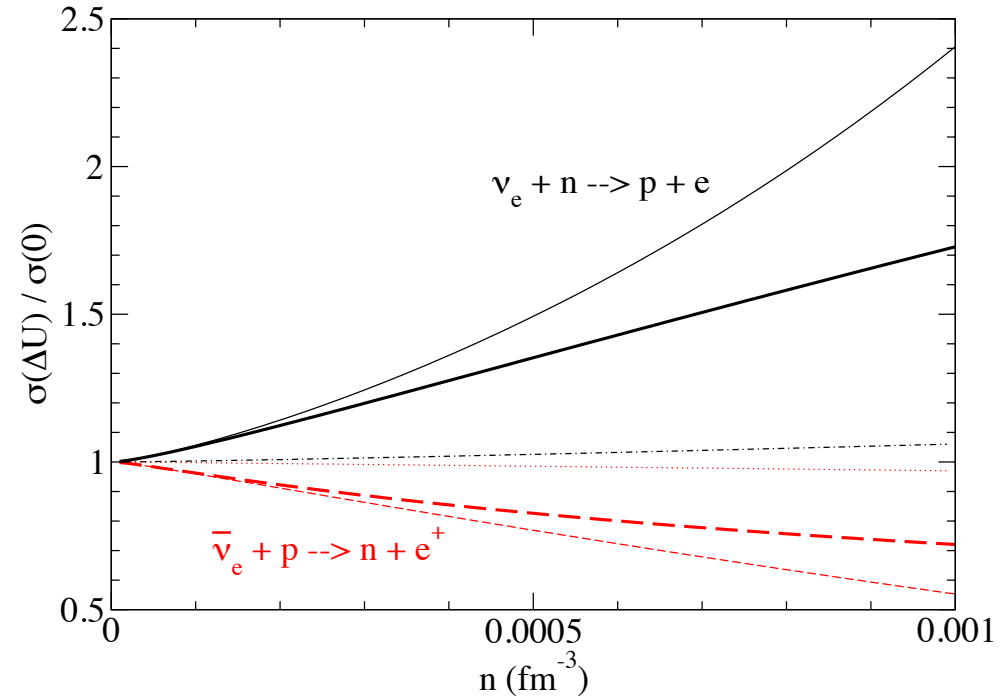
# Virial expansion and other MF models predictions

- ratio of cross sections

$$\frac{\sigma_{\nu_e}(\Delta U)}{\sigma_{\nu_e}(0)} = \frac{(E_\nu + \Delta U)^2 [1 - f(E_\nu + \Delta U)]}{E_\nu^2 [1 - f(E_\nu)]}$$

$$\frac{\sigma_{\bar{\nu}_e}(\Delta U)}{\sigma_{\bar{\nu}_e}(0)} = \frac{(E_{\bar{\nu}} - \Delta U)^2 \Theta(E_{\bar{\nu}} - \Delta U)}{E_{\bar{\nu}}^2}$$

- effect larger than MF, due to n-p correlation
- May reduce electron fractions in neutrino driven wind compared to MF
- Feedback: larger  $Y_e$  in  $\nu$ -sphere from larger  $E_{\text{sym}}$
- Need consistent simulation with EOS and virial response



$Z < 0.16$

# Summary & Outlook

- Unified framework for equations of state of nuclei and dense nuclear matter, for extreme astrophysics scenarios.
- Neutrino spectra from PNS receives medium correction, which is dependent on EOS ( $E_{\text{sym}}$ ).

*Thank you !*